STAR - Laplacian Spectral Kernels and Distances for Geometry Processing and Shape Analysis

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Abstract

In geometry processing and shape analysis, several applications have been addressed through the properties of the spectral kernels and distances, such as commute-time, biharmonic, diffusion, and wave distances. Our survey is intended to provide a background on the properties, discretization, computation, and main applications of the Laplace-Beltrami operator, the associated differential equations (e.g., harmonic equation, Laplacian eigenproblem, diffusion and wave equations), Laplacian spectral kernels and distances (e.g., commute-time, biharmonic, wave, diffusion distances). While previous work has been focused mainly on specific applications of the aforementioned topics on surface meshes, we propose a general approach that allows us to review Laplacian kernels and distances on surfaces and volumes, and for any choice of the Laplacian weights. All the reviewed numerical schemes for the computation of the Laplacian spectral kernels and distances are discussed in terms of robustness, approximation accuracy, and computational cost, thus supporting the reader in the selection of the most appropriate method with respect to shape representation, computational resources, and target application.

Categories and Subject Descriptors (according to ACM CCS): I.3.3 [Numerical Analysis]: Approximation/Image Generation—Special function approximations I.3.5 [Computer Graphics]: Computational Geometry and Object Modeling/Curve, surface, solid, and object representations—I.3.6 [Computer Graphics]: Methodology and Techniques—

Keywords: Laplace-Beltrami operator, Laplacian spectrum, harmonic equation, Laplacian eigenmproblem, heat equation, diffusion geometry, Laplacian spectral distance and kernels, spectral geometry processing, shape analysis, numerical analysis.

1. Introduction

In geometry processing and shape analysis, several applications have been addressed through the properties of the spectral kernels and distances, such as commute-time, biharmonic, diffusion, and wave distances. Spectral distances are easily defined through a filtering of the Laplacian eigenpairs and include random walks [FPS05, RS13], heat diffusion [BBK*10, BBOG11, CL06, GBAL09, LKC06, LSW09], biharmonic [LRF10, Rus11b], and wave kernel [BB11a, ASC11] distances. Laplacian spectral distances have been applied to shape segmentation [dGGV08] and comparison [BBOG11, GBAL09, Mem09, OMMG10, SOG09] with multi-scale and isometry-invariant signatures [DRW10, LKC06, MS05, Mem11, RBBK10, Rus07, MS09]. In fact, they are intrinsic to the input shape, invariant to isometries, multi-scale, and robust to noise and tessellation. Biharmonic [LRF10, Rus11b] distances provide a trade-off between a nearly geodesic behavior for small distances and the encoding of global surface properties for large distances, thus guaranteeing an intrinsic and multi-scale characterization of the input shape. The heat kernel [BBG94] is also central

in diffusion geometry [BN03, CL06, GK06, Sin06], dimensionality reduction with spectral embeddings [BN03, XHW10], and data classification [SK03]. As main applications, we mention the multi-scale approximation of functions [PF10] and gradients [LSW09], shape segmentation and comparison through heat kernel shape descriptors, auto-diffusion functions, and diffusion distances. The diffusion kernel and distance also play a central role in several applications, such as dimensionality reduction with spectral embeddings [BN03, XHW10]; data visualization [BN03, HAvL05, RS00, TSL00], representation [CWS03, SK03, ZGL03], and classification [NJW01, SM00, ST07].

STAR topics and contributions Our survey is intended to provide a common background on the definition and computation of Laplacian spectral kernels and distances for geometry processing and shape analysis. All the reviewed numerical schemes are discussed and compared in terms of robustness, approximation accuracy, and computational cost, thus supporting the reader in the selection of the most appropriate with respect to shape representation, computational resources, and target application. Indeed, our review is complementary to previous work, which has been focused mainly on specific applications, such as mesh filtering [Tau99], surface coding and spectral partitioning [KG00], 3D shape deformation based on differential coordinates [Sor06], spectral methods [ZvKD07] and

Laplacian eigenfunctions [Lev06] for geometry processing and diffusion shape analysis [BCA12].

Firstly, we define a unified representation of the isotropic and anisotropic discrete Laplacian on surfaces and volumes (Sect. 2); then, we introduce the associated differential equations. For the harmonic equation (Sect. 3) and the Laplacian eigenproblem (Sect. 4), we focus on the stability and accuracy of numerical solvers, also presenting their main applications. This discussion provides the background for a detailed analysis of the heat equation (Sect. 5) and allows us to identify the main limitations (e.g., computational cost, storage overhead, selection of user-defined parameters) of previous work on the approximation of the diffusion distances, which is based mainly on the evaluation of the Laplacian spectrum and on linear approximations of the exponential matrix. For the heat equation, we discuss the selection of the time scale and the main approaches for the computation of the solution to the heat equation, such as linear, polynomial, and rational approximations.

Filtering the Laplacian spectrum, we introduce the Laplacian spectral distances (Sect. 6), which generalize the commute-time, biharmonic, diffusion and wave distances, and their discretization in terms of the Laplacian spectrum. The growing interest on these distances is motivated by their capability of encoding local geometric properties (e.g., Gaussian curvature, geodesic distance) of the input shape, their intrinsic and multi-scale definition with respect to the input shape, their invariance to isometries, shape-awareness, robustness to noise and tessellation. While previous work has been focused mainly on surfaces discretized as triangle meshes, we introduce a unified representation of the spectral distances and kernels, which is independent of the selected Laplacian weights, of the surface or volume representation as polygonal mesh, point set, tetrahedral or voxel grid. From this general representation, we show that the main properties of the spectral distances are guided mainly by the filter that is applied to the Laplacian eigenpairs.

The expensive cost for the computation of the Laplacian spectrum and the sensitiveness of multiple Laplacian eigenvalues to surface discretization generally preclude an accurate evaluation of the spectral kernels and distances on large data sets. To discuss these problems, we review and compare different methods for the numerical evaluation of the spectral distances and kernels. In particular, we detail their *spectrum-free computation*, which is defined through a polynomial or rational approximation of the filter function. The resulting computational scheme only requires the solution of sparse linear systems, is not affected by the Gibbs phenomenon, is independent of the representation of the input domain, the selected Laplacian weights, and the evaluation of the Laplacian spectrum.

As main applications (Sect. 7), we detail the Laplacian smoothing and the definition of basis functions for geometry processing and shape analysis. Finally (Sect. 8), we conclude our review with a discussion of open questions and challenges. [†]
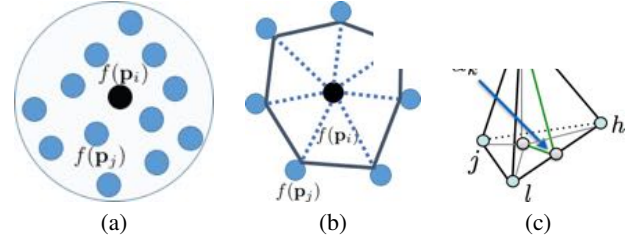

[†] Additional material is available at http://pers.ge.imati.cnr.it/patane/EG2016-STAR/EG2016-STAR.html.

Figure 1: *Neighbor and Laplacian stencil for a (a) point set, (b) triangle and (c) tetrahedral mesh.*

2. Laplace-Beltrami operator and related equations

We review the isotropic and anisotropic Laplace-Beltrami operators and introduce a unified representation of the corresponding Laplacians for surfaces and volumes. Additional results have been presented in [Sor06, Tau99, KG00, ZvKD07].

Let $\mathcal{N}$ be a smooth surface, possibly with boundary, equipped with a Riemannian metric and let us consider the *scalar product* $\langle f, g \rangle_2 := \int_{\mathcal{N}} f(\mathbf{p}) g(\mathbf{p}) d\mathbf{p}$ defined on the space $\mathcal{L}^2(\mathcal{N})$ of square integrable functions on $\mathcal{N}$ and the corresponding norm $\|\cdot\|_2$. Then, the *intrinsic smooth Laplace-Beltrami operator* $\Delta := -\text{div}(\text{grad})$ satisfies the following properties [Ros97]:

- *self-adjointness*: $\langle \Delta f, g \rangle_2 = \langle f, \Delta g \rangle_2$, $\forall f, g$;
- *positive semi-definiteness*: $\langle \Delta f, f \rangle_2 \geq 0$, $\forall f$. In particular, the Laplacian eigenvalues are positive;
- *null eigenvalue*: the smallest Laplacian eigenvalue is null and the corresponding eigenfunction ϕ, $\Delta \phi = 0$, is constant;
- *locality*: the value $\Delta f(\mathbf{p})$ does not depend on $f(\mathbf{q})$, for any couple of distinct points $\mathbf{p}$, $\mathbf{q}$;
- *linear precision*: if $\mathcal{N}$ is planar and f is linear, then $\Delta f = 0$.

The *anisotropic Laplace-Beltrami operator* [ARAC14] is defined as $\Delta_{\mathbf{D}} f = \text{div}(\mathbf{D} \nabla f)$, where $\mathbf{D}$ is a 2×2 matrix applied to vectors belonging to the tangent plane and controls the direction and strength of the deviation from the isotropic case. The tensor $\mathbf{D} := \text{diag}(\varphi_\alpha(\kappa_m), \varphi_\alpha(\kappa_M))$ takes into account the directions and the values κ_m, κ_M of low and high curvature, where the filter is $\varphi_\alpha(s) := (1 + \alpha|s|)^{-1}$, $\alpha > 0$. As $\alpha \to 0$, we get the isotropic Laplace-Beltrami operator (i.e., $\mathbf{D} := \mathbf{I}$). The alternative definition [KTT13] of the anisotropic Laplace-Beltrami operator applies a non-linear factor $\mathbf{D}(\mathbf{v})$, which modifies the magnitude of $\mathbf{D}(\mathbf{v})$ without changing its direction.

We now introduce a unified representation of the Laplacian matrix on surfaces and volumes, which is independent of the underlying discretization.

Discrete Laplacians and spectral properties Let us consider a (triangular, polygonal, volumetric) mesh $\mathcal{M} := (\mathcal{P}, T)$, which discretizes a domain $\mathcal{N}$, where $\mathcal{P} := \{\mathbf{p}_i\}_{i=1}^n$ is the set of n vertices and T is the connectivity graph (Fig. 1). On $\mathcal{M}$, a piecewise linear scalar function $f : \mathcal{M} \to \mathbb{R}$ is defined by linearly interpolating the values $\mathbf{f} := (f(\mathbf{p}_i))_{i=1}^n$ of f at the vertices using barycentric coordinates. For point sets, f is defined only at $\mathcal{P}$ and T is the k-nearest neighbor graph.

We represent the Laplace-Beltrami operator on surface and volume meshes in a unified way as $\tilde{\mathbf{L}} := \mathbf{B}^{-1}\mathbf{L}$, where $\mathbf{B}$ is a sparse, symmetric, positive definite matrix (*mass matrix*) and $\mathbf{L}$ is sparse, symmetric, and positive semi-definite (*stiffness matrix*). We also assume that the entries of $\mathbf{B}$ are positive and that the sum of each row of $\mathbf{L}$ is null. In particular, we consider the $\mathbf{B}$-scalar product $\langle \mathbf{f}, \mathbf{g} \rangle_{\mathbf{B}} = \mathbf{f}^{\top}\mathbf{B}\mathbf{g}$ and the induced norm $\|\mathbf{f}\|_{\mathbf{B}}^2 := \mathbf{f}^{\top}\mathbf{B}\mathbf{f}$. Analogously to the continuous case, the Laplacian matrix satisfies the following properties.

- *self-adjointness*: $\tilde{\mathbf{L}}$ is adjoint with respect to the $\mathbf{B}$-scalar product; i.e., $\langle \tilde{\mathbf{L}}\mathbf{f}, \mathbf{g} \rangle_{\mathbf{B}} = \langle \mathbf{f}, \tilde{\mathbf{L}}\mathbf{g} \rangle_{\mathbf{B}} = \mathbf{f}^{\top}\mathbf{L}\mathbf{g}$. If $\mathbf{B} := \mathbf{I}$, then this property reduces to the symmetry of $\mathbf{L}$;
- *positive semi-definiteness*: $\langle \tilde{\mathbf{L}}\mathbf{f}, \mathbf{f} \rangle_{\mathbf{B}} = \mathbf{f}^{\top}\mathbf{L}\mathbf{f} \geq 0$. In particular, the Laplacian eigenvalues are positive;
- *null eigenvalue*: by construction, we have that $\tilde{\mathbf{L}}\mathbf{1} = \mathbf{0}$;
- *locality*: since the weight $w(i,j)$ is not null for each edge (i,j), the value $(\tilde{\mathbf{L}}\mathbf{f})_i$ depends only on the f-values at $\mathbf{p}_i$ and its 1-star neighbor $\mathcal{N}(i) := \{j : (i,j) \text{ edge}\}$.

For a detailed discussion of these properties with respect to the selected Laplacian weights, we refer the reader to [WMKG07].

Laplacian matrix on graphs, triangle and polygonal meshes Associating a set $\{w(i,j)\}_{i,j}$ of positive weights with the edges (i,j) of T, the entries of the stiffness matrix are defined as $L(i,j) = \sum_{k \neq i} w(i,k) - w(i,j)$. The entries of the mass matrix $\mathbf{B}$ are normalization coefficients that take into account the geometry of the input domain.

On *graphs* [Chu97], the weights of the stiffness matrix are equal to 1 for each edge and zero otherwise; each diagonal entry of the mass matrix is equal to the valence of the corresponding node. On *triangle meshes*, the *stiffness matrix* $\mathbf{L}$ and the *mass matrix* $\mathbf{B}$ of the linear FEM Laplacian weights [RWP06, VL08] are defined as

$$L(i,j) := \begin{cases} w(i,j) := -\frac{\cot\alpha_{ij} + \cot\beta_{ij}}{2} & j \in N(i), \\ -\sum_{k \in N(i)} w(i,k) & i = j, \end{cases}$$

$$B(i,j) := \begin{cases} \frac{|t_r| + |t_s|}{12} & j \in N(i), \\ \frac{\sum_{k \in N(i)}|t_k|}{6} & i = j, \end{cases}$$

where $N(i)$ is the 1-star of the vertex i; α_{ij}, β_{ij} are the angles opposite to the edge (i,j) (Fig. 1b); t_r, t_s are the triangles that share the edge (i,j); and $|t|$ is the area of the triangle t. Lumping the mass matrix $\mathbf{B}$ to the diagonal matrix $\mathbf{D}$, $D(i,i) = \frac{1}{3}\sum_{t \in N(i)}|t|$, whose entries are the areas of the Voronoi regions, $\tilde{\mathbf{L}}$ reduces to the Laplacian matrix $\mathbf{D}^{-1}\mathbf{L}$ with *Voronoi-cotangent weights* [DMSB99], which extend the cotangent weights introduced in [PP93] ($\mathbf{B} := \mathbf{I}$). The *mean-value weights* [Flo03] have been derived from the mean value theorem for harmonic functions and are always positive. In [CLB*09], the weak formulation of the Laplacian eigenproblem is achieved by selecting a set of volumetric test functions, which are defined as $k \times k \times k$ B-splines (e.g., $k := 4$) and restricted to the input shape. For the *anisotropic Laplacian* [ARAC14], the entries of $\mathbf{L}$ are a variant of the cotangent weights (i.e., with respect to different angles) and the entries of the diagonal mass matrix $\mathbf{B}$ are the areas of the Voronoi regions.

While the Laplace-Beltrami operator depends only on the Reimannian metric (*intrinsic property*), its discretization is generally affected by the quality of the input triangulation [She02, HPW06]. For instance, two (simplicial) isometric surfaces with two different triangulations are associated with two different Laplacian matrices. According to [BS07], the cotangent weights are non-negative if and only if the input triangulation is Delaunay and the corresponding Laplacian matrix is more accurate than the one evaluated on the original mesh. We briefly recall [DZM07, LXH15, LXFH15] that a triangulation of a piecewise flat surface is a Delaunay triangulation if and only if all its interior edges are locally Delaunay (i.e., the sum of the angles opposite to an edge in the adjacent triangles does not exceed π). Furthermore, the minimum of the Dirichlet energy of a piecewise linear function, on all the possible triangulations of a piecewise flat surface $\mathcal{M}$, is attained at the Delaunay triangulation of $\mathcal{M}$ and the corresponding discrete Laplace-Beltrami operator is intrinsic to the input surface.

On *polygonal meshes*, the Laplacian discretization in [AW11, HKA15] generalizes the Laplacian matrix with cotangent weights to surface meshes with non-planar, non-convex faces. Finally, an approximation of the Laplace-Beltrami operator with point-wise convergence has been proposed in [BSW08].

Laplacian matrix for point sets In [BN03, BN06, BN08, BSW09], the Laplace-Beltrami operator on a point set $\mathcal{P}$ has been discretized as the Laplacian matrix

$$L(i,j) := \frac{1}{nt(4\pi t)^{3/2}} \begin{cases} \exp\left(-\frac{\|\mathbf{p}_i - \mathbf{p}_j\|_2}{4t}\right) & i \neq j, \\ -\sum_{k \neq i}\exp\left(-\frac{\|\mathbf{p}_i - \mathbf{p}_k\|_2}{4t}\right) & i = j. \end{cases}$$

To guarantee the sparsity of the Laplacian matrix, for each point $\mathbf{p}_i$ we consider only the entries $L(i,j)$ related to the points $\{\mathbf{p}_j\}_{j \in \mathcal{N}_{\mathbf{p}_i}}$ that are closest to $\mathbf{p}_i$ with respect to the Euclidean distance. In this case, we select either the k-nearest neighbor or the points that belong to a sphere centered at $\mathbf{p}_i$ and with radius σ. As described in [DS05, MN03], the choice of σ can be adapted to the local sampling density $\varepsilon := k(\pi\sigma^2)^{-1}$ and the curvature of the surface underlying $\mathcal{P}$. The computation of the k- or σ-nearest neighbor graph takes $\mathcal{O}(n\log n)$-time [AMN*98, Ben75], where n is the number of input points.

Starting from this approach, a new discretization [LPG12] has been achieved through a finer approximation of the local geometry of the surface at each point through its Voronoi cell. More precisely, as $t \to 0$ the stiffness and mass matrix are defined as

$$L(i,j) := \begin{cases} \frac{1}{4\pi t^2}\exp\left(-\frac{\|\mathbf{p}_i - \mathbf{p}_j\|_2^2}{4t}\right) & i \neq j, \\ -\sum_{j \neq i} L(i,j) & i = j, \end{cases} \qquad B(i,i) = v_i,$$

and v_i is the area of the Voronoi cell associated with the point $\mathbf{p}_i$. The Voronoi cell of $\mathbf{p}_i$ is approximated by projecting the points of a neighbor of $\mathbf{p}_i$ on the estimated tangent plane to $\mathcal{M}$ at $\mathbf{p}_i$. If $\mathbf{B} := \mathbf{I}$, then this approximation reduces to the previous one and both approaches converge to the Laplace-Beltrami operator, as $t \to 0^+$.

Laplacian matrix on volumes Representing the input domain as a tetrahedral mesh [ACSYD05, LTDZ09, TLHD03], the entries of the stiffness matrix are (Fig. 1c) $L(i,j) := w(i,j) := \frac{1}{6}\sum_{k=1}^{n} l_k \cot\alpha_k$ for each edge (i,j), $L(i,i) := -\sum_{j \in N(i)} w(i,j)$, and zero otherwise;

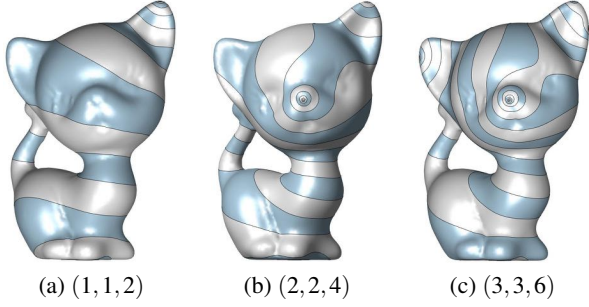| (a) $(1,1,2)$	(b) $(2,2,4)$	(c) $(3,3,6)$

Figure 2: *Level sets and critical points (m,M,s) of harmonic functions with (a) two, (b) four, and (c) six Dirichlet boundary conditions. The insertion of new initial constraints locally affects the resulting harmonic function.*

ϕ_1: $(2,2,4)$	ϕ_2: $(4,4,8)$	ϕ_3: $(5,3,8)$

Figure 3: *Level sets and number of critical points of different Laplacian eigenfunctions (linear FEM weights).*

the diagonal mass matrix **B** encodes the tetrahedral volume at each vertex.

3. Harmonic functions

The *harmonic function* $h : \mathcal{N} \to \mathbb{R}$ is the solution of the Laplace equation $\Delta h = 0$ with Dirichlet boundary conditions $h|_{\mathcal{S}} = h_0$, $\mathcal{S} \subset \mathcal{N}$. We recall that a harmonic function

- minimizes the *Dirichlet energy* $\mathcal{E}(h) := \int_{\mathcal{N}} \|\nabla h(\mathbf{p})\|_2^2 d\mathbf{p}$;
- satisfies the *locality property*; i.e., if $\mathbf{p}$ and $\mathbf{q}$ are two distinct points, then $\Delta h(\mathbf{p})$ is not affected by the value of h at $\mathbf{q}$;
- verifies $h(\mathbf{p}) = (2\pi R)^{-1} \int_{\Gamma} h(s) ds = (\pi R^2)^{-1} \int_{\mathcal{B}} h(\mathbf{q}) d\mathbf{q}$, where $\mathcal{B} \subseteq \mathcal{N}$ is a disc of center $\mathbf{p}$, radius R, and boundary Γ (*mean-value theorem*).

According to the *maximum principle* [Ros97], a harmonic function has no local extrema other than at constrained vertices. In the case that all constrained minima are assigned the same global minimum value and all constrained maxima are assigned the same global maximum value, all the constraints will be extrema in the resulting field. Harmonic and poly-harmonic (i.e., $\Delta^i h = 0$) functions have been applied to volumetric parameterization [LGW*07, LXW*10], to the definition of shape descriptors with pairs of surface points [ZTZX13] and coupled biharmonic bases [KBB*13], to shape approximation [FW12] and deformation [JMD*07, JBPS14, WPG12].

Discrete harmonic functions The harmonic equation is approximated at the vertices of $\mathcal{M}$ as the homogeneous linear system $\mathbf{Lf} = \mathbf{0}$, with initial conditions $f(\mathbf{p}_i) = a_i$, $i \in \mathcal{I} \subseteq \{1, \ldots, n\}$. According to the *Euler formula* $\chi(\mathcal{M}) = m - s + M$, the number of minima m, maxima M, and saddles s of a harmonic function depends on the Dirichlet boundary conditions, which determine the maxima and minima of the resulting harmonic function. In particular, a harmonic function with one maximum and one minimum has a minimal number of $2g$ saddles, where g is the genus of $\mathcal{M}$ (Fig. 2). Harmonic functions are efficiently computed in $\mathcal{O}(n)$ time with iterative solvers of sparse linear systems; their computation is stable for the mean-value weights while negative Voronoi cotangent weights generally induce local undulations in the resulting harmonic function. Main applications include surface
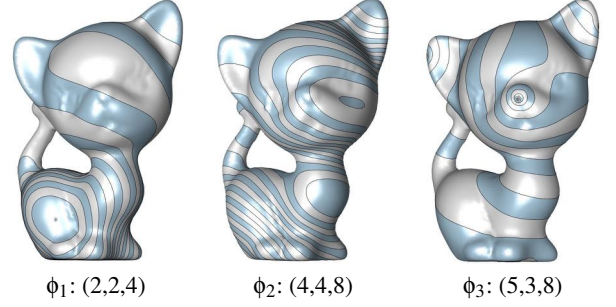
quadrangulation [DKG05, NGH04], the definition of volumetric mappings [LGQ09, LXW*10, MCK08, MC10], and biharmonic distances [OBCS*12, LRF10, Rus11b] (Sect. 6.2.2).

In the paper examples, the level sets of a given function, or kernel, or distance are associated with iso-values uniformly sampled in its range. For spectral distances, the minimum and the maximum values are depicted in blue and red, respectively. For all the other functions, colors begin with red, pass through yellow, green, cyan, blue, and magenta, and return to red. Finally, the color coding represents the same scale of values for multiple shapes.

4. Laplacian eigenfunctions

We introduce the Laplacian eigenpairs (Sect. 4.1), their discretization (Sect. 4.2), and the stability of their computation (Sect. 4.3).

4.1. Laplacian eigenmpas

Since the Laplace-Beltrami operator is self-adjoint and positive semi-definite, it has an *orthonormal eigensystem* $\mathcal{B} := \{(\lambda_n, \phi_n)\}_{n=0}^{+\infty}$, $\Delta \phi_n = \lambda_n \phi_n$, in $\mathcal{L}^2(\mathcal{N})$. In the following, we assume that the Laplacian eigenvalues are increasingly ordered; in particular $\lambda_0 = 0$. Using the orthonormality and completeness of the Laplacian eigenfunctions in $\mathcal{L}^2(\mathcal{N})$, any function can be represented as a linear combination of the eigenfunctions as $f = \sum_{n=0}^{+\infty} \langle f, \phi_n \rangle_2 \phi_n$, where $\langle f, \phi_n \rangle_2 \phi_n$ is the projection of f on ϕ_n. Furthermore, the function Δf is expressed in terms of the Laplacian spectrum as $(\Delta f)(\mathbf{p}) = \sum_{n=0}^{+\infty} \lambda_n \langle f, \phi_n \rangle_2 \phi_n(\mathbf{p})$ (*spectral decomposition theorem*). A deeper discussion of the analogies between the heat kernel, the Fourier analysis, and wavelets has been presented in [HVG11, BEKB15].

The Laplacian eigenfunctions are intrinsic to the input shape and those ones related to smaller eigenvalues correspond to smooth and slowly-varying functions. Increasing the eigenvalues, the corresponding eigenfunctions generally show rapid oscillations (Fig. 3). From the Laplacian spectrum, we can estimate geometric and topological properties of the input shape. For instance, we can compute the surface area, as the sum of the Laplacian eigenvalues; estimate the Euler characteristic of a surface with genus $g \geq 2$ through the relation [Nad88] $m_j \leq 2j - 2\chi(\mathcal{M}) + 3$, where m_j is the multiplicity of λ_j; and evaluate the total Gaussian curvature [RWP06]. If two shapes are isometric, then they have the same Laplacian

spectrum (*iso-spectral property*); however, the viceversa does not hold [GS02, ZGLG12] and we cannot recover the metric of a given surface.

4.2. Discrete Laplacian eigenpairs

To introduce the discrete Laplcian eigenpairs, the Laplacian eigenproblem is converted to its weak formulation $\langle \Delta \phi, \psi \rangle_2 = \lambda \langle \phi, \psi \rangle_2$ [All07], where ψ is a test function. The weak formulation is then discretized as $\mathbf{L}\mathbf{x} = \lambda \mathbf{B}\mathbf{x}$. Here, $\mathbf{L}$, $L(i, j) := \langle \Delta \psi_i, \psi_j \rangle_2$, is the stiffness matrix and $\mathbf{B}$, $B(i, j) := \langle \psi_i, \psi_j \rangle_2$, is the mass matrix. The *generalized Laplacian eigensystem* $\{(\lambda_i, \mathbf{x}_i)\}_{i=1}^n$ $(\lambda_1 = 0)$ satisfies the identity $\mathbf{L}\mathbf{x}_i = \lambda_i \mathbf{B}\mathbf{x}_i$ and the eigenvectors are orthonormal with respect to the $\mathbf{B}$-scalar product; i.e., $\langle \mathbf{x}_i, \mathbf{x}_j \rangle_{\mathbf{B}} = \mathbf{x}_i^\top \mathbf{B}\mathbf{x}_j = \delta_{ij}$. In particular, the spectral decomposition theorem becomes $\tilde{\mathbf{L}}\mathbf{f} = \sum_{i=1}^n \lambda_i \langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}} \mathbf{x}_i = \mathbf{X}\Gamma\mathbf{X}^\top \mathbf{B}\mathbf{f}$, where $\mathbf{X}$ is the eigenvectors' matrix and Γ is the diagonal matrix of the eigenvalues. The discrete Laplacian eigenfunctions generally have a global support (i.e., they are null only at some isolated points) and eigenfunctions with a compact support can be calculated by minimizing the corresponding ℓ_1 norm [NVT*14].

For the computation of the Laplacian eigenvectors, numerical methods generally exploit the sparsity of the Laplacian matrix and reduce the high-dimensional eigenproblem to one of lower dimension, by applying a coarsening step. The solution is efficiently calculated in the low-dimensional space and then mapped back to the initial dimension through a refinement step. Main examples include the algebraic multi-grid method [Fal06], Arnoldi iterations [LS96, Sor92], and the Nystrom method [FBCM04]. Even though the eigenvalues and eigenvectors are computed in superlinear time [VL08], this computational cost and the required $\mathcal{O}(n^2)$ storage are expensive for densely sampled domains. Indeed, modifications of the Laplacian eigenproblem are applied to locally compute specific sub-parts of the Laplacian spectrum. For instance, the shift method evaluates the spectrum $(\lambda_i - \lambda, \mathbf{x}_i)_{i=1}^n$ of $(\tilde{\mathbf{L}} - \lambda \mathbf{I})$ to calculate the eigenpairs associated with a spectral band centered around a value λ. To swap the Laplacian spectrum, the inverse method considers the spectrum $(\lambda_i^{-1}, \mathbf{x}_i)_{i=2}^n$ of the pseudo-inverse $\tilde{\mathbf{L}}^\dagger$. The power method computes the eigenpairs $(\lambda_i^k, \mathbf{x}_i)_{i=2}^n$ of the sequence of matrices $(\tilde{\mathbf{L}}^k)_{k \geq 1}$ and controls the convergence speed through the selection of k. Finally, pre-conditioners of the Laplacian matrix tailored to computer graphics' applications have been proposed in [KFS13].

Laplacian eigenfunctions on surfaces In *spectral graph theory*, the Laplacian eigenvectors have been applied to graph partitioning [Fie73, MP93, Kor03] into sub-graphs, which are handled in parallel [AKY99], to graph/mesh layout [DPS02, Kor03], to the reduction of the bandwidth of sparse matrices [BPS93]. In *machine learning*, the Laplacian spectrum have been used for clustering [SS02] (§ 14) and dimensionality reduction [BN03, XHW10] with spectral embeddings. For instance, a common way to measure the dissimilarity between two graphs is to compute the corresponding spectral decomposition in their own [LD08] or joint [Ume88, CK04] eigenspaces.

In *geometry processing*, the spectral properties of the uniform discrete Laplacian have been used to design low-pass filters [Tau95]. Successively, this formulation has been refined to include the local geometry of the input surface [DMSB99, KR05, PP93] and it has been applied to implicit mesh fairing [DMSB99, KR05, ZF03] and to fairing functionals [KCVS98, Mal89], which optimize the triangles' shape and/or the surface smoothness [NISA06]. Further applications include mesh watermarking [OTMM01, OMT02], geometry compression [KG00, SCOT03], the computation of the gradient [LSW09] and the multi-scale approximation of functions [Pat13, PS13a, PF09]. The Laplacian eigenvectors have been also used for embedding a surface of arbitrary genus into the plane [ZSGS04, ZKK02] and mapping a closed genus zero surface into a spherical domain [GGS03].

In *shape analysis*, the Laplacian spectrum has been applied to shape [LZ07, ZL05] segmentation and analysis through nodal domains [RBG*09], correspondence [JZ07, JZvK07], and comparison [MPSF11, RWP06, JZ07]. Mesh Laplacian operators are also associated with a set of differential coordinates for surface deformation [SLCO*04] and quadrangulation with Laplacian eigenfunctions [DKG05]. As detailed in Sect. 6, the Laplacian spectrum is also fundamental to define random walks [RS13], commute-time [BB11a], biharmonic [OBCS*12, Rus11b], wave kernel [BB11a, ASC11], and diffusion distances [BBK*10, BBOG11, CL06, GBAL09, LKC06, LSW09, PS13b].

Laplacian eigenfunctions on volumes Laplacian eigenfunctions on a discrete volumetric domain $\mathcal{M}$ are computed either by diagonalizing the corresponding Laplacian matrix or by extending the values of the eigenfunctions computed on the boundary of $\mathcal{M}$ to its interior with barycentric coordinates or non-linear methods (e.g., moving least-squares, radial basis functions) [PSF09, PS12]. The computational cost, which is generally high in case of volumetric meshes, is effectively reduced but associated with a lower approximation accuracy. Volumetric Laplacian eigenfunctions have been applied to shape retrieval [JZ07] and to the definition of volumetric [Rus11a] shape descriptors.

4.3. Stability of the Laplacian spectrum

Theoretical results on the sensitivity of the Laplacian spectrum against geometry changes, irregular sampling density and connectivity have been presented in [HPW06, Xu07]. Here, we briefly recall that the instability of the computation of the Laplacian eigenpairs is generally due to repeated or close eigenvalues, with respect to the numerical accuracy of the solver of the eigen-equation. While repeated eigenvalues are quite rare, numerically close or switched eigenvalues can be present in the spectrum and in spite of the regularity of the input discrete surface. The following discussion will be useful also for the definition of the conditions on the filter function that induces the spectral distances (Sect. 6.4).

To show that the computation of single eigenvalues is numerically stable, we perturb the Laplacian matrix $\tilde{\mathbf{L}}$ by $\varepsilon\mathbf{E}$, $\varepsilon \to 0$, and compute the eigenpair $(\lambda(\varepsilon), \mathbf{x}(\varepsilon))$ of the new problem $(\mathbf{B}^{-1}\mathbf{L} + \varepsilon\mathbf{E})\mathbf{x}(\varepsilon) = \lambda(\varepsilon)\mathbf{x}(\varepsilon)$, with initial conditions $\mathbf{x}(0) = \mathbf{x}$, $\lambda(0) = \lambda$. The size of the derivative of $\lambda(\varepsilon)$ indicates the variation that it undergoes when the matrix $\tilde{\mathbf{L}}$ is perturbed in the direction $(\mathbf{E}, \varepsilon)$. By differentiating the previous equation and evaluating the

result at $\varepsilon = 0$, we obtain that $\mathbf{BEx} + \mathbf{Lx}'(0) = \lambda'(0)\mathbf{Bx} + \lambda \mathbf{Bx}'(0)$. Multiplying both sides of this last relation with $\mathbf{x}^\top$, the perturbed eigenvalue $|\lambda'(0)| = |\mathbf{x}^\top \mathbf{BEx}| \leq \|\mathbf{Ex}\|_\mathbf{B}$ is bounded by the $\mathbf{B}$-norm of $\mathbf{Ex}$. Indeed, the computation of the Laplacian eigenvalue with multiplicity one is stable.

Assuming that λ_k is an eigenvalue with multiplicity m_k and rewriting the characteristic polynomial as $p_{\tilde{\mathbf{L}}}(\lambda) = (\lambda - \lambda_k)^{m_k} q(\lambda)$, where $q(\cdot)$ is a polynomial of degree $n - m_k$ and $q(\lambda_k) \neq 0$, we get that $(\lambda - \lambda_k)^{m_k} = O(\varepsilon)/q(\lambda)$; i.e., $\lambda = \lambda_k + O(\varepsilon^{\frac{1}{m_k}})$. It follows that a perturbation $\varepsilon := 10^{-m_k}$ produces a change of order 0.1 in λ_k and this amplification becomes more and more evident while increasing the multiplicity of the eigenvalue. According to [GV89] (§ 7), repeated eigenvalues are generally associated with a numerical instability in the computation of the corresponding eigenvectors; in fact, the ℓ_2-norm of the difference between the generalized eigenvectors $\mathbf{x}_i, \mathbf{x}_j$ related to the eigenvalues λ_i, λ_j is bounded as

$$\|\mathbf{x}_i - \mathbf{x}_j\|_2 \leq \varepsilon \sum_{j \neq i} \left| \frac{\mathbf{x}_i^\top \mathbf{Ex}_j}{\lambda_i - \lambda_j} \right| + O(\varepsilon^2).$$

Indeed, the computation of the eigenvectors related to multiple or close Laplacian eigenvalues might be unstable. Finally, the Laplacian eigenvalues might be locally switched (i.e., we are not able to numerically distinguish two consecutive eigenvalues) and this situation happens independently of the quality of the discretized surface in terms of point density, angles, and connectivity.

5. Heat and wave equations

We introduce the heat (Sect. 5.1), wave and mean curvature flow (Sect. 5.2) equations; then, we discuss their discretization (Sect. 5.3), the selection of the time scale (Sect. 5.4), and the computation of their solution (Sect. 5.5).

5.1. Heat equation

The *scale-based representation* $H : \mathcal{N} \times \mathbb{R}^+ \to \mathbb{R}$ of the function $h : \mathcal{N} \to \mathbb{R}$ is the solution to the *heat diffusion equation* $(\partial_t + \Delta)H(\mathbf{p},t) = 0$, $H(\cdot,0) = h$. The function $H(\mathbf{p},t)$ represents the heat distribution at the point $\mathbf{p}$ and at time t, where h is the initial distribution. The solution to the heat equation is written as

$$H(\mathbf{p},t) = \langle K_t(\mathbf{p},\cdot),h \rangle_2 = \sum_{n=0}^{+\infty} \exp(-\lambda_n t)\langle h, \phi_n \rangle_2 \phi_n(\mathbf{p}), \quad (1)$$

where $K_t(\mathbf{p},\mathbf{q}) = \sum_{n=0}^{+\infty} \exp(-\lambda_n t)\phi_n(\mathbf{p})\phi_n(\mathbf{q})$ is the spectral representation of the *heat diffusion kernel*. The heat diffusion and the Laplace-Beltrami operators have the same eigenfunctions $\{\phi_n\}_{n=0}^{+\infty}$ and $(\exp(-\lambda_n t))_{n=0}^{+\infty}$ are the eigenvalues of the heat operator. The heat kernel is invariant to isometries and verifies the *semi-group* $\langle K_{t_1}, K_{t_2} \rangle_2 = K_{t_1+t_2}$ and *inversion* $K_t^{-1} = K_{-t}$ *properties*. The spectral representation (1) shows the smoothing effect on the initial condition h; as the scale increases, the component of h along the eigenfunctions associated with the larger Laplacian eigenvalue becomes null. We also notice that the normalized function $\mathcal{A}_\mathcal{N}^{-1} H(\cdot,t)$ with respect to the surface area $\mathcal{A}_\mathcal{N}$ minimizes the weighted least-squares error $\int_\mathcal{N} K_t(\mathbf{p},\mathbf{q})|h(\mathbf{q}) - g(\mathbf{p})|^2 d\mathbf{q}$ on $\mathcal{L}^2(\mathcal{N})$, for a given h.

Heat equation on surfaces On surfaces, the heat kernel satisfies the following properties [SOG09, Gri06]:

- for an isometry $\Phi : \mathcal{N} \to \mathcal{Q}$ between two manifolds $\mathcal{N}, \mathcal{Q}$,

$$K_t^\mathcal{N}(\mathbf{p},\mathbf{q}) = \mathcal{K}_t^\mathcal{Q}(\Phi(\mathbf{p}),\Phi(\mathbf{q})), \qquad \forall \mathbf{p},\mathbf{q} \in \mathcal{N}, \forall t \in \mathbb{R}^+; \quad (2)$$

- if Φ is surjective and Eq. (2) holds, then Φ is an isometry;
- if D is a compact set of $\mathcal{N}$, then $\lim_{t \to 0} K_t^D(\mathbf{p},\mathbf{q}) = K_t^\mathcal{N}(\mathbf{p},\mathbf{q})$;
- if $D_1 \subseteq D_2 \subseteq \mathcal{N}$, then $K_t^{D_1}(\mathbf{p},\mathbf{q}) \leq K_t^{D_2}(\mathbf{p},\mathbf{q})$;
- on smooth and polygonal surfaces, the heat kernel fully determines the Riemannian metric [ZGLG12].

For small values of t [SOG09, Var67], the auto-diffusivity function

$$K_t(\mathbf{p},\mathbf{p}) \approx \begin{cases} (4\pi t)^{-1}(1 + 1/3t\kappa(\mathbf{p})) + \mathcal{O}(t^2), \\ (4\pi t)^{3/2}(1 + 1/6s(\mathbf{p})), \end{cases} \quad t \to 0,$$

encodes the Gaussian $\kappa(\mathbf{p})$ and total $s(\mathbf{p})$ curvature at $\mathbf{p}$. For large t, $K_t(\mathbf{p},\mathbf{q})$ is dominated by the Fiedler vector ϕ_1 [Fie73], which encodes the global structure of the input shape. According to [SOG09, dGGV08], the surface $\mathcal{N}$ at $\mathbf{p}$ can be characterized in terms of the average squared diffusion distance at $\mathbf{p}$ (*eccentricity*), which is defined as

$$\text{ecc}_t(\mathbf{p}) = \mathcal{A}_\mathcal{N}^{-1} \int_\mathcal{N} d_t(\mathbf{p},\mathbf{q})d\mathbf{q} = K_t(\mathbf{p},\mathbf{p}) + E_\mathcal{N}(t) - 2\mathcal{A}_\mathcal{N}^{-1}, \quad t \to 0,$$

where $E_\mathcal{N}(t) := \sum_{n=0}^{+\infty} \exp(-\lambda_n t)$ is the sum of the eigenvalues of the heat kernel. Since the area and trace are independent of the evaluation point, the functions ecc_t and $K_t(\mathbf{p},\cdot)$ have the same level sets and extrema on $\mathcal{N}$. In particular, for small scales the extrema of the eccentricity are localized at the curvature extrema.

Heat equation on volumes The analytical representation of the volumetric heat kernel $K_t(\mathbf{p},\mathbf{q}) := (4\pi t)^{-3/2} \exp(-\|\mathbf{p}-\mathbf{q}\|_2^2/4t)$ allows us to solve the heat equation as $F(\cdot,t) = K_t \star h$ and without computing the Laplacian spectrum (Sect. 5.5.4).

5.2. Wave equation and mean curvature flow

The heat equation is strictly related to the *Schroedinger (wave) equation* $(i\Delta + \partial_t)H(\cdot,t) = 0$, with initial condition $H(\cdot,0) = h$, which represents the physical model of a quantum particle with initial energy h. The spectral representation of the solution is $H(\cdot,t) = \sum_{n=0}^{+\infty} \exp(i\lambda_n t)\langle f, \phi_n \rangle_2 \phi_n$; i.e., a complex wave function with oscillatory behavior. This periodic effect is due to the real and complex parts of the filter $\exp(i\lambda_n t) = \cos(\lambda_n t) + i\sin(\lambda_n t)$. The norm of the solution is the probability $P_t(\mathbf{p})$ to find a point $\mathbf{p}$ after a time t; in fact, the following identity holds

$$P_t(\mathbf{p}) = \lim_{T \to +\infty} \int_0^T |H(\mathbf{p},t)|^2 dt$$
$$= \sum_{n=0}^{+\infty} |\langle f, \phi_n \rangle_2|^2 |\phi_n(\mathbf{p})|^2 = \|H(\mathbf{p},t)\|_2^2.$$

Finally, the heat equation is related to the *mean curvature flow* [CPS13, KSB12] $(\partial_t + \Delta_t)\Phi_t = 0$, where $\Phi_t : \mathcal{M} \to \mathbb{R}^3$ is a family of immersions and Δ_t is the Laplace-Beltrami operator associated with the metric induced by the immersion at time t.
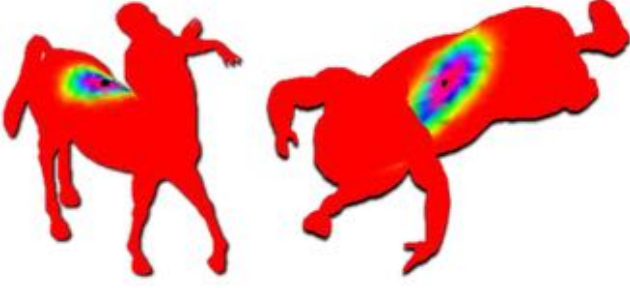
Figure 4: *Anisotropic heat kernel centered at a (black) seed point on a coarse triangle mesh.*

Table 1: *Main properties of the discrete heat kernel: sparsity, positive definiteness, and symmetry. The full ● and empty ○ circle means that the corresponding property is or is not satisfied, respectively.*

Heat Ker.	Matrix $\mathbf{K}_t$	Sp.	Pos. Def.	Sym.	Cov.	Inv.
Std	$\mathbf{X}\mathbf{D}_t\mathbf{X}^\top$	○	●	●	○	○
Vor.-cot	$\mathbf{X}\mathbf{D}_t\mathbf{X}^\top\mathbf{D}$	○	●	○	●	○
wFEM	$\mathbf{X}\mathbf{D}_t\mathbf{X}^\top\mathbf{B}$	○	●	○	●	○

5.3. Discrete heat equation and kernel

We briefly introduce the weak formulation [All07] of the heat equation; similar results apply to the equations previously introduced. Chosen a set $\mathcal{B} := \{\psi_i\}_{i=1}^n$ of linearly independent functions on $\mathcal{N}$, we approximate the solution $\tilde{H}(\mathbf{p},t) := \sum_{i=1}^n a_i(t)\psi_i(\mathbf{p})$ to the weak heat equation as $\langle \partial_t \tilde{H}(\cdot,t), \psi_i \rangle_2 + \langle \Delta \tilde{H}(\cdot,t), \psi_i \rangle_2 = 0$, $i = 1,\ldots,n$. Introducing the matrices $\mathbf{L} := (\langle \Delta \psi_i, \psi_j \rangle_2)_{i,j=1}^n$ and $\mathbf{B} := (\langle \psi_i, \psi_j \rangle_2)_{i,j=1}^n$, the discrete heat equation becomes $(\mathbf{B}\partial_t + \mathbf{L})\mathbf{a}(t) = \mathbf{0}$, $\mathbf{a}(t) := (a_i(t))_{i=1}^n$. An analogous relation can be derived for the boundary condition $H(\mathbf{p},0) = h(\mathbf{p})$. Since $\mathbf{B}$ is the Gram matrix associated with $\mathcal{B}$, it is invertible and the previous system of equations is $(\partial_t + \mathbf{B}^{-1}\mathbf{L})\mathbf{a}(t) = \mathbf{0}$. Then, the solution to the discrete heat equation is expressed as a linear combination of the Laplacian eigensystem as $\mathbf{F}(t) = \sum_{i=1}^n \exp(-\lambda_i t)\langle \mathbf{f}, \mathbf{x}_i \rangle_\mathbf{B} \mathbf{x}_i$.

Properties The solution to the discrete heat equation is $\mathbf{F}(t) = \mathbf{K}_t \mathbf{f}$ (Fig. 4), where $\mathbf{K}_t := \mathbf{X}\mathbf{D}_t\mathbf{X}^\top\mathbf{B}$, $\mathbf{D}_t := \text{diag}(\exp(-\lambda_i t))_{i=1}^n$, is the *heat kernel matrix* (Table 1). Lumping the linear FEM mass matrix $\mathbf{B}$, the heat kernel becomes equal to the *Voronoi-cotangent heat kernel* $\mathbf{K}_t^\star := \mathbf{X}\mathbf{D}_t\mathbf{X}^\top\mathbf{D}$, $\mathbf{L}\mathbf{X} = \mathbf{X}\Gamma$. Choosing $\mathbf{B} := \mathbf{I}$, we get the *heat kernel* $\tilde{\mathbf{K}}_t := \mathbf{X}\mathbf{D}_t\mathbf{X}^\top$ with cotangent weights. Analogously to the results in Sect. 5.1, the heat kernel matrix satisfies the following relations: $\mathbf{K}_{t_1} \times \mathbf{K}_{t_2} = \mathbf{K}_{t_2} \times \mathbf{K}_{t_1} = \mathbf{K}_{t_1+t_2}$ (*commutative* and *semi-group properties*), $\mathbf{K}_t^{-1} = \mathbf{K}_{-t}$ (*inversion property*). If $\mathbf{B}$ is the linear FEM mass matrix or the diagonal matrix of the Voronoi areas, then the heat kernel matrix $\mathbf{K}_t$ is intrinsically *scale-covariant*; i.e., rescaling the points of $\mathcal{M}$ by a factor α, $\alpha > 0$, and indicating the new surface as $\alpha\mathcal{M}$, we get that only the time component of the kernel is rescaled. In fact, the rescaling changes the matrix $\mathbf{B}$ and the eigensystem $\{(\lambda_i,\mathbf{x}_i)\}_{i=1}^n$ of $\mathcal{M}$ into $\alpha^2 \mathbf{B}$ and $\{(\alpha^{-2}\lambda_i, \alpha^{-1}\mathbf{x}_i)\}_{i=1}^n$, respectively. Indeed, $\mathbf{K}_t(\alpha\mathcal{M}) = \mathbf{K}_{\alpha^{-2}t}(\mathcal{M})$ without an *a-posteriori* normalization. The scale-covariance of $\mathbf{K}_t$ is guaranteed by the mass matrix, which
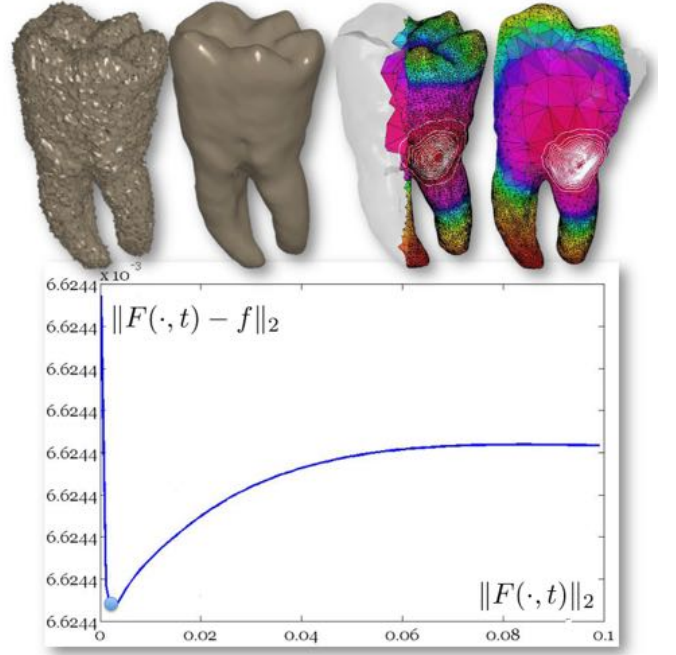

Figure 5: *(c) Selection of the optimal scale ($t_{opt} = 0.0032$) and corresponding volumetric diffusion smoothing (a), Padé-Chebyshev approximation of degree $r = 7$) on the noisy volumetric model of a teeth (b).*

changes according to the surface rescaling and compensates the variation of the corresponding Laplacian spectrum. The kernel becomes *scale-invariant* (i.e., $\mathbf{K}_t(\alpha\mathcal{M}) = \mathbf{K}_t(\mathcal{M})$) by normalizing each eigenvalue by λ_n, which is efficiently computed using the inverse method [GV89, VL08]. Alternatively, the scale-invariance and covariance of the heat kernel is achieved in the Fourier domain [BK10]. In [BBB*10, BBC*10], the matching performances of heat kernel descriptors have been tested against shape transformation, sampling, and noise.

5.4. Selection of the time scale

For shape analysis, the real line is uniformly sampled in order to consider both small and large scales. For geometry processing, the *optimal time value* is defined as the value of t that provides the best compromise between a small residual error $\|F(\cdot,t) - f\|_2^2 = \sum_{n=0}^{+\infty} |1 - \exp(-2\lambda_n t)|^2 |\langle f,\phi_n \rangle_2|^2$ and a low energy $\|F(\cdot,t)\|_2^2 = \sum_{n=0}^{+\infty} \exp(-2\lambda_n t)|\langle f,\phi_n \rangle_2|^2$. If t tends to zero, then the residual becomes null and the energy converges to $\|f\|_2$. If t becomes large, then the residual increases until it converges to $|\langle f,\phi_0 \rangle_2|$ and the solution norm decreases until it converges to $(\|f\|_2^2 - |\langle f,\phi_0 \rangle_2|^2)^{1/2}$. According to these properties, the plot (*L-curve*) of the energy (*y*-axis) versus the residual (*x*-axis) is L-shaped [HO93] and its minimum provides the optimal time value (Fig. 5). For the computation of the optimal time value, we mention the corner detection based on cubic B-splines approxima-

Table 2: *Numerical computation of the solution to the heat equation; $\tau(n)$ is the cost for the solution of a sparse linear system.*

Method	Numerical scheme	Scales	Comput. cost	References
Linear approximation				
Trunc. spec. approx.	$\mathbf{F}_k(t) = \sum_{i=1}^k \exp(-\lambda_i t)\langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}} \mathbf{x}_i$	Any	$\mathcal{O}(kn)$	[GV89, VBCG10]
Euler backw. approx.	$(t\tilde{\mathbf{L}}+\mathbf{I})\mathbf{F}_{k+1}(t) = \mathbf{F}_k(t)$	Small	$\mathcal{O}(\tau(n))$	[CDR00, DMSB99, ZH08]
I order Taylor approx.	$\mathbf{BF}(t) = (\mathbf{B}-t\mathbf{L})\mathbf{f}$	Small	$\mathcal{O}(\tau(n))$	[CDR00, DMSB99]
Krylov/Schur approx.	Projection on	Any	$\mathcal{O}(m\tau(n)), \mathbf{B} \neq \mathbf{I}$	[GV89, Saa92, ZH08]
	$\{\mathbf{g}_i := (\mathbf{B}^{-1}\mathbf{L})^i \mathbf{f}\}_{i=1}^m$		$\mathcal{O}(n), \mathbf{B} = \mathbf{I}$	
Polynomial approximation				
Power approx.	$\mathbf{F}(t) = \sum_{i=0}^m \mathbf{g}_i/i!$	Any	$\mathcal{O}(m\tau(n)), \mathbf{B} \neq \mathbf{I}$	
	$\mathbf{g}_i := \tilde{\mathbf{L}}^i \mathbf{f}$		$\mathcal{O}(n), \mathbf{B} = \mathbf{I}$	[GV89]
Rational approximation				
Padé-Cheb. approx.	$\mathbf{F}(t) = \alpha_0 \mathbf{f} + \sum_{i=1}^r \mathbf{g}_i$	Any	$\mathcal{O}(r\tau(n))$	[CRV84, Sid98, Saa92]
	$(t\mathbf{L} + \theta_i \mathbf{B})\mathbf{g}_i = -\alpha_i \mathbf{Bf}$			[Pat13, Pat14, Pat16]
Contour integral approx.	$\mathbf{F}(t) = \sum_{i=1}^r \alpha_i \mathbf{g}_i$	Any	$\mathcal{O}(r\tau(n))$	[Pus11]
	$(\alpha_i)_{i=1}^r$ quadr. coeff.			

tion [HO93], the evaluation of the curvature of the graph of the L-curve or its adaptive pruning [HO93].

5.5. Computation of the discrete heat kernel

For the computation of the solution to the discrete heat equation and kernel, we consider linear (Sect. 5.5.1), polynomial (Sect. 5.5.2), and rational (Sect. 5.5.3) approximations. On volumes (Sect. 5.5.4), we discuss the solution to the heat equation based on the analytic representation of the heat kernel. With the exception of the truncated spectral method, all the previous approximations are independent of the evaluation of the Laplacian spectrum and reduce to a set of sparse linear systems (Table 2). The polynomial and rational approximations generally provide the best compromise between approximation accuracy and computational cost.

5.5.1. Linear approximation

For the solution to the heat equation, we review the truncated spectral approximation, the Euler backward method, the first order Taylor approximation, the Krylov and Schur methods.

Truncated spectral approximation and power method The computational bottleneck for the evaluation of the whole Laplacian spectrum imposes on us to consider only a small subset of the Laplacian spectrum. Since the decay of the *filter factor* $\exp(-\lambda_i t)$ increases with λ_i, in the spectral representation of the solution to the heat equation we consider only the contribution related to the first k eigenpairs; i.e., $\mathbf{F}_k(t) = \sum_{i=1}^k \exp(-\lambda_i t)\langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}} \mathbf{x}_i$. The truncated approximation is accurate only if the exponential filter decays fast (e.g., large values of time) and the effect of the selected eigenpairs on the approximation accuracy cannot be estimated without computing the whole spectrum. The *multi-resolution prolongation operators* [VBCG10] prolongate the values of the truncated spectral approximation, computed on a low-resolution representation of the input shape, to the initial resolution through a hierarchy of simplified meshes. In this case, the number of eigenpairs are heuristically adapted to the surface resolution and its global/local features.

Euler backward method In [CDR00, DMSB99], the solution to the heat equation is computed through the Euler backward method $(t\tilde{\mathbf{L}}+\mathbf{I})\mathbf{F}_{k+1}(t) = \mathbf{F}_k(t)$, $\mathbf{F}_0 = \mathbf{f}$. The resulting functions are over-smoothed and converge to a constant function, as $k \to +\infty$.

First order Taylor approximation Since the derivative of $\mathbf{K}_t$ at $t = 0$ equals the Laplacian matrix (i.e., $(\mathbf{I} - \mathbf{K}_t)/t \to \mathbf{B}^{-1}\mathbf{L}$, $t \to 0$), the heat kernel $\mathbf{K}_t$ is approximated by $(\mathbf{I} - t\mathbf{B}^{-1}\mathbf{L})$ and $\mathbf{F}(t) = \mathbf{K}_t\mathbf{f}$ solves the sparse linear system $\mathbf{B}(\mathbf{K}_t\mathbf{f}) = (\mathbf{B} - t\mathbf{L})\mathbf{f}$. This last relation gives an approximation of $\mathbf{F}(t)$ that is independent of the Laplacian spectrum and is valid only for small values of t. For an arbitrary value of t, the "power" method applies the identity $(\mathbf{K}_{t/m})^m = \mathbf{K}_t$, where m is chosen in such a way that t/m is sufficiently small to guarantee that the approximation $\mathbf{K}_{t/m} \approx (\mathbf{I} - t/m\tilde{\mathbf{L}})$ is accurate. However, the selection of m and its effect on the approximation accuracy cannot be estimated a-priori.

Krylov and Schur approximations The Krylov subspace projection [GV89, Saa92] computes an approximation of $\exp(-t\mathbf{A})\mathbf{f}$ in the space generated by the vectors $\mathbf{f}, \mathbf{Af}, \dots, \mathbf{A}^{m-1}\mathbf{f}$, thus processing a $m \times m$ matrix instead of a $n \times n$ matrix, where m is much lower than n (e.g., $m \approx 20$). This approximation [ZH08] becomes computationally expensive when the dimension of the Krylov space increases, still remaining much lower than n (e.g., $n \approx 5K$). In both cases, the vector $\tilde{\mathbf{L}}^i \mathbf{f} = (\mathbf{B}^{-1}\mathbf{L})^i \mathbf{f}$ must be computed without inverting the mass matrix; to this end, we notice that the vector $\mathbf{g}_i := (\mathbf{B}^{-1}\mathbf{L})^i \mathbf{f}$ satisfies the linear system $\mathbf{Bg}_i = \mathbf{Lg}_{i-1}$, $\mathbf{Bg}_1 = \mathbf{Lf}$. Since the coefficient matrix $\mathbf{B}$ is sparse, symmetric, and positive definite, the vectors $(\mathbf{g}_i)_{i=1}^m$ are evaluated in linear time by applying iterative solvers (e.g., conjugate gradient) or pre-factorizing $\mathbf{B}$.

5.5.2. Polynomial approximations

The exponential of a matrix $\mathbf{A}$ is defined as the exponential *power series* $\exp(\mathbf{A}) = \sum_{n=0}^{+\infty} \mathbf{A}^n/n!$, which converges for any square matrix $\mathbf{A}$. Even though the input matrix $\mathbf{A}$ is sparse, its exponential $\exp(-t\mathbf{A})$ is always full ($t \neq 0$) and can be computed or stored

---

Algorithm 1 Padé-Chebyshev approximation of the solution to the heat equation.

Require: A function $f : \mathcal{P} \to \mathbb{R}$, $\mathbf{f} := (f(\mathbf{p}_i))_{i=1}^{n}$.

Ensure: The approximate solution $\mathbf{F}(t) = \mathbf{K}_t\mathbf{f}$ of $\mathbf{f}$ to the heat equation.

1: Select the value of t (e.g., optimal value, Sect. 5.4).
2: **for** $i = 1, \ldots, r - 1$ **do**
3:     Compute $\mathbf{g}_i$: $(t\mathbf{L} + \theta_i\mathbf{B})\mathbf{g}_i = -\alpha_i\mathbf{B}\mathbf{f}$.
4: **end for**
5: Approximate $\mathbf{K}_t\mathbf{f}$ as $\alpha_0\mathbf{f} + \sum_{i=1}^{r} \mathbf{g}_i$.
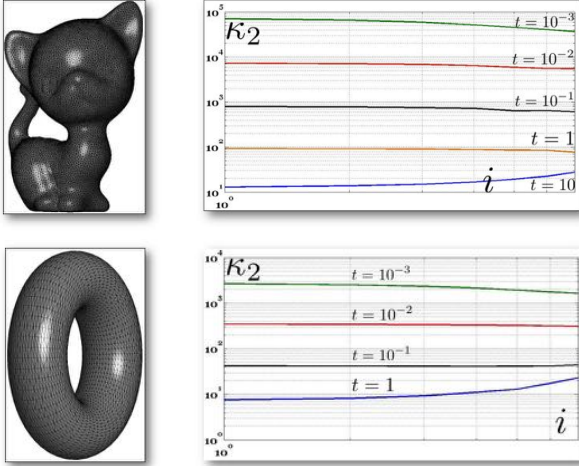
Figure 6: *Conditioning number κ_2 (y-axis) of the matrices $\{(t\mathbf{L}+\theta_i\mathbf{B})\}_{i=1}^{7}$, for several values the time parameter t; the indices of the coefficients $\{\theta_i\}_{i=1}^{7}$ are reported on the x-axis.*

only if $\mathbf{A}$ has a few hundred rows and columns only. In particular, for computer graphics applications we can consider 3D shapes only with a small number of samples (i.e., few hundreds) or evaluate the heat kernel on a set of seed points that are representative of the geometry and features of the input shape.

5.5.3. Rational approximation

The exponential of an arbitrary matrix $\mathbf{A}$ is equal to the complex contour integral $\exp(t\mathbf{A}) = (2\pi i)^{-1} \int_\Gamma \exp(z)(z\mathbf{I} - t\mathbf{A})^{-1}\mathrm{d}z$, where Γ is a closed contour winding once around the spectrum of $t\mathbf{A}$ [GV89] (§ 11), [Rud87] (§ 10). From this identity, we introduce two accurate and computationally efficient approximations of the exponential of the Laplacian matrix.

Padé-Chebyshev approximation The rational approximation of the exponential function of order (k, k) and with simple poles is $r_{kk}(z) := p_k(z)/q_k(z) = \alpha_0 + \sum_{i=1}^{k} \alpha_i(z - \theta_i)^{-1}$, where p_k, q_k are polynomials of order k, $\alpha_0 = \lim_{z \to +\infty} r_{kk}(z)$, θ_i is a pole, and α_i is the residual at θ_i. Applying this last relation to $t\mathbf{A}$, we get $\exp(t\mathbf{A}) = \alpha_0\mathbf{I} + \sum_{i=1}^{k} \alpha_i(t\mathbf{A} - \theta_i\mathbf{I})^{-1}$. Among the rational approximations of the exponential function, we focus on its best approximation $r_{kk}(\cdot)$ of order k with respect to the ℓ_∞ norm; i.e., the unique $r_{kk}(z) = p_k(z)/q_k(z)$ that minimizes the error $\|\pi(z) - \exp(-z)\|_\infty$
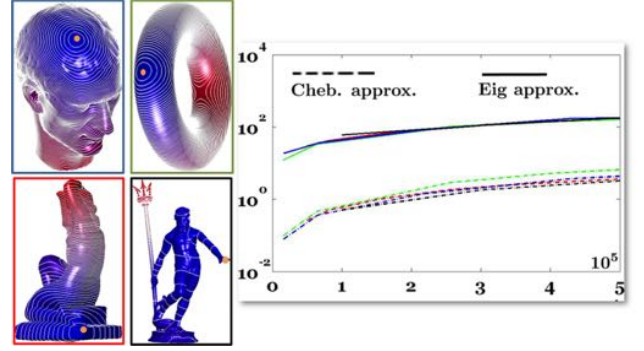

Figure 7: *Cost (in seconds, y-axis, log-scale) for the evaluation of the diffusion distances on 3D shapes with n samples (x-axis), approximated with $k = 500$ eigenpairs and the Padé-Chebyshev approximation. Colors from the source (orange) point vary from blue (null distance) to red (maximum distance).*

in the space $\Pi_{kk} \ni \pi$ of rational polynomials of order k. Here, the main difficulty is the evaluation of the coefficients and poles of the rational approximation of the exponential function for a given k, which is generally affected by the ill-conditioned computation of the polynomial roots. These coefficients and poles have been computed with a different accuracy and for different orders of the rational polynomial [CRV84, CMV69, MVL03, Sid98, Saa92]. These approximations are also included in standard numerical libraries for signal processing. Finally, we recall that in spectral graph theory [OSV12], the Padé-Chebyshev and the Lanczos methods have been applied to the approximation of $\exp(-\mathbf{A})\mathbf{f}$, where $\mathbf{A}$ is a symmetric and positive semi-definite matrix.

The idea behind the *spectrum-free computation* [Pat13, Pat14] is to apply the (r, r)-degree Padé-Chebyshev rational approximation to the exponential representation $\mathbf{F}(t) = \exp(-t\tilde{\mathbf{L}})\mathbf{f}$ of the solution to the heat equation $(\partial_t + \tilde{\mathbf{L}})\mathbf{F}(t) = \mathbf{0}$, $\mathbf{F}(0) = \mathbf{f}$ (Algorithm 1). In this case, the solution $\mathbf{F}(t) = \alpha_0\mathbf{f} + \sum_{i=1}^{r} \mathbf{g}_i$ is the sum of the solutions of r sparse linear systems $(t\mathbf{L} + \theta_i\mathbf{B})\mathbf{g}_i = -\alpha_i\mathbf{B}\mathbf{f}$, $i = 1, \ldots, r$. The resulting approximation belongs to the linear space generated by $\mathbf{f}$ and $\{\mathbf{g}_i\}_{i=1}^{r}$, which are calculated as a minimum norm residual solution [GV89], depend on the input domain, the initial condition, and the selected time value. In comparison, the Laplacian eigenfunctions only encode the domain geometry and it is difficult to select the number of eigenpairs necessary to achieve a given approximation of $\mathbf{F}(t)$ with respect to t and $\mathbf{f}$.

This approximation is independent of the computation of the Laplacian spectrum, user-defined parameters, and multi-resolutive prolongation operators [VBCG10], which heuristically adapt the number of eigenpairs to the surface resolution. The sparse and well-conditioned matrices of the previous linear systems have the same structure and sparsity of the connectivity matrix of the input domain, properly encode the local and global features in the heat kernel, and can be computed for any representation of the input domain and of the Laplacian weights. Finally, the accuracy of the Padé-Chebychev approximation is lower than 10^{-r} (e.g., $r = 5, 7$).

The value of t influences the conditioning number of the ma-
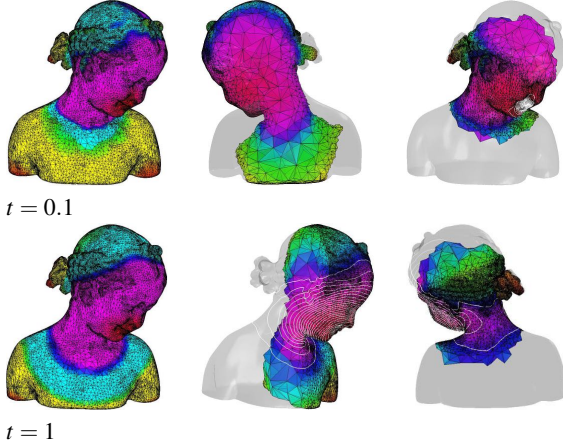
$t = 0.1$

$t = 1$

Figure 8: *Volumetric heat kernel ($r = 7$). Level-sets correspond to iso-values uniformly sampled in the range of the solution restricted to the volume boundary.*

trices $(t\mathbf{L} + \theta_i\mathbf{B})$, $i = 1, \ldots, r$. Experiments (Fig. 6, [Pat14]) have shown that the linear systems associated with the Padé-Chebyshev approximation are generally well-conditioned; in any case, preconditioners and regularization techniques [GV89] can be applied to attenuate numerical instabilities. Finally, timings on surfaces and volumes (Fig. 7) are reduced from 20 up to 1200 times with respect to the approximation based on a fixed number of Laplacian eigenpairs. Laplacian eigenvectors have been computed with the Arnoldi iteration method [LS96, Sor92].

Rational approximation from contour integrals Since the exponential factor rapidly decays to zero as $Re(z) \to +\infty$, in [Pus11] the complex contour integral has been efficiently computed with quadrature rules. In this case, $\alpha_0 = 0$, the poles $\theta_i := \phi(x_i)$ are evaluated at the quadrature points $\{x_i\}_i$, $\alpha_i := -(2\pi i)^{-1}h\exp(\phi(x_i))\phi'(x_i)$ are the weights of the quadrature rules, and h is the interval length in the quadrature scheme. The resulting approximation accuracy is guided by the degree of the quadrature rule; low degrees (e.g., $k = 2$, $k = 4$) generally provide a satisfactory approximation accuracy.

5.5.4. Special case: heat equation on volumes

On a volume, the function $\mathbf{F}(t) = \sum_{i=1}^n \alpha_i V_i K_t(\mathbf{p}_i, \cdot)f(\mathbf{p}_i)$ is approximated as a linear combination of the basis functions $\{K_t(\mathbf{p}_i, \cdot)\}_{i=1}^n$. Here, $\mathbf{V} = \text{diag}(V_i)_{i=1}^n$ is the diagonal matrix of the volumes V_i at $\mathbf{p}_i$, $\mathbf{K}_t$ is the Gram matrix for the Gaussian kernel, and the unknowns $\alpha = (\alpha_i)_{i=1}^n$ are determined by imposing the condition $F(\mathbf{p}_i, 0) = f(\mathbf{p}_i)$, $i = 1, \ldots, n$. To overcome the time-consuming solution of the $n \times n$ linear system $\mathbf{VK}_t\alpha = \mathbf{f}$, the number of conditions is reduced or the coefficient matrix is sparsified according to the exponential decay of its entries. Alternatively, the volumetric heat equation is solved by discretizing the Laplace-Beltrami operator with finite elements [All07, RWSN09], or with finite differences on a 6-neighborhood stencil [LBB11, LBB12, RBBK10], or with a geometry-driven approximation of the gradient field [LTDZ09, TLHD03].

While a discretization of the heat kernel on a voxel grid is accurate enough for the evaluation of diffusion descriptors [LBB11, RBBK10], which are quantized and clustered in bags-of-features, the computation of the solution to the volumetric heat equation generally requires a more accurate discretization of the input domain. The prolongation of the Laplacian [Rus11a, Rus11b], harmonic [LXW*10, MCK08], and diffusion functions from the volume boundary to its interior, through barycentric coordinates or non-linear approximation, achieves a low accuracy of the solution in a neighbor of the boundary. The multi-resolution simplification of the input volume is also time-consuming, and the selection of the volume resolution with respect to the expected approximation accuracy are generally guided by heuristics. Indeed, these methods do not intend to approximate the heat kernel quantitatively, but provide alternative approaches that qualitatively behave like the heat kernel on volumes. To improve the accuracy, we consider the volumetric Laplacian matrix of the input domain and compute the Padé-Chebyshev approximation of the induced heat kernel (Fig. 8).

6. Laplacian spectral distances

Distances can be defined directly on the input domain $\mathcal{M}$ (e.g., geodesic distances) or in the space of functions on $\mathcal{M}$ (e.g., random walks, biharmonic distances, diffusion and wave distances). For geometry processing and shape analysis, the distance

- must satisfy the following properties: *nullity* ($d(\mathbf{p}, \mathbf{q}) = 0$ if and only if $\mathbf{p} \equiv \mathbf{q}$); *symmetry* ($d(\mathbf{p}, \mathbf{q}) = d(\mathbf{q}, \mathbf{p})$); *triangular inequality* ($d(\mathbf{p}, \mathbf{q}) \leq d(\mathbf{p}, \mathbf{r}) + d(\mathbf{r}, \mathbf{q})$);
- should be *multi-scale* and *geometry-aware*, through the encoding of local/global features and geometric properties;
- should be *isometry-invariant* through a proper filtering of the Laplacian spectrum, *robust* to noise and domain discretization.

We introduce the spectral distances (Sect. 6.1), by filtering the Laplacian spectrum and as a generalization of the commute-time, biharmonic, diffusion and wave distances (Sect. 6.2). Then, we discuss their spectral discretization (Sect. 6.3), computation (Sect. 6.4), and comparison (Sect. 6.5).

6.1. Laplacian spectral kernels and distances

Starting from recent work on the geodesic and heat diffusion distances [CWW13, Pat13], we address the definition of spectral distances on a manifold $\mathcal{N}$ by filtering its Laplacian spectrum [BB11b, Pat14]. Given a strictly positive *filter function* $\varphi : \mathbb{R}^+ \to \mathbb{R}$, let us consider the power series $\varphi(s) = \sum_{n=0}^{+\infty} \alpha_n s^n$. Noting that $\Delta^i f = \sum_{n=0}^{+\infty} \lambda_n^i \langle f, \phi_n \rangle_2 \phi_n$, we define the *spectral operator* as

$$\Phi(f) = \sum_{n=0}^{+\infty} \alpha_n \Delta^n f = \sum_{n=0}^{+\infty} \varphi(\lambda_n)\langle f, \phi_n \rangle_2 \phi_n. \quad (3)$$

According to [Pat16], if the function $\tilde{\varphi}(s) := s^{1/2}\varphi(s)$ is integrable on $\mathbb{R}^+$ then the spectral operator is well-defined, linear, continuous, and $\Phi(f) = K \star f$, where $K(\mathbf{p}, \mathbf{q}) = \sum_{n=0}^{+\infty} \varphi(\lambda_n)\phi_n(\mathbf{p})\phi_n(\mathbf{q})$ is the *spectral kernel*. Through the spectral operator, in $\mathcal{L}_2(\mathcal{N})$ we introduce the *spectral scalar product* and *distance* as

$$\begin{cases} \langle f, g \rangle := \langle \Phi(f), \Phi(g) \rangle_2 = \sum_{n=0}^{+\infty} \varphi^2(\lambda_n)\langle f, \phi_n \rangle_2 \langle g, \phi_n \rangle_2 \\ d^2(f, g) = \|f - g\|^2 = \sum_{n=0}^{+\infty} \varphi^2(\lambda_n)|\langle f - g, \phi_n \rangle|^2. \end{cases} \quad (4)$$
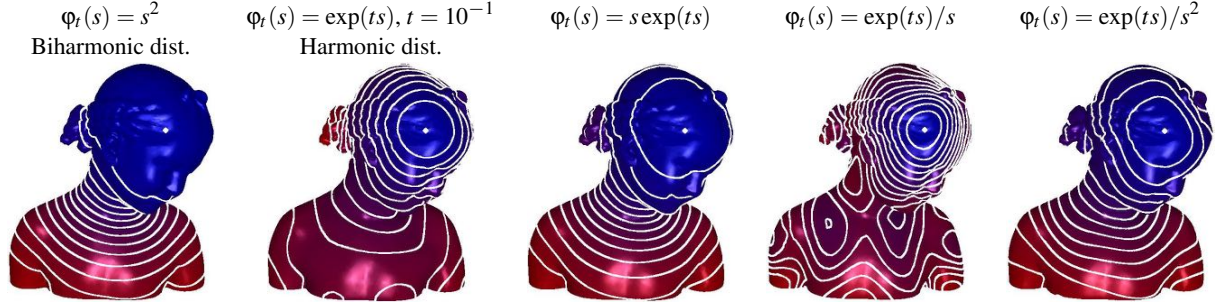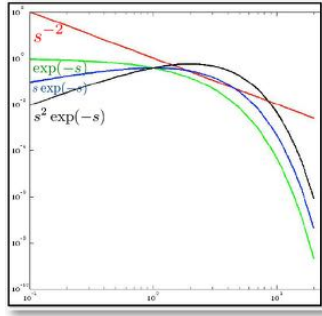
Figure 9: *Level-sets of the spectral distances from a source point (white dot) induced by the filter φ and evaluated with the Padé-Chebyshev approximation ($r = 5$).*

φ^2	Spect. dist.	$d^2(\mathbf{p},\mathbf{q}) = \sum_{n=0}^{+\infty} \varphi^2(\lambda_n)|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$	
		Associated equation	Kernel
s^{-1} s^{-2} s^{-k}	Comm.-time Biharm. Poly-harm.	$(\partial_t + \Delta^r)F(\cdot,t) = 0$ *Poly-harm. eq.*	$\sum_{n=0}^{+\infty} t^k e^{-\lambda_n t}|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$
e^{-st}	Heat diff.	$(\partial_t + \Delta)F(\cdot,t) = 0$ *Heat diffusion eq.*	$\sum_{n=0}^{+\infty} e^{-\lambda_n t}|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$
e^{-ist}	Wave ker.	$(\partial_t + i\Delta)F(\cdot,t) = 0$ *Schroedinger eq.*	$\sum_{n=0}^{+\infty} e^{-i\lambda_n t}|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$

Figure 10: *Spectral distances and kernels induced by the filter function φ (log-scale on the t- and y-axis) applied to the Laplacian eigenvalues.*

Indicating with $\delta_\mathbf{p}$ the function that takes value 1 at $\mathbf{p}$ and 0 otherwise, the *spectral distance* between $\mathbf{p}, \mathbf{q}$ is (Fig. 9)

$$d^2(\mathbf{p},\mathbf{q}) := \|\delta_\mathbf{p} - \delta_\mathbf{q}\|^2 = \sum_{n=0}^{+\infty} \varphi^2(\lambda_n)|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$$

$$= \|K(\mathbf{p},\cdot) - K(\mathbf{q},\cdot)\|_2^2 = K(\mathbf{p},\mathbf{p}) - 2K(\mathbf{p},\mathbf{q}) + K(\mathbf{q},\mathbf{q}),$$

where the first row provides the spectral representation and the second row expresses the spectral distances in terms of the corresponding kernel.

Properties Analogously to the diffusion kernel, the spectral kernel satisfies the following properties: *non-negativity* ($K(\mathbf{p},\mathbf{p}) \geq 0$), *symmetry* ($K(\mathbf{p},\mathbf{q}) = K(\mathbf{q},\mathbf{p})$), and *positive semi-definiteness*:

$$0 \leq \langle \Phi(f), f \rangle_2 = \int_{\mathcal{N} \times \mathcal{N}} K(\mathbf{p},\mathbf{q})f(\mathbf{p})f(\mathbf{q})\mathrm{d}\mathbf{p}\mathrm{d}\mathbf{q}$$

$$= \sum_{n=0}^{+\infty} \varphi(\lambda_n)|\langle f, \phi_n \rangle_2|^2.$$

We also mention the *square integrability* $\|K\|_2^2 = \sum_{n=0}^{+\infty}|\varphi(\lambda_n)|^2$, which is equivalent to the Parseval's equality and the *conservation*: $\int_{\mathcal{N}} K(\mathbf{p},\mathbf{q})\mathrm{d}\mathbf{p} = 1$, which is a consequence of the Perron-Frobenious theorem.

Through the selected filter function and the Laplacian spectrum, we define the *spectral embedding* $\mathcal{E}: \mathcal{M} \to \ell_2$, which maps each point $\mathbf{p}$ to the sequence $\mathcal{E}(\mathbf{p}) := (\varphi(\lambda_n)\phi_n(\mathbf{p}))_{n=0}^{+\infty}$. The equality $d(\mathbf{p},\mathbf{q}) = \|\mathcal{E}(\mathbf{p}) - \mathcal{E}(\mathbf{q})\|_2$ shows that the spectral distances can be interpreted as Euclidean distances in the embedding space. Finally, the *spectral shape descriptor* $SD(\mathbf{p}) := \sum_{n=0}^{+\infty} |\varphi^2(\lambda_n)|\phi_n^2(\mathbf{p})$ and *signature* $SE(\mathbf{p}) := (\varphi^{-1/2}(\lambda_n)\phi_n(\mathbf{p}))_{n=0}^{+\infty}$ generalize the diffusion descriptor and signature [DLL*10, SOG09] (Sect. 6.2.1).

Selection of the filter function The filter function is learned from a training data set [ABBK11] or chosen in such a way that the corresponding spectral distances satisfy the properties introduced at the beginning of Sect. 6. For instance (Fig. 10), selecting $\varphi_t(s) := \exp(-st), \exp(-ist)$ or $\varphi(s) := s^{-k/2}, s^{-1/2}$, we get the *heat diffusion*, *wave*, or *poly-harmonic*, *commute-time distances*, respectively. Mexican hat wavelets [HQ12] are generated by the filter $\varphi(s) := s^{1/2}\exp(-s^2)$ and in [BB11a, ASC11] the filter function $\varphi(s) := \exp(is), s \in [0, 2\pi]$, defines the wave kernel signature. The spectral distances associated with this periodic filter identify local shape features by separating the contribution of different frequencies and of the corresponding eigenfunctions.

Similarly to random walks [RS13], we introduce multi-scale kernels by integrating the moment of order k of the differential operator $\Delta^\alpha \exp(-t\Delta^\alpha)$. In this case, the filter function is $\varphi(s) = t^k s^\alpha \exp(-ts^\alpha)$, where k scales the rate of diffusion and α controls the decay of the Laplacian eigenvalues to zero. The selection of the parameters α, k makes the multi-scale kernels more robust to geometric and topological noise; the integral over time also avoids the selection of the heat diffusion rate. The filter functions $\varphi_t(s) := [\cos^{-1/2}(\sqrt{st}), s^{-1/4}\sin^{1/2}(\sqrt{st})]$ and $\varphi(s,t) = \exp(s^r t)$ are associated with the diffusion equations $(\partial_t^2 + \Delta)F(\cdot,t) = 0$ and
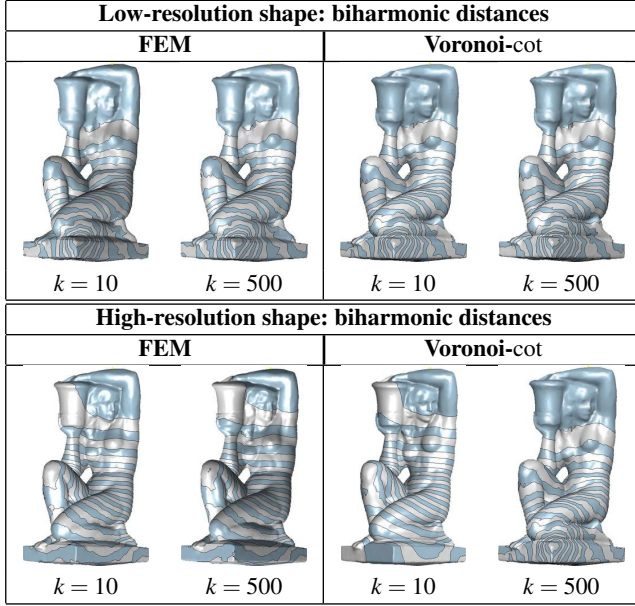
Low-resolution shape: biharmonic distances			
FEM		Voronoi-cot	
$k = 10$	$k = 500$	$k = 10$	$k = 500$
High-resolution shape: biharmonic distances			
FEM		Voronoi-cot	
$k = 10$	$k = 500$	$k = 10$	$k = 500$

Figure 11: *Biharmonic distance on a surface at different resolutions, with different Laplacian weights and k eigenpairs.*

$(\partial_t + \Delta^r)F(\cdot, t) = 0$, respectively. Finally, the filter function can be learned from a set of retrieval examples [ABBK11, BMM*15].

6.2. Main examples of spectral distances

As special cases, we consider the diffusion (Sect. 6.2.1), commute-time and biharmonic distances (Sect. 6.2.2), and the approximation of the geodesic and transportation distances (Sect. 6.2.3).

6.2.1. Diffusion distances

The heat kernel induces the *diffusion distances*, whose spectral representation is $d_t^2(\mathbf{p}, \mathbf{q}) = \sum_{n=0}^{+\infty} \exp(-\lambda_n t)|\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2$. Recalling that the heat kernel is self-adjoint with respect to the scalar product induced by the mass matrix $\mathbf{B}$, we define the *diffusive scalar product* $\langle \mathbf{f}, \mathbf{g} \rangle_t := \langle \mathbf{K}_t \mathbf{f}, \mathbf{g} \rangle_{\mathbf{B}}$ and express the *discrete diffusion distances* as $d_t(\mathbf{p}_i, \mathbf{p}_j) = \|\delta_{\mathbf{p}_i} - \delta_{\mathbf{p}_j}\|_t = \|\mathbf{K}_t(\mathbf{e}_i - \mathbf{e}_j)\|_{\mathbf{B}}$, where $\|\mathbf{f}\|_t^2 = \mathbf{f}^\top \mathbf{X} \mathbf{D}_t \mathbf{X}^\top \mathbf{B} \mathbf{f} = \sum_{i=1}^n \exp(-\lambda_i t)|\langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}}|^2$ is the *diffusion norm*.

Through the heat kernel, a shape is associated with a diffusion metric that measures the rate of connectivity among its points with paths of length t and characterizes the local/global geometric behavior with small/large values of t. This property has been used to define a multi-scale and isometry-invariant signatures [BBK*10, BK10, BBOG11, CL06, DRW10, GBAL09, LKC06, MS05, Mem09, Mem11, OMMG10, RBBK10, Rus07, MS09, SOG09] and to rewrite the shape similarity problem as the comparison of two metric spaces. Main examples include the *heat kernel signature* $HKS(\mathbf{p}) := \sum_{n=0}^{+\infty} \exp(-\lambda_n t)|\phi_n(\mathbf{p})|^2$ and *descriptor* $HKD(\mathbf{p}) := (\lambda_n^{-1/2}\phi_n(\mathbf{p}))_{n=0}^{+\infty}$, and the *wave kernel signature* $WKS(\mathbf{p}) := \sum_{n=0}^{+\infty} \exp(-i\lambda_n t)|\phi_n(\mathbf{p})|^2$. Furthermore, the heat diffusion distance and kernel have been successfully applied to
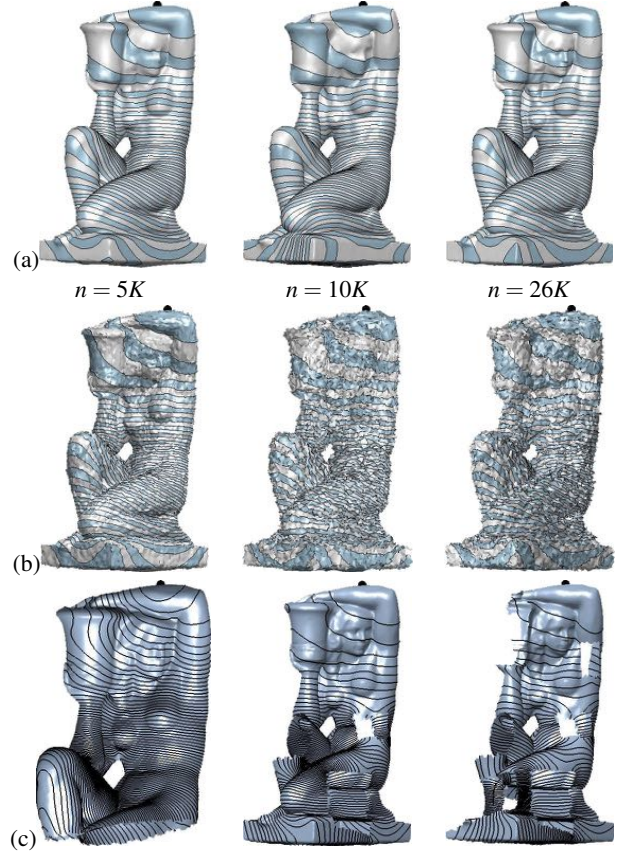

Figure 12: *Stability of the biharmonic distance from a source (black) point with respect to (a) sampling, (b) noise, (c) holes.*

shape segmentation [dGGV08]; the computation of the gradient of discrete functions [LSW09]; and the multi-scale approximation of functions [PF10]. The diffusion distance and kernel also play a central role in several applications, such as dimensionality reduction with spectral embeddings [BN03, XHW10]; data visualization [BN03, HAvL05, RS00, TSL00], representation [CWS03, SK03, ZGL03], and classification [NJW01, SM00, ST07].

6.2.2. Commute-time and biharmonic distances

Integrating the diffusion distances with respect to t, we get the *commute-time distance*

$$d^2(\mathbf{p}, \mathbf{q}) = \frac{1}{2} \int_0^{+\infty} d_t^2(\mathbf{p}, \mathbf{q}) \mathrm{d}t = \sum_{n=0}^{+\infty} \lambda_n^{-1} |\phi_n(\mathbf{p}) - \phi_n(\mathbf{q})|^2,$$

which is induced by the filter $\varphi(s) := s^{1/2}$ and is scale-invariant. While the diffusion distance estimates the connection of two points with respect to any random walk of length t, the commute-time distance measures this connection with respect to arbitrary random walks. The *biharmonic distances* [OBCS*12, LRF10, Rus11b] are induced by $\varphi(s) := s$ and provide a trade-off between a nearly-geodesic behavior for small distances and global shape-awareness for large distances, thus guaranteeing an intrinsic multi-scale characterization of the input shape. In Fig. 11, the approximation of
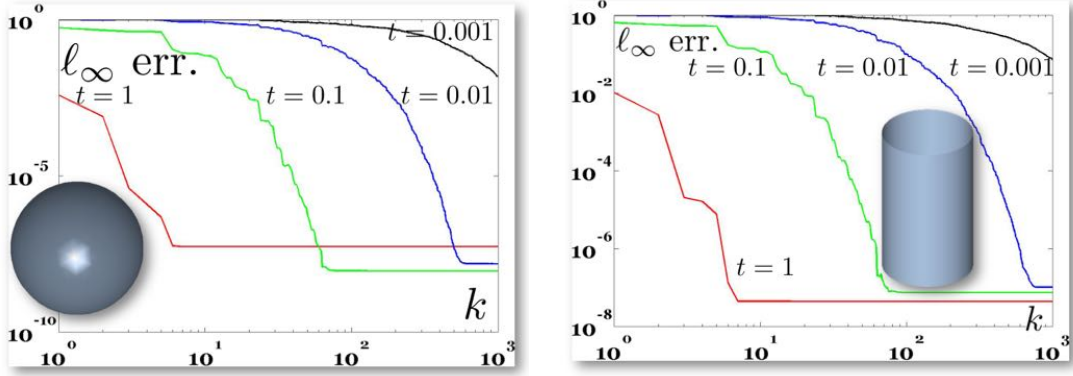
Figure 13: ℓ_∞ *error (y-axis) for the diffusion distance approximated with k (x-axis) Laplacian eigenpairs. For the Padé-Chebyshev method* $(r = 5)$ *and all the scales, the* ℓ_∞ *error with respect to the ground-truth is lower than* 8.9×10^{-6}.
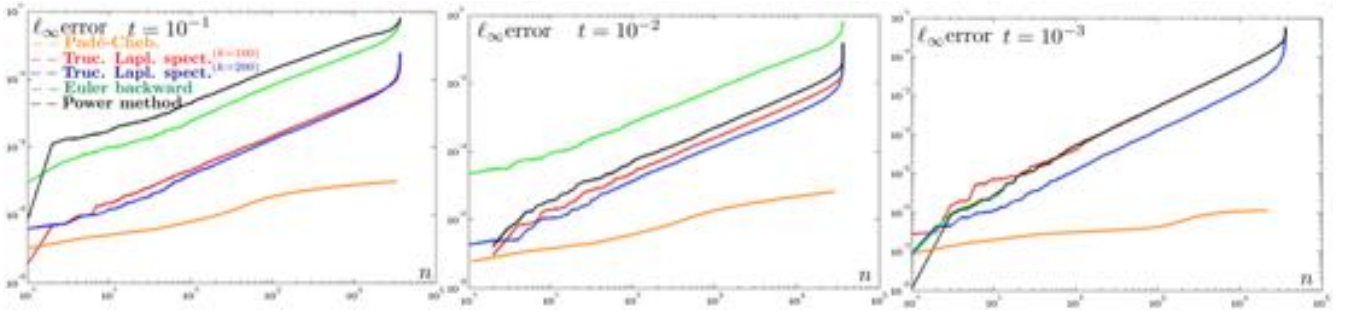

Figure 14: ℓ_∞ *error (y-axis) between the ground-truth diffusion distances on the cylinder, with a different sampling (x-axis). For different scales, the accuracy of the Padé-Chebyshev method* $(r = 5, \text{orange})$ *remains almost unchanged and higher than the truncated approximation with* 100 *and* 200 *eigenpairs (red, blue), the Euler backward (green) and power (black) methods.*

the biharmonic kernel and distance with a subset of the Laplacian spectrum presents local artifacts, which are represented by isolated level sets and are reduced by increasing the number of eigenpairs without disappearing. In Fig. 12, the smooth and uniform distribution of the level sets of the biharmonic distance around the anchor point (black dot) confirms the stability of the spectrum-free approximation with respect to surface sampling, noise, and missing parts.

6.2.3. Approximating geodesics and transportation distances with the heat kernel

In [CWW13], the relation $d_G(\mathbf{p}, \mathbf{q}) = -\lim_{t \to 0} (4t \log K_t(\mathbf{p}, \mathbf{q}))$ has been applied to compute the geodesic distance $d_G(\mathbf{p}, \mathbf{q})$ from the heat kernel values $K_t(\mathbf{p}, \mathbf{q})$ as the scale tends to zero. More precisely, the geodesic distance d_G on $\mathcal{N}$ is approximated by computing the solution $F(\cdot, t)$ to the heat equation $(\partial_t + \Delta) F(\cdot, t) = 0$ on $\mathcal{N}$, as $t \to 0^+$, normalizing the corresponding gradient $X = \nabla F(\cdot, t) / \|\nabla F(\cdot, t)\|_2$, and solving the equation $\Delta d_G = \operatorname{div}(X)$. This approximation is computationally efficient, capable of identifying different types of features by selecting different diffusion models, and robust to noise. The main difficulty is the tuning of the time scale with respect to the shape features; in fact, the selection of a large scale is generally associated with an over-smoothing of the geodesic values and local details.

In [SdGP*15], the *optimal transportation distances* have been approximated using the iterative Sinkhorn's method [Sin64] and the entropic regularization, thus reducing their computation to the solution of two sparse matrix equations that involve the heat kernel matrix. Instead of approximating the heat kernel with an implicit Euler integration [DMSB99], we can apply the Padé-Chebyshev approximation (Sect. 5.5.3) in order to improve the approximation accuracy at small scales and without modifying the overall approach.

6.3. Discrete spectral distances

Inserting the generalized eigensystem $\mathbf{LX} = \mathbf{BX\Lambda}$, with orthonormal eigenvectors $\mathbf{X}^\top \mathbf{BX} = \mathbf{I}$, in Eq. (3), its discretization is $\mathbf{K} = \mathbf{X}\varphi(\Lambda)\mathbf{X}^\top \mathbf{B}$, $\varphi(\Lambda) := \operatorname{diag}(\varphi(\lambda_i))_{i=1}^n$, and the corresponding *discrete spectral distances* are

$$d^2(\mathbf{p}_i, \mathbf{p}_j) = \|\mathbf{K}(\mathbf{e}_i - \mathbf{e}_j)\|_{\mathbf{B}}^2 = \sum_{l=1}^n \varphi^2(\lambda_i)|\langle \mathbf{x}_l, \mathbf{e}_i - \mathbf{e}_j\rangle_{\mathbf{B}}|^2. \quad (5)$$

The spectral representation of the kernel provides its link between the Laplacian matrix; i.e., $\tilde{\mathbf{L}}$ and $\mathbf{K}$ have the same eigenvectors and $(\varphi(\lambda_i))_{i=1}^n$ are the (filtered) Laplacian eigenvalues of $\mathbf{K}$.

In previous work, the spectral distances have been discretized as $d(\mathbf{p}_i, \mathbf{p}_j) = \|\mathbf{K}^\star(\mathbf{e}_i - \mathbf{e}_j)\|_2$ and with respect to the Euclidean scalar

(a) $t = 10^{-1}$ (b) $t = 10^{-2}$ (c) $t = 10^{-3}$

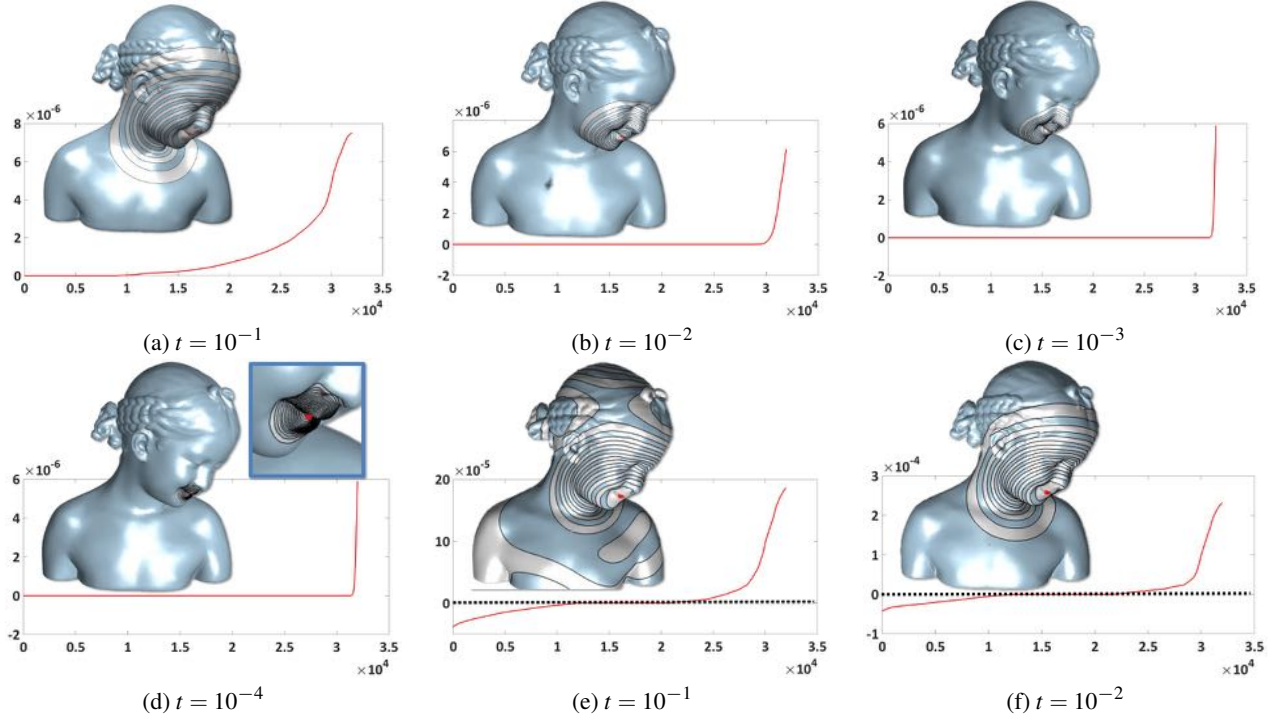(d) $t = 10^{-4}$ (e) $t = 10^{-1}$ (f) $t = 10^{-2}$

Figure 15: *(a-d) Robustness of the Padé-Chebyshev approximation of the diffusion distances and (e,f) sensitiveness of truncated spectral approximation to the Gibbs phenomenon. At all scales (a-d), the distance values (red curve) computed with the Padé-Chebyshev approximation are positive; at large scales (e,f), the truncated spectral approximation is affected by the Gibbs phenomenon, as represented by the part of the plot below the zero line (black curve).*

product, where $\mathbf{K}^{\star} := \mathbf{X}\varphi(\Lambda)\mathbf{X}^{\top}$ is the corresponding kernel. This last discretization does not take into account the intrinsic $\mathbf{B}$-scalar product, thus disregarding the geometry of the input data and the underlying generalized eigenproblem. Considering the linear FEM mass matrix $\mathbf{B}$ and noting that $B(i, j) = \langle 1_{\mathbf{p}_i}, 1_{\mathbf{p}_j} \rangle_2$, where $1_{\mathbf{p}}$ is the function that takes value 1 at $\mathbf{p}$ and 0 otherwise, the $\mathbf{B}$-scalar product is the counterpart of the $\mathcal{L}_2(\mathcal{N})$ scalar product on the space of discrete functions on $\mathcal{M}$. The orthogonality of the Laplacian eigenvectors with respect to the $\mathbf{B}$-scalar product is crucial to encode the geometry of the surface underlying $\mathcal{M}$ in the spectral distances and makes its evaluation robust to surface sampling.

6.4. Computation of the spectral distances

Recalling that the computation of the Laplacian eigenpairs is numerically unstable in case of repeated eigenvalues (Sect. 4.3), the filter function should be chosen in such a way that the filtered Laplacian matrix does not have additional (if any) repeated eigenvalues. This condition is generally satisfied by choosing an injective filter. The selection of periodic filters, the expensive cost of the computation of the Laplacian spectrum, and the sensitiveness of multiple Laplacian eigenvalues to surface discretization are the main motivations for the definition of alternative approaches for the evaluation of the spectral distances and kernels. Among them, we discuss the truncated (Sect. 6.4.1) and spectrum-free (Sect. 6.4.2) approximations.

Algorithm 2 Computation of the spectral distances.

Require: A surface or volume $\mathcal{M}$, a filter function $\varphi : \mathbb{R} \to \mathbb{R}$.
Ensure: The spectral distance $d(\mathbf{p}_i, \mathbf{p}_j)$ in Eq. (5), $\mathbf{p}_i, \mathbf{p}_j \in \mathcal{M}$.
1: Compute $(\mathbf{L}, \mathbf{B})$, which define the Laplacian $\tilde{\mathbf{L}} := \mathbf{B}^{-1}\mathbf{L}$.
2: Define the vector $\mathbf{f} = \mathbf{e}_i - \mathbf{e}_j$.
3: *CASE I - Arbitrary filter: polynomial approximation*
4: Compute the polynomial approx. $p_r(s) = \sum_{i=0}^{r} \alpha_i s^i$ of φ.
5: Compute $\mathbf{g}_1$: $\mathbf{B}\mathbf{g}_1 = \mathbf{L}\mathbf{f}$.
6: **for** $i = 1, \dots, r - 1$ **do**
7:      Compute $\mathbf{g}_{i+1}$: $\mathbf{B}\mathbf{g}_{i+1} = \mathbf{L}\mathbf{g}_i$
8: **end for**
9: Compute $\mathbf{u} = \mathbf{K}\mathbf{f} \approx p_r(\tilde{\mathbf{L}}) = \alpha_0 \mathbf{f} + \sum_{i=1}^{r} \alpha_i \mathbf{g}_i$ (c.f., Eq. (6)).
10: Compute the distance $d(\mathbf{p}_i, \mathbf{p}_j) = \|\mathbf{u}\|_{\mathbf{B}}$.
11: *CASE II - Arbitrary filter: Padé-Chebyshev approximation*
12: Compute the P.C. approx. $p_r(s) = \sum_{i=1}^{r} \alpha_i (1 + \beta_i s)^{-1}$ of φ.
13: **for** $i = 1, \dots, r$ **do**
14:      Compute $\mathbf{g}_i$: $(\mathbf{B} + \beta_i \mathbf{L})\mathbf{g}_i = \mathbf{B}\mathbf{f}$ (c.f., Eq. (7))
15: **end for**
16: Compute $\mathbf{u} = \mathbf{K}\mathbf{f} \approx p_r(\tilde{\mathbf{L}})\mathbf{f} = \sum_{i=1}^{r} \alpha_i \mathbf{g}_i$.
17: Compute the distance $d(\mathbf{p}_i, \mathbf{p}_j) = \|\mathbf{u}\|_{\mathbf{B}}$.

6.4.1. Truncated approximation

The computational limits for the evaluation of the whole Laplacian spectrum and the decay of the coefficients in Eq. (4) are the main reasons behind the approximation of the solution to the spectral
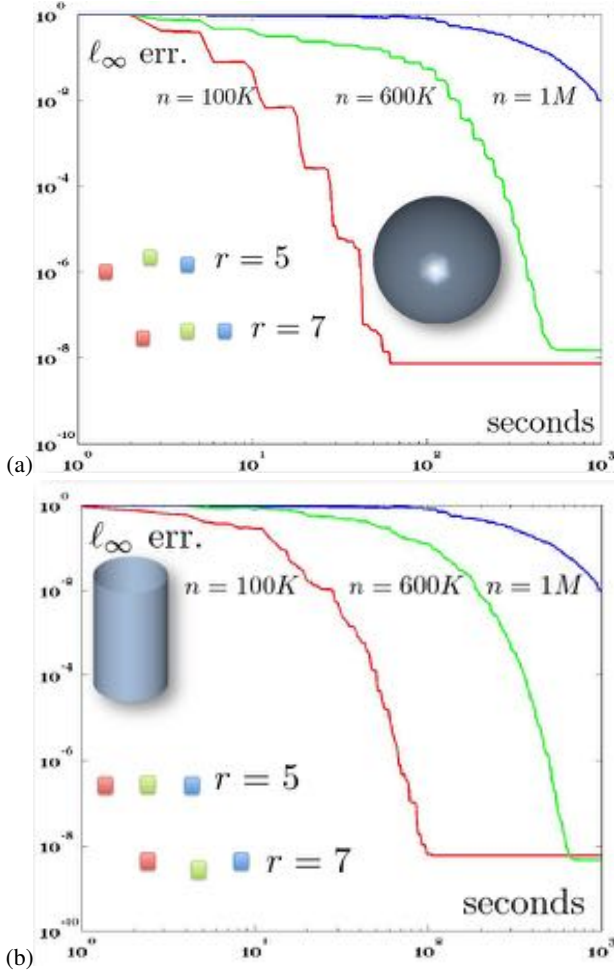
(a)

(b)

Figure 16: *Trade-off between accuracy (y-axis) and time (x-axis) for the Padé-Chebyshev (r = 5, 7) and truncated approximations (k = 50 eigenpairs) on the (a) sphere and (b) cylinder.*

Figure 17: *Timings (in seconds) for the evaluation of the heat kernel on a domain with n points, approximated with $k = 100, 500$ eigenpairs* (Eigs) *and the Padé-Chebyshev approximation (r = 7).*

distances as a truncated sum; i.e.,

$$\begin{cases} \Phi_k \mathbf{f} = \sum_{i=1}^{k} \varphi(\lambda_i) \langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}} \mathbf{x}_i \\ d^2(\mathbf{p}_i, \mathbf{p}_j) = \sum_{l=1}^{k} \varphi^2(\lambda_l) |\mathbf{x}_l^\top \mathbf{B} \mathbf{e}_i - \mathbf{x}_l^\top \mathbf{B} \mathbf{e}_j|^2, \end{cases}$$

where k is the number of selected eigenpairs. Even though the first k Laplacian eigenpairs are computed in super-linear time [VL08], the evaluation of the whole Laplacian spectrum is unfeasible for storage and computational cost, which are quadratic in the number of surface samples. Furthermore, the selection of filters that are periodic or do not decrease to zero motivates the need of defining a spectrum-free computation of the corresponding kernels and distances, which cannot be accurately approximated with the contribution of only a subpart of the Laplacian spectrum. The number of selected eigenpairs is heuristically adapted to the decay of the filter function and the approximation accuracy cannot be estimated without computing the whole spectrum.
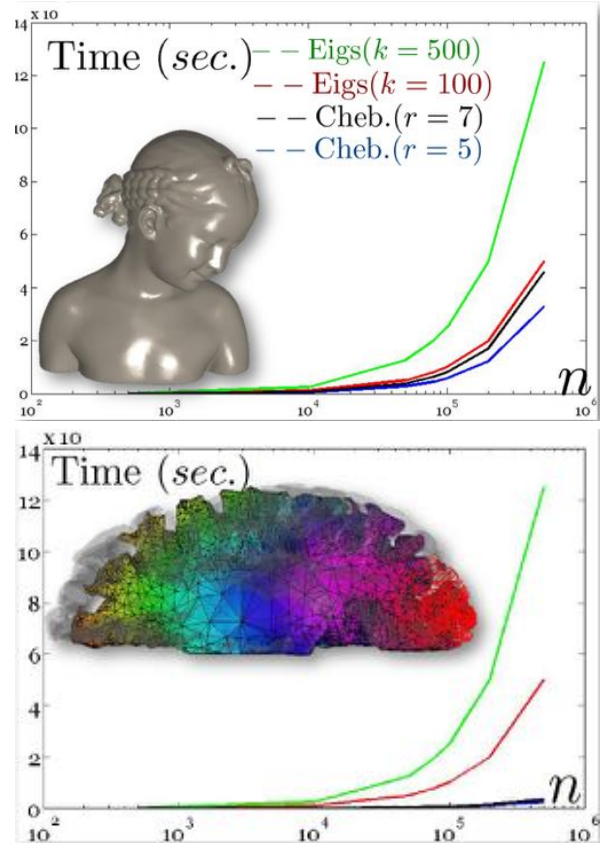
6.4.2. Spectrum-free approximation

We now introduce the spectrum-free evaluation of the spectral distances, which is based on a polynomial or rational approximation of the filter.

Arbitrary filter: polynomial approximation For an arbitrary filter φ, the matrix $\varphi(\mathbf{A})$ is approximated by selecting a new function g such that the matrix $\tilde{\mathbf{A}} = g(\mathbf{A})$ approximates $\mathbf{A}$ and can be easily calculated. One of the main approaches for the approximation of a matrix function is through the truncated Taylor approximation [GV89]. More precisely, given the power series representation $\varphi(s) = \sum_{n=0}^{+\infty} \alpha_n s^n$ defined on an open disk containing the spectrum of $\mathbf{A}$, we have that $\varphi(\mathbf{A}) = \sum_{n=0}^{+\infty} \alpha_n \mathbf{A}^n$. In this case, it is enough to consider the contribution of the first k terms in the sum and to compute the powers $(\mathbf{A}^i)_{i=1}^{k}$, through a binary powering [VL79].

Let $[0, \lambda]$ be an interval that contains the spectrum of $\tilde{\mathbf{L}}$, where λ is the maximum eigenvalue, which is computed by the Arnoldi method [GV89], or is set equal to the upper bound [LS96, Sor92] $\lambda_n \leq \min\{\max_i\{\sum_j \tilde{L}(i, j)\}, \max_j\{\sum_i \tilde{L}(i, j)\}\}$. Applying the Taylor approximation $\varphi(s) \approx p_r(s) := \sum_{n=0}^{r} \alpha_n s^n$ to the Laplacian ma-
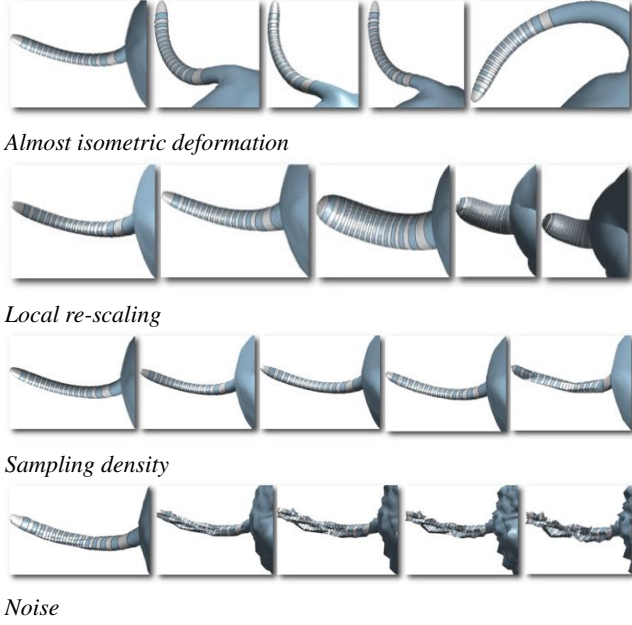
Almost isometric deformation

Local re-scaling

Sampling density

Noise

Figure 18: *Robustness of the computation of the linear FEM heat kernel from a seed point placed on the spike of the tail. The transformation strength increases from left to right.*

(a) $t = 0.1$ (b) $t = 1$

(c) $t = 0.1$ (d) $t = 1$

Figure 19: *Level sets of the linear FEM diffusion distance, computed using the Padé-Chebyshev approximation ($r := 7$), from a source point (black dot), with different values of t, on a (a,b) smooth and (c,d) noisy surface.*

trix in $[0, \lambda]$, $\mathbf{Ke}_i$ is evaluated as (Algorithm 2)

$$\mathbf{Ke}_i \approx \sum_{n=0}^{r} \alpha_n (\mathbf{B}^{-1}\mathbf{L})^n \mathbf{e}_i = \alpha_0 \mathbf{e}_i + \sum_{n=1}^{r} \alpha_n \mathbf{g}_n, \quad (6)$$

where $\mathbf{g}_n$ satisfies the linear system $\mathbf{Bg}_{n+1} = \mathbf{Lg}_n$, $\mathbf{Bg}_1 = \mathbf{Le}_i$.

From the upper bound [Pat14]

$$\left\| \varphi(\tilde{\mathbf{L}}) - \sum_{n=0}^{r} \alpha_n \tilde{\mathbf{L}}^n \right\|_2 \leq \frac{n}{(r+1)!} \left[\frac{\lambda_{\max}(\mathbf{L})}{\lambda_{\min}(\mathbf{B})} \right]^{r+1} \|\varphi^{(r+1)}(\tilde{\mathbf{L}})\|_2,$$

it follows that the approximation accuracy is mainly controlled by the degree of the Taylor approximation and the variation of the ratio between the maximum eigenvalue of $\mathbf{L}$ and the minimum eigenvalue of $\mathbf{B}$. If necessary, a higher approximation accuracy is achieved by slightly increasing the degree r. Finally, this computation of both the spectral kernel and distance is independent of the discretization of the input surface as a polygonal mesh or a point cloud. In case of a complex kernel, it is enough to apply the previous discussion to its real and imagery parts; e.g., for the wave kernel we consider the series $\sin(\tilde{\mathbf{L}}) = \sum_{n=0}^{+\infty} (-1)^n \tilde{\mathbf{L}}^{2n+1}/(2n+1)!$ and $\cos(\tilde{\mathbf{L}}) = \sum_{n=0}^{+\infty} (-1)^n \tilde{\mathbf{L}}^{2n}/(2n)!$.

Arbitrary filter: Padé-Chebyshev approximation For an arbitrary filter, we consider the rational Padé-Chebyshev approximation $p_r(s) = \frac{a_r(s)}{b_r(s)}$ of φ [GV89] (Ch. 11) with respect to the $\mathcal{L}_\infty$ norm. Here, $a_r(\cdot)$ and $b_r(\cdot)$ are polynomials of degree equal to or lower than r. Let $p_r(s) = \sum_{i=1}^{r} \alpha_i (1 + \beta_i s)^{-1}$ be the partial form of the Padé-Chebyshev approximation, where $(\alpha_i)_{i=1}^r$ are the weights and $(\beta_i)_{i=1}^r$ are the nodes of the r-point Gauss-Legendre quadrature rule [GV89] (Ch. 11). The weights and nodes are precomputed
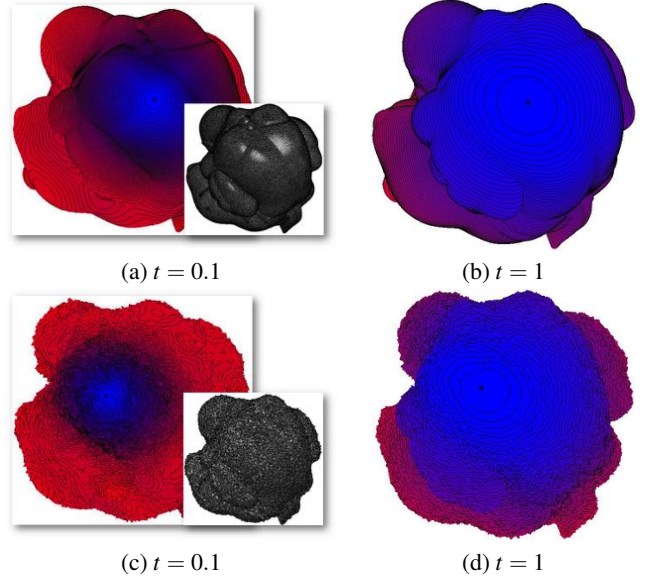
for any degree of the rational polynomial [CRV84]. Applying this approximation to the spectral kernel, we get that

$$\mathbf{u} = \mathbf{Kf} \approx p_r(\tilde{\mathbf{L}})\mathbf{f} = \sum_{i=1}^{r} \alpha_i \left(\mathbf{I} + \beta_i \tilde{\mathbf{L}} \right)^{-1} \mathbf{f} = \sum_{i=1}^{r} \alpha_i \mathbf{g}_i,$$

where $\mathbf{g}_i$ solves the symmetric and sparse linear system

$$(\mathbf{B} + \beta_i \mathbf{L}) \mathbf{g}_i = \mathbf{Bf}, \quad i = 1, \dots, r. \quad (7)$$

The Padé-Chebyshev approximation generally provides an accuracy higher than the polynomial approximation, as a matter of its uniform convergence to the filter.

Properties According to [MVL03], the approximation of the matrix $\varphi(\tilde{\mathbf{L}})$ might be numerically unstable if $\|\tilde{\mathbf{L}}\|_2$ is large. From the bound $\|\mathbf{B}^{-1}\mathbf{L}\|_2 \leq \lambda_{\min}^{-1}(\mathbf{B})\lambda_{\max}(\mathbf{L})$, a well-conditioned mass matrix $\mathbf{B}$ guarantees that $\|\mathbf{B}^{-1}\mathbf{L}\|_2$ is bounded. Recalling that $\mathbf{X}^\top (\mathbf{B} + \beta_i \mathbf{L})\mathbf{X} = (\mathbf{I} + \beta_i \Lambda)$, $\{1 + \beta_i \lambda_j\}_{j=1}^n$ are the eigenvalues of $(\mathbf{B} + \beta_i \mathbf{L})$ and its conditioning number is bounded by the constant $(1 + \beta_{\max} \lambda_n)$, $\beta_{\max} := \max_{i=1,\dots,n} |\beta_i|$. Indeed, the coefficient matrices in Eq. (6) are well-conditioned and specialized preconditioners [KFS13] can be applied to further attenuate numerical instabilities.

Approximating an arbitrary filter function with a rational or a polynomial function of degree r, the evaluation of the corresponding spectral distance between two points is reduced to solve r sparse, symmetric, linear systems (c.f., Eq. (6)), whose coefficient matrices have the same structure and sparsity of the connectivity matrix of the input triangle mesh or of the k-nearest neighbor graph for a point set. Applying iterative solvers, such as the Jacobi, Gauss-Seidel, minimum residual methods [GV89], and without extracting the Laplacian spectrum, the computational cost is $\mathcal{O}r\tau(n)$,
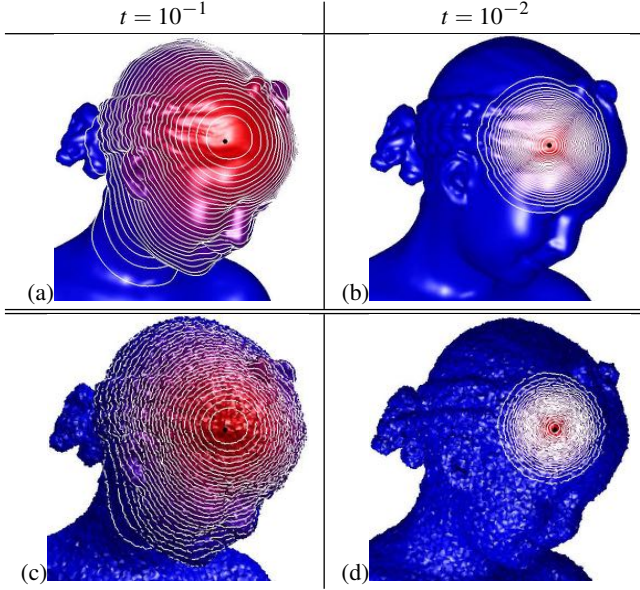
Figure 20: *Robustness of the Padé-Chebyshev approximation (r = 7) of the (a,c) diffusion kernel $\mathbf{K}_t\mathbf{e}_i$ and (b,d) distance at $\mathbf{p}_i$ (black dot) on a smooth and noisy triangulated surface.*

Figure 21: *(b,c) Robustness of the Padé-Chebyshev approximation (r = 5) of the heat kernel at different scales ($t = 10^{-1}, 10^{-2}$) with respect to (a) surface sampling (n = 5K, 20K).*

where $\tau(n)$ is the cost for the solution of a sparse linear system, which varies from $\mathcal{O}(n)$ to $\mathcal{O}(n^2)$, according to the sparsity of the coefficient matrix, and it is $\mathcal{O}(n\log n)$ in the average case.

The spectrum-free computation of the one-to-all distances $\{d(\mathbf{p}_i, \mathbf{p}_j)\}_{j=1}^n$ takes $\mathcal{O}(rn\tau(n))$ time; in fact, we solve the sparse linear system (6) with n different right-hand vectors $(\mathbf{e}_i - \mathbf{e}_j)$, $j = 1, \ldots, n$. Computing a fixed number k of eigenpairs in $\mathcal{O}(kn^2)$ time, the truncated spectral approximation of the one-to-all distance is evaluated in constant time for any filter. Indeed, the spectrum-free approach is competitive with respect to the truncated spectral approximation with $k(n) \geq r\tau(n)/n$ Laplacian eigenpairs. In the average case, $\tau(n) \approx n\log n$ and $k(n) \geq k_n$, $k_n = r\log n$. For instance, for a surface with $n = 10^4, 10^5, 10^6$ points and a degree $r = 5$, the number of eigenpairs is $k_n = 46, 58, 69$; in particular, this growth of k_n with respect to n is slow, as a matter of the logarithm in k_n.

6.5. Comparison and discussion

Fig. 13 reports the ℓ_∞ discrepancy (y-axis) between the diffusion distance on the sphere/cylinder and its approximation computed with the Padé-Chebyshev method and the truncated spectral approximation. In this case, the analytical expression of the Laplacian eigenfunctions on the sphere and cylinder has been used to compute the ground-truth distances [Pat16]. For small scales (e.g., $t = 10^{-2}, 10^{-3}$), the approximation error remains higher than 10^{-2}, with $k \leq 280$ eigenpairs; in fact, local shape features encoded by the heat kernel are recovered for a small t using the eigenvectors associated with high frequencies, thus requiring the computation of a large part of the Laplacian spectrum. For large scales (e.g., $t = 1, 10^{-1}$), increasing k strongly reduces the approximation error until it becomes almost constant and close to zero.
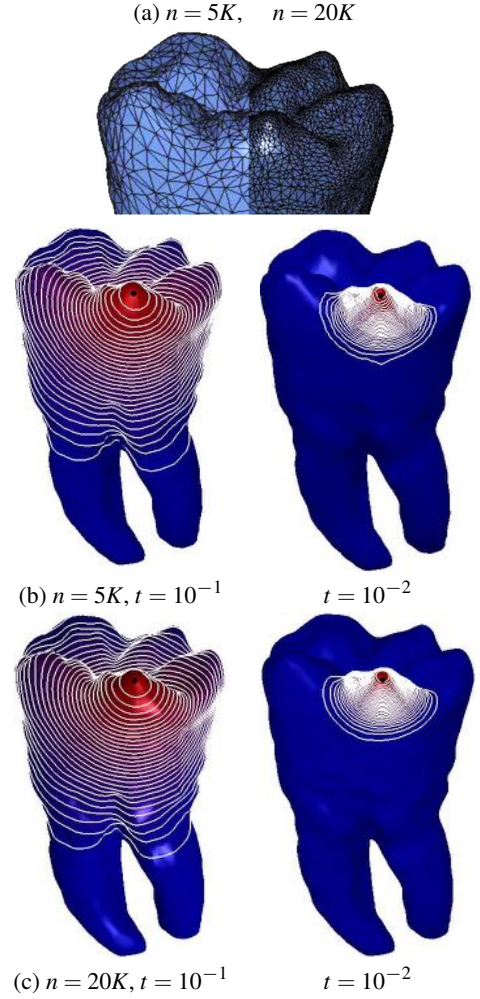
In this case, the behavior of the heat kernel is mainly influenced by the Laplacian eigenvectors related to the smaller eigenvalues. Indeed, the truncated spectral representation generally requires a high number of eigenpairs and does not achieve the approximation accuracy of our approach, which remains lower than 8.9×10^{-6} for all the scales. According to [VBCG10], there are no theoretical guarantees on the approximation accuracy of the heat kernel provided by multi-resolution prolongation operators. Furthermore, a low-resolution sampling of the input surface might affect the resulting accuracy.

For all the scales (Fig. 14), the accuracy of the Padé-Chebyshev method is higher than the truncated approximation with k eigenpairs, $k = 1, \ldots, 10^3$, the Euler backward method, and the power method. Reducing the scale, the accuracy of the Padé-Chebyshev remains almost unchanged while the other methods are affected by a larger discrepancy and tend to have an analogous behavior ($t = 10^{-4}$). Finally, the Euler backward method tends to over-
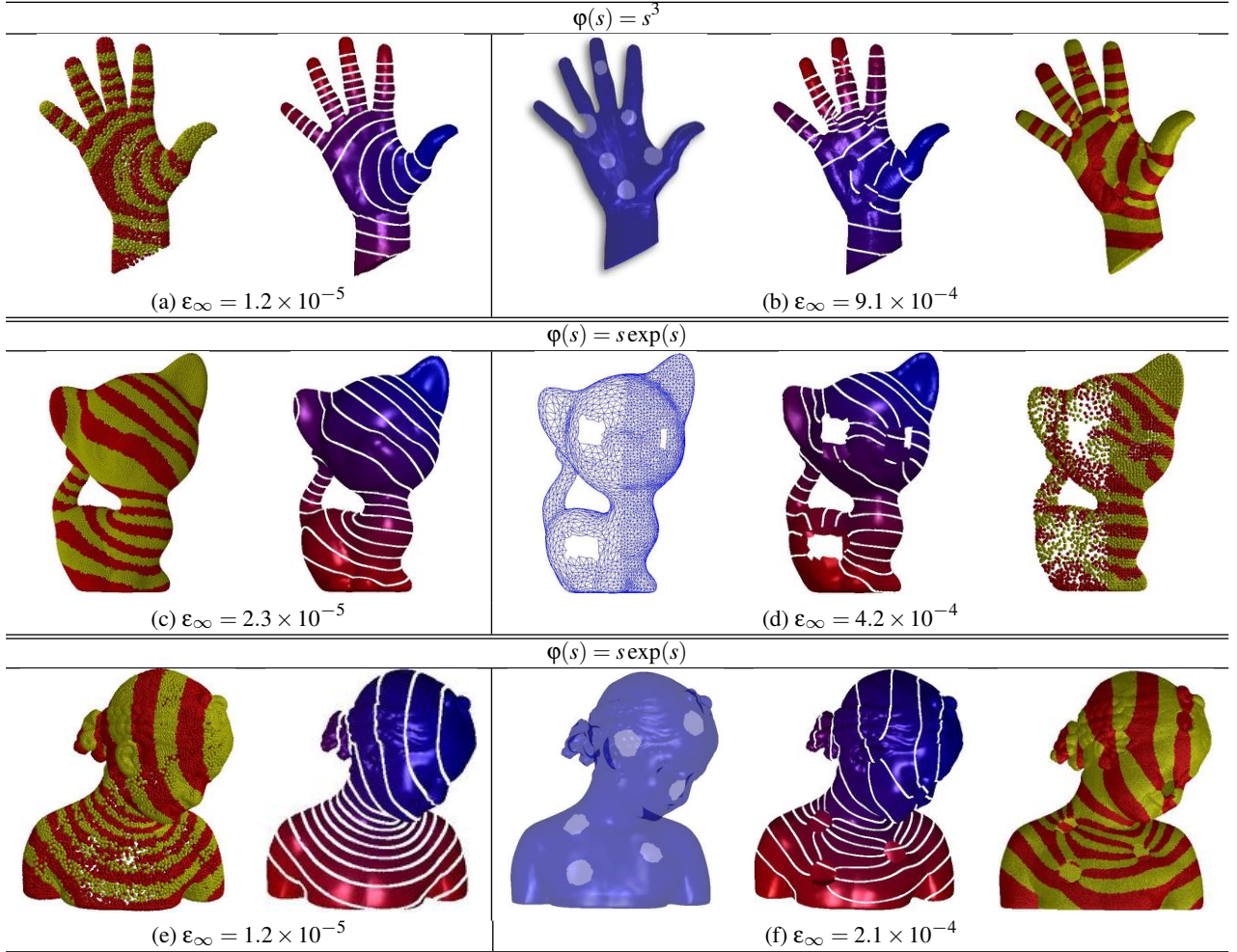
Figure 22: *Distances computed with the Padé-Chebyshev method ($r = 5$) on (a,c,e) regularly-sampled and (b,d,f) irregularly-sampled (left) meshes and (right) point sets with holes. To improve the visualization, points are represented as spheres.*

smooth the solution, which converges to a constant as $k \to +\infty$, and the selection of the power m is guided by heuristics.

The truncated spectral approximation of the diffusion distance is generally affected by the Gibbs phenomenon; i.e., small negative distance values. This phenomenon is more evident at small cases, which induce diffusion distances that decrease fast to zero and that are largely affected by small negative values. In fact, at small scales the diffusion distances decrease fast to zero and the negative values are no more compensated by the Laplacian eigenvectors related to smaller eigenvalues, as they are not included in the approximation (Fig. 15(e,f)). For the Padé-Chebyshev approximation (Fig. 15(a-d)), the distance values are positive at all the scales; in fact, we approximate the filter function without selecting a sub-part of the Laplacian spectrum.

Results in Figs. 16, 17 confirm that the diffusion distances at small scales generally require a number of eigenpairs that is much higher than the estimated value k_n. All tests have been performed on a 2.7 GHz Intel Core i7 Processor, with 8 GB memory. This case makes our computation of the one-to-all distance competitive with respect to its truncated approximation and useful to evaluate the distances for slowly-increasing (e.g., diffusion distances at small scales) or periodic filters or among seed points, as if happens for the evaluation of shape descriptors [OFCD02] and bags-of-features [BB11a, BBOG11]. Here, the number of seeds is much lower than the number of samples and the higher accuracy of our computation improves the discrimination capabilities of descriptors based on spectral distances.

In our experiments, the analogous behavior of the level-sets of the heat kernel and diffusion distance confirm the robustness of the Padé-Chebyshev of the approximation with respect to sampling, discrretization (Figs. 18, 19) and noise (Figs. 20, 21). A higher resolution of $\mathcal{M}$ improves the quality of the level-sets, which are always uniformly distributed and an increase of the noise magnitude does not affect the shape and distribution of the level sets.

Fixing the number of Laplacian eigenpairs makes the truncated spectral approximation of the one-to-all distances faster than ours
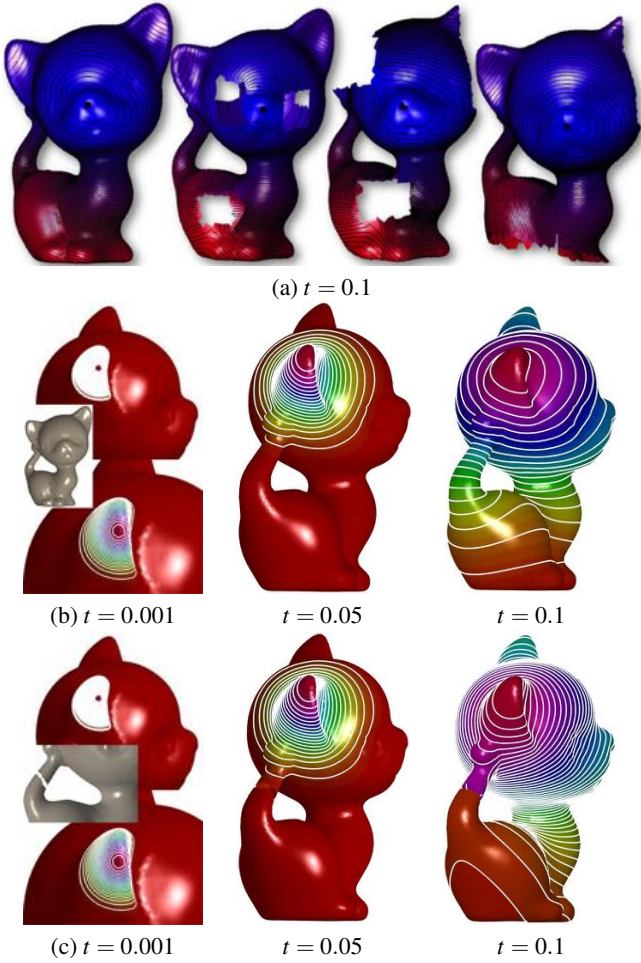
(a) $t = 0.1$

(b) $t = 0.001$ $t = 0.05$ $t = 0.1$

(c) $t = 0.001$ $t = 0.05$ $t = 0.1$

Figure 23: *Robustness of the Padé-Chebyshev approximation of the linear FEM (a) diffusion distance on partially-sampled surfaces and (b,c) heat kernel on smooth and topologically noisy surfaces (cut on the kitten tail), respectively.*

Figure 24: *(a) Noisy (m = 127, M = 57, s = 188) and (b) smoothed scalar function (m = 12, M = 14, s = 30). (c) Zoom-in on the level sets of the input (left) and smoothed (right) function.*

but generally provides a lower approximation accuracy. Slowly-increasing filters and small scales for the diffusion distances also require the computation of a large number of Laplacian eigenpairs, thus reducing the gap between the computational cost of the proposed approximation of the one-to-all distances and previous work. An analogous discussion applies to prolongation operators, which compute the truncated spectral approximation on a lower resolution of the input shape. Furthermore, previous work has not addressed methods for the selection of the proper number of eigenpairs with respect to the target approximation accuracy, which cannot be estimated without computing the whole Laplacian spectrum. Finally, Figs. 22, 23 show the robustness of the spectrum-free computation with respect to a different shape discretization, non-manifold and bordered surfaces, topological noise. At large scales only (e.g., $t = 1$), the shape of the level sets of the heat kernel changes in a neighbor of the topological cut.
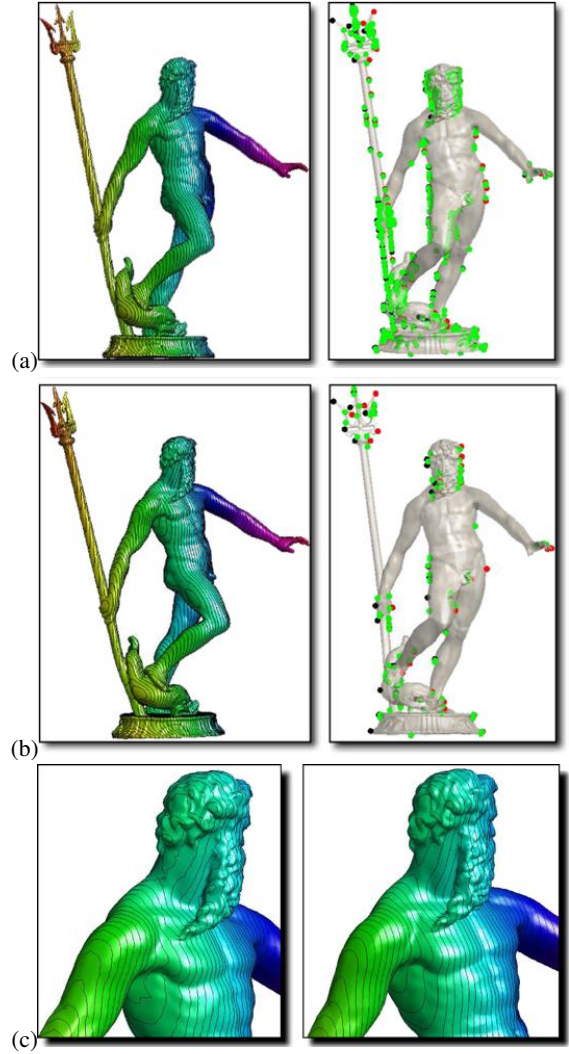
7. Applications

We now show how the Laplacian spectral properties and kernels have been used for smoothing in geometry processing (Sect. 7.1) and the definition of geometric basis functions (Sect. 7.2).

7.1. Smoothing

In real applications, the noisy component of the input data is due to a low quality of the discrete representations, unstable computations, and numerical approximations. Smoothing typically works in the function space and applies isotropic Laplacian filters [DBG*06, NGH04, Tau95] or bilateral smoothing operators to the function itself [LZ07]. The isotropy of the Laplacian matrix indiscriminately smooths noise and topological features [DBG*06, NGH04, Tau95] without constraints on their relocations or cancellations. Constrained least-squares techniques [SCOIT05] have been efficiently
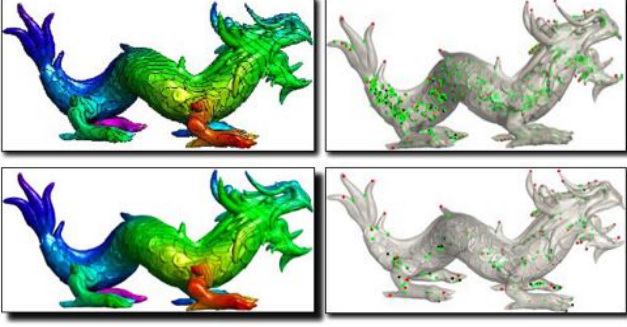
Figure 25: *Smoothing with interpolating constraints. The level sets and critical points of the input and smoothed scalar function are shown in the first and second row. The $\mathcal{L}_\infty$-error is 0.08.*

used to define compression schemes based on the selection of a set of anchors. While in [SCOIT05] the choice of the constrained vertices is guided by the final approximation accuracy of the reconstructed surface, in [PF09] the emphasis is on the preservation of the differential properties of f through the simplification of its critical points.

Unconstrained Laplacian smoothing According to [PF09], the smooth approximation $\tilde{f}$ of a noisy scalar function $f : \mathcal{M} \rightarrow \mathbb{R}$ can be computed as the compromise between approximation accuracy and smoothness of the solution, we minimize the energy $\mathcal{F}(\tilde{\mathbf{f}}) := \varepsilon\|\tilde{\mathbf{f}} - \mathbf{f}\|_{\mathbf{B}}^2 + \|\mathbf{L}\tilde{\mathbf{f}}\|_2^2$, whose normal equation is $(\mathbf{L}^\top \mathbf{L} + \varepsilon\mathbf{B})\tilde{\mathbf{f}} = \varepsilon\mathbf{B}\mathbf{f}$. Since the coefficient matrix is sparse and positive definite, $\tilde{\mathbf{f}}$ is uniquely defined and it is efficiently computed through direct or iterative solvers of sparse linear systems [GV89]. Finally, the spectral representation is $\tilde{\mathbf{f}} = \sum_{i=1}^n (\lambda_i^2 + \varepsilon)^{-1} \langle \mathbf{f}, \mathbf{x}_i \rangle_{\mathbf{B}} \mathbf{x}_i$, where the smoothing term $(\lambda_i^2 + \varepsilon)^{-1}$ filters out the contributions to the solution corresponding to the high eigenvalues (Fig. 24).

Laplacian smoothing with interpolating constraints According to [PSF07], we can include a set of constraints on the topological features of f to be preserved (Fig. 25). In this case, we define the smooth scalar function $\tilde{f}$ as the solution of the constrained minimization problem $\min_{\tilde{\mathbf{f}} \in \mathbb{R}^n} \|\mathbf{L}\tilde{\mathbf{f}}\|_2$, $\tilde{f}(\mathbf{p}_i) := f(\mathbf{p}_i), i \in \mathcal{I}$. This problem is equivalent to minimizing the least-squares error $\|\tilde{\mathbf{L}}\mathbf{x} - \mathbf{b}\|_2$, $\mathbf{x} \in \mathbb{R}^{n-k}$. Here, $\tilde{\mathbf{L}}$ is the $(n-k) \times (n-k)$ matrix achieved by removing the ith-row and i^{th}-column of $\mathbf{L}$, $i \in \mathcal{I}$, and the entries of the constant term $\mathbf{b} \in \mathbb{R}^{n-k}$ are $\sum_{j \in N(i) \cap \mathcal{I}} l_{ij} f(\mathbf{p}_j), i \in \mathcal{I}^C$.

Smoothing medical data In medical applications, the heat kernel is central in diffusion filtering and smoothing of images [ALM92, FCC92, PM90, SKS07, TWBO02, Wit83], 3D shapes [BX02, GSS99], and anatomical surfaces [CRD*05, KCS*12, WZS*13, Pat15]. Fig. 26 compares the diffusion smoothing of a noisy data set computed with the Padé-Chebyshev approximation of degree $r = 7$ and the truncated approximation with k Laplacian eigenparis. A low number of eigenpairs does not preserve shape details; increasing k reconstructs the surface noise. The ℓ_∞ error between (a) and the smooth approximation of (b) is lower than 1% for (c) the
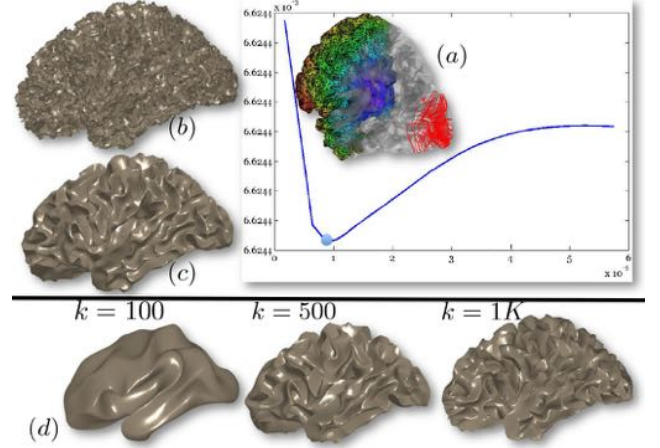

Figure 26: *(a) Input mesh and L-curve of the approximation accuracy (y-axis) versus the solution smoothness (x-axis). (b) Data set achieved by adding a Gaussian noise to (a). Diffusion smoothing computed with (c) the Padé-Chebyshev approximation ($r = 7$) and (d) the truncated approximation with k Laplacian eigenparis. A lower number of eigenpairs smooths local details; increasing k reconstructs the noisy component. The ℓ_∞ error between the ground-truth (a) and the smooth approximation of (b) is lower than 1% for the Padé-Chebyshev method (c) and varies from 12% ($k = 100$) to 13% ($k = 1K$) for the truncated approximation (d).*

Padé-Chebyshev method and (d) varies from 12% ($k = 100$) up to 13% ($k = 1K$) for the truncated spectral approximation.

7.2. Laplacian and diffusion basis functions

Even though the Laplacian eigenvectors are intrinsic to the input surface, they can be computed only for a small set of eigenvalues and do not provide a flexible alignment of the function behavior to specific shape features. The geometry-aware functions [SCOIT05] provide a computationally efficient way to encode the local geometric information of $\mathcal{M}$; and a similar approach can be applied to define more general classes of basis functions on a given shape. Applying the heat kernel matrix, we can define the *diffusion basis* $\mathcal{B} := \{\mathbf{K}_t \mathbf{e}_i\}_{i=1}^n$, whose elements have a smooth behavior on $\mathcal{M}$ and are intrinsically defined by $\mathcal{M}$ (Fig. 18). To define a set of *shape-driven* canonical basis functions, as feature points $\{\mathbf{p}_i\}_{i \in \mathcal{A}}$ of a 3D shape we select the maxima and minima of the Laplacian eigenfunctions related to the smallest eigenvalues [RPSS10] or of the auto-diffusion functions [GBAL09]. Finally, the definition of different basis function is also fundamental to define functions between shapes [GCO06, GMGP05, HK03, LG05, MS05, OFCD02, OBCS*12, RPSS10].

8. Conclusions

Our survey provides a common background on the definition and numerical computation of Laplacian spectral kernels and distances for geometry processing and shape analysis, as a generalization of

the well-known biharmonic, diffusion, and wave distances. To support the reader in the selection of the most appropriate with respect to shape representation, computational resources, and target application, all the reviewed numerical schemes have been discussed and compared in terms of robustness, approximation accuracy, and computational cost.

From the numerical point of view, the evaluation of full shape descriptors (e.g., heat kernel values among all the input points) is partially limited in case of densely sampled shapes, due to the expensive computational time and storage overhead. Indeed, the appropriate selection of seed points on the input domain and the conversion of the spectral descriptor to a sparse approximation are still crucial steps for the evaluation of full shape descriptors. Furthermore, the robustness of the spectrum-free computation to sampling and missing parts suggests the use of the spectral distances and descriptor for partial shape matching.

From the point of view of the definition of the spectral kernels and distances, we have discussed the general properties of the filter that guarantee their well-posedness, intrinsic and invariance properties. A deeper analysis of the filter properties with respect to the induced spectral distances would be beneficial to improve current results on surface watermarking and shape comparison. Finally, learning the filter from the geometric properties of a given class of data is an efficient, but partially unexplored, way to address shape segmentation, comparison, and more generally manifold learning.

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APPENDIX - Giuseppe Patanè's biography

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Giuseppe Patanè is researcher at CNR-IMATI (2001-today) and member of the Shape Modelling Group. He received a Ph.D. in "Mathematics and Applications" from the University of Genova (2005) and a Post-Lauream Degree Master from the "F. Severi National Institute for Advanced Mathematics" (2000), University of Milano-Bicocca. From 2001, his research activities have been focused on the definition of paradigms and algorithms for modeling and analyzing digital shapes and multidimensional data. The main results of these research activities have been published in 37 journal papers, in 27 conference papers, and presented in 2 SIGGRAPH ASIA Courses [PLG14, PLG13], 1 Eurographics Course [BFF*07], and 1 SMI Course [PGLS12]. Finally, he was organizer of the following courses

- "*An Introduction to Ricci Flow and Volumetric Approximation with Applications to Shape Modeling*" [PLG14], at SIGGRAPH ASIA 2014, December, 3-4, 2013, Shenzen, China. *Speakers:* G.Patanè, X.D. Gu, X.S. Li.
- "*Surface- and Volume-Based Techniques for Shape Modeling and Analysis*" [PLG13], at SIGGRAPH ASIA 2013, November, 19-22, 2013, Hong Kong, China. *Speakers:* G.Patanè, X.D. Gu, X.S. Li.
- "*3D Shape Modeling and Analysis Through Functional and Metric Spaces*", during the *3-D Geometry/Imaging Summer School*, July, 16-August, 3, 2012, Kunming University of Science and Technology, Yunnan, China. *Speaker:* G.Patanè.
- "*Spectral, Curvature Flow Surface- and Volume-Based Techniques for Shape Modeling and Analysis*", Shape Modeling International, 2012, May, 22-25, 2012, College Station, Texas, USA. *Speakers:* G.Patanè, X.D. Gu, X.S. Li, M. Spagnuolo.

For details, we refer to the personal web-page ‡ and CV §.

Main papers of the author related the STAR topics

- Patanè G., *Accurate and Efficient Computation of Laplacian Spectral Distances and Kernels*. In: Computer Graphics Forum. In press, 2016.
- Patanè G., *Volumetric Heat Kernel: Padè-Chebyshev Approximation, Convergence, and Computation*. In: Computer&Graphics, Volume 46, February 2015, Pages 64-71.
- Patanè G., *Diffusive Smoothing of 3D Segmented Medical Data*. In: Journal of Advanced Research, Elsevier, Volume 6, Issue 3, May 2015, Pages 425-431.
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