

# 1 Optimization of a New Reactive Force Field for Silver-Based 2 Materials

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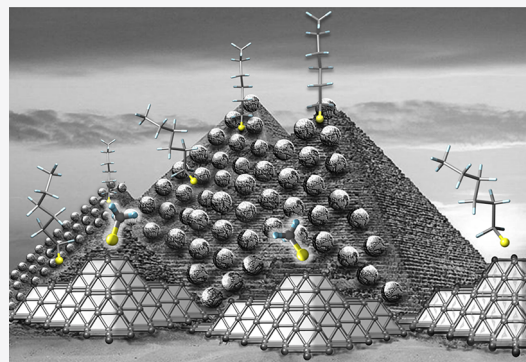


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4 **ABSTRACT:** A new reactive force field based on the ReaxFF formalism is  
5 effectively parametrized against an extended training set of quantum chemistry  
6 data (containing more than 120 different structures) to describe accurately  
7 silver and silver–thiolate systems. The results obtained with this novel  
8 representation demonstrate that the novel ReaxFF paradigm is a powerful  
9 methodology to reproduce more appropriately average geometric and  
10 energetic properties of metal clusters and slabs when compared to the earlier  
11 ReaxFF parametrizations dealing with silver and gold. ReaxFF cannot describe  
12 adequately specific geometrical features such as the observed shorter distances  
13 between the under-coordinated atoms at the cluster edges. Geometric and  
14 energetic properties of thiolates adsorbed on a silver  $\text{Ag}_{20}$  pyramid are  
15 correctly represented by the new ReaxFF and compared with results for gold.  
16 The simulation of self-assembled monolayers of thiolates on a silver (111)  
17 surface does not indicate the formation of staples in contrast to the results for gold–thiolate systems.



## 1. INTRODUCTION

18 Self-assembled monolayers (SAMs) of organic molecules on  
19 metallic surfaces can be used for creating nanoimprints,  
20 biosensors, protective layers, biomimetic covers, and stabilizers  
21 to favor specific morphologies. They also serve to manipulate  
22 topological, chemical, and functional features and to design  
23 complex structures that are important in a wide variety of  
24 applications. Among SAMs, thiolate layers on gold and silver  
25 surfaces have received major attention.<sup>1,2</sup> In particular, these  
26 SAMs can stabilize effectively the metal nanoparticles and  
27 improve their properties when functionalized with hydrophilic  
28 or hydrophobic moieties.<sup>3</sup> Recent investigations on SAMs  
29 adsorbed on gold surfaces have revealed the formation of S–  
30 Au–S staple motifs, which were already predicted by the  
31 theoretical methodologies.<sup>4,5</sup>

32 The computational characterization of SAMs by quantum  
33 chemical methods such as DFT<sup>6–8</sup> is usually limited to  
34 nanoparticles where the size is lower than 3 nm or small  
35 periodic slabs, but it becomes impractical and too demanding  
36 when both slab and nanoparticle dimensions increase to 5 nm  
37 and beyond. In this case, two different approaches are usually  
38 employed: simulations based on force fields<sup>9–13</sup> and tight-  
39 binding DFT calculations.<sup>14–16</sup> However, due to the special  
40 features of SAMs, to the metals and their mutual interactions,  
41 *ad hoc* force field parameters and formulations have been  
42 developed to obtain realistic representations. Bond-order  
43 reactive force fields are the most appropriate solution and  
44 can be designed for describing and predicting the correct  
45 adsorption locations of the SAMs, the formation of staples,

surface reconstructions, and other effects<sup>17–21</sup> important for  
the properties of these functionalized interfaces. 47

Regarding the gold/thiolate system,<sup>19–21</sup> three reactive force  
fields have been proposed in the past and used for various  
types of simulations, including bioconjugate materials for  
biomedical applications.<sup>22</sup> The excellent results, in line with  
the experiments, allowed the disclosure of different aspects of  
thiolate and staple formation. The success of this type of  
description for many other metals, such as copper,<sup>23</sup> silver<sup>24</sup>  
and cobalt<sup>25</sup> is apparent from the literature, as well. To the best  
of our knowledge, no ReaxFF force field has been developed  
yet to characterize comprehensively silver–thiolate materials  
and, more specifically, silver/thiolates nanocrystals. This would  
open up new atomistic simulations of realistic silver nano-  
particles in various types of environmental conditions close to  
the experimental setup and will allow one to better understand  
their properties and all the mechanisms connected to their  
diffusion, reactivity, and response to external stimuli. 63  
Furthermore, it will lead on the long-term to a possible  
improvement and tuning of their characteristics important for  
many applications. Silver nanoparticles are indeed widely used  
for their antimicrobial and plasmonic properties.<sup>26,27</sup> 67

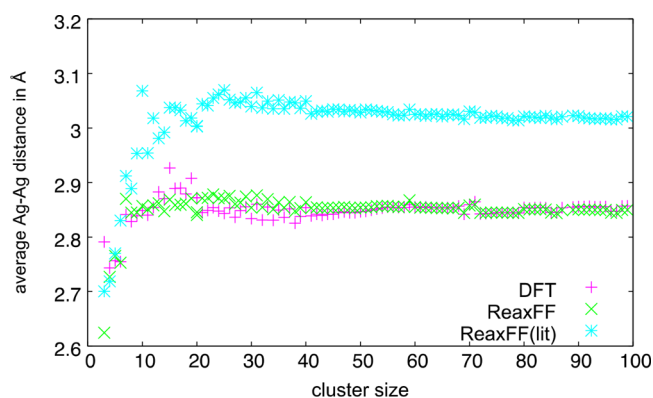
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68 Gold and silver have many similar properties. Due to  
69 relativistic effects, the distance between the atoms in bulk gold  
70 and silver with fcc structure is close (2.87 Å).<sup>28</sup> Moreover, the  
71 adsorption of thiolates on these metals follows similar  
72 mechanisms occupying, preferentially, the bridge sites on flat  
73 (111) surfaces in DFT calculations.<sup>8</sup> However, it is still an  
74 open question if the reorganization usually observed for gold<sup>5</sup>  
75 also occurs for silver. Moreover, for gold and silver, essential  
76 differences appear. For example, gold is more deformable than  
77 silver, and thus gold clusters show a broader range of atomic  
78 nearest-neighbors distances compared to silver.<sup>29</sup> Moreover,  
79 SAM thiolates on silver do not adopt the classical hexagonal  
80 superstructure observed on gold.<sup>1</sup>

81 The reproduction of all these features is a challenge that  
82 cannot be solved if only classical nonreactive force fields are  
83 considered. This is why we have decided to develop a new  
84 reactive force field for silver/thiolate systems and to compare  
85 the simulations results with published data to validate the  
86 improvements of our description.

87 The initial ReaxFF parameters of silver were extracted from  
88 the literature.<sup>24</sup> A preliminary check of the old parameters  
89 revealed that the silver clusters obtained at the DFT level could  
90 not be reproduced by the published force field (see Figure 1).



**Figure 1.** Average first-neighbor distances between silver atoms as a function of the cluster size (from 2 to 99 silver atoms) for DFT, the published reactive force field (ReaxFF(lit)),<sup>24</sup> and the one developed here (ReaxFF). The root-mean-square deviations of the ReaxFF(lit) and ReaxFF results with respect to the DFT data are 0.171 and 0.023 Å, respectively.

91 Thus, all the parameters were reoptimized against a database of  
92 DFT data consisting of published structures and new  
93 optimized systems. Besides the metal, we optimized the  
94 parameters of Ag–S using new DFT data. The results for silver  
95 and silver–thiolate systems are compared with those obtained  
96 for the published gold reactive force fields. Finally, the results  
97 of the molecular dynamics simulations of SAMs on gold and  
98 silver are presented.

## 2. METHODS

99 **2.1. The ReaxFF Model.** In the ReaxFF approach, the  
100 energy of a system is calculated with the following  
101 equation:<sup>30,31</sup>

$$E_{\text{system}} = E_{\text{bond}} + E_{\text{angle}} + E_{\text{tors}} + E_{\text{auto}} + E_{\text{Coulomb}} + E_{\text{vdWaal}} \\ + E_{\text{Hbond}} + E_{\text{conjugation}} + E_{\text{lon-pair}} + E_{\text{over}} + E_{\text{under}}$$

$E_{\text{bond}}$ ,  $E_{\text{angle}}$  and  $E_{\text{tors}}$  correspond to the two-, three- and four-  
body intramolecular terms, respectively.  $E_{\text{auto}}$  is the energy  
difference of neutral and charged atoms.  $E_{\text{Coulomb}}$ ,  $E_{\text{vdWaal}}$ , and  
 $E_{\text{Hbond}}$  represent the electrostatic, van der Waals and hydrogen-  
bond interactions. The term  $E_{\text{conjugation}}$  represents a possible  
conjugation between three or four atoms connected by  
alternating double bonds.  $E_{\text{lon-pair}}$ ,  $E_{\text{over}}$  and  $E_{\text{under}}$  represent  
nonbinding terms, a surplus, and a lack of valence electrons  
relative to the isolated species, respectively.

This force field differs from classical force fields by the fact  
that all the contributions, apart from the van der Waals and  
electrostatic terms, are expressed as a function of the bond  
order. This allows bond breaking and formation. A detailed  
explanation of all the terms can be found in the literature.<sup>30,31</sup>

In the case of metals some of the terms are set to zero (for  
example,  $E_{\text{angle}}$  and  $E_{\text{tors}}$ ).

Atomic charges can be calculated by the EEM method  
(*electronegativity equalization method*) introduced by Mortier et  
al.<sup>32,33</sup> or adopting the atom-condensed Kohn–Sham DFT  
approximated to second-order (ACKS2) description.<sup>34</sup> In this  
case, all the old parameters should be revised to guarantee  
compatibility with the new representation. Considering that  
the previous gold/silver force fields were based on the EEM  
method, we chose this for the new development. All the  
simulations and force field optimizations were carried out with  
the ReaxFF modules integrated in the SCM/ADF<sup>35</sup> and  
LAMMPS<sup>36,37</sup> packages.<sup>38</sup>

**2.2. The Parameterization of the AgSCH-ReaxFF  
Potential with the MCFF Method.** The force field  
optimization performed in this investigation was carried out  
with the Monte Carlo algorithm (MCFF)<sup>39</sup> implemented in  
the SCM/ADF code. Even though a new method, dubbed  
covariance matrix adaptation evolution strategy (CMA-ES),<sup>40</sup>  
has recently been published, its comparison with the MCFF  
methodology<sup>41</sup> has revealed that both tend toward very similar  
results when the training set is sufficiently large; thus, the  
MCFF procedure was chosen to stay in line with the recent  
Au/S/C/H parametrization strategy reported in ref 21. A brief  
description is provided to give an idea of the number of  
possible choices that are inherent to this procedure and could  
bias the final results. Briefly, the generic principle of MCFF is  
as follows:

- (1) A training set containing experimental and theoretical  
data is prepared (for example, molecular geometries  
optimized at different levels of accuracy, their relative  
energies, surfaces of potential energy representing the  
dissociation curve between two atoms, the atomic  
charges of atoms, etc.). A part of the results is reserved  
for a validation set which shows the reliability of the  
parametrized force field.
- (2) A set of parameters is chosen and the selected  
parameters are optimized by minimizing the unitless  
error function:

$$\text{Error} = \sum_{i=1}^n \left[ \frac{X_{i,TS} - X_{i,ReaxFF}}{\sigma_i} \right]^2$$

with  $n$  the number of points present in the training set,  $X_{i,TS}$  the  
value of point  $i$  in the training set,  $X_{i,ReaxFF}$  the corresponding  
value obtained by ReaxFF calculations, and  $\sigma_i$  the desired  
precision that the point  $i$  shall take.  $\sigma_i$  is equal to  $1/\omega_i$  with  $\omega_i$   
the actual weight of the point in the total error.

160 The acceptance of the new set of parameters is determined  
161 by a probability defined as follows:

$$P_{\text{acceptation}} = \min[1, \exp(-\beta \Delta \text{Error})]$$

162 with  $\beta$  a unitless parameter controlling the sampling in the  
163 parameter space and  $\Delta \text{Error} = \text{Error}_{\text{new}} - \text{Error}_{\text{previous}}$  the  
164 difference between the new error of the system and the  
165 previous one.<sup>42,43</sup> The initial value for  $\beta$  and the linear  
166 temperature increase factor are determined by the method  
167 proposed by Shchygol et al.<sup>41</sup> Preliminary tests have shown  
168 that the use of eight replica is appropriate to increase the  
169 chances of getting closer to a minimum error.<sup>40</sup>

170 **2.3. DFT Calculations.** DFT calculations were carried out  
171 with the PBE functional.<sup>44</sup> The modLAN2DZ and 6-311+  
172 +G(2d,2p) basis sets were used for Au/Ag<sup>45</sup> and S, C, and  
173 H,<sup>46</sup> respectively. The calculations were performed with the  
174 Gaussian09 package.<sup>47</sup> Gaussian and not ADF has been used  
175 here for coherence with previous investigations.

176 These selections were based on data found in the literature.  
177 Indeed, in 2014 Muniz-Miranda et al.<sup>48</sup> investigated system-  
178 atically the behavior of more than 25 exchange-correlation  
179 functionals (GGA, meta-GGA, hybrid, ...), including PBE, on  
180 three aggregates of Au covered with sulfur and phosphorus  
181 ligands, with different morphologies, charges, and experimental  
182 structures. This study showed that the PBE functional, coupled  
183 with the basis of Gaussian functions with pseudopotentials  
184 modLAN2DZ, leads to the results closest to the experimental  
185 data, for both structures and energy differences between the  
186 highest occupied and the lowest unoccupied molecular orbitals  
187 (HOMO–LUMO). These data suggest that PBE is appropri-  
188 ate for the present molecular systems.

189 **2.4. Simulation Methods.** Molecular dynamics simu-  
190 lations in the NVE ensemble were carried out using the  
191 LAMMPS program.<sup>36,37</sup> The three force fields for gold–  
192 thiolate were available in the SCM/ADF package. The force  
193 field developed for the silver–thiolate systems is included in  
194 the [Supporting Information](#). In all the molecular dynamics  
195 simulations, a time step of 0.25 fs was employed and the  
196 temperature was controlled by the Berendsen thermostat with  
197 a damping constant of 5 fs. The total simulation time was  
198 approximately 100 ps.

### 3. RESULTS AND ANALYSIS

199 **3.1. Force Field Optimization for Silver.** 3.1.1. *Training*  
200 *Set and Selected FF Parameters.* To reparametrize the  
201 potential, we used the following data:

- 202 • Several optimized clusters structures with 2 to 99 silver  
203 atoms were obtained from the [Supporting Information](#)  
204 of the article by Chen et al.<sup>50</sup> (PBE/LANL2DZ). The  
205 clusters with 2 to 20 silver atoms were reoptimized at the  
206 PBE/modLANL2DZ level of theory. All the Ag–Ag first  
207 neighbor distances (<3.2 Å) and their Ag–Ag–Ag  
208 angles within the clusters made of 2 to 20 silver atoms  
209 were included in the training set together with the first  
210 neighbor distances in the 21–99 atom clusters (number  
211 of points in the training set: ~22000).
- 212 • The cluster energies per atom for Ag<sub>2</sub> to Ag<sub>99</sub> were  
213 included in the training set. These are calculated by  
214 considering a reference. Two different references are  
215 used corresponding to the minimum energy structures  
216 for the clusters with <20 atoms and that of the larger  
217 ones.

$$E_{\text{cluster}} = \frac{E_{\text{Ag}_n}}{n} - \frac{E_{\text{Ag}_{n,\text{ref}}}}{n_{\text{ref}}}$$

- 218 • The potential energy surfaces (PES) describing the  
219 dissociation of the Ag<sub>2</sub> dimer were included in the  
220 training set together with a second PES, which describes  
221 the dissociation of a vertex atom of an Ag<sub>20</sub> pyramid.  
222 This was obtained by shifting the vertex atom of the  
223 pyramid along the central axis passing from the vertex.  
224 This was meant to reproduce the desorption of an  
225 adatom from a crystalline (111) surface of silver.  
226
- 227 • The equations of state of body-centered-cubic, face-  
228 centered-cubic, and hexagonal closed packed solids were  
229 added to the training set to obtain a good description of  
230 bulk silver. These data were obtained by Lloyd et al.<sup>24</sup>  
231 who carried out DFT calculations at the PBE/DZP level  
232 (see ref 24 for more details on these computations).<sup>51</sup>

233 The training set and its geometric files are given in the  
234 [Supporting Information](#). The general parameters for the  
235 optimization are reported in [Table 1](#), which shows the weights

**Table 1. Chosen Accuracy in Terms of  $\sigma_i$  and Weights of the Data Points in the Training Set Consisting of Interatomic Distances (Å), Angles (deg), and Energies (kcal/mol)**

type	$\sigma_i$	wt on the total error ( $1/\sigma_i^2$ )
distances	0.1 (Å)	100.0 (Å <sup>-2</sup> )
angles	5°	(1/5°) <sup>2</sup>
energies	1.0 (kcal/mol)	1.0 ((kcal/mol) <sup>-2</sup> )

usually used in this kind of parametrization.<sup>21,39,41</sup> The values  
236 of  $\sigma_i$  correspond to the optimization accuracy. The weights of  
237 some data have been modified to obtain the best reproduction  
238 of the DFT values (see trainsets in [Supporting Information](#) for  
239 the actual values). To ensure that the minimum of the PES is  
240 well reproduced, only those values close to the minimum were  
241 weighted 1 (kcal/mol)<sup>-2</sup>, whereas at shorter and larger  
242 distances smaller weights were used. To reach a good  
243 reproduction of the energies for large clusters, the weight of  
244 these data has also been increased.  
245

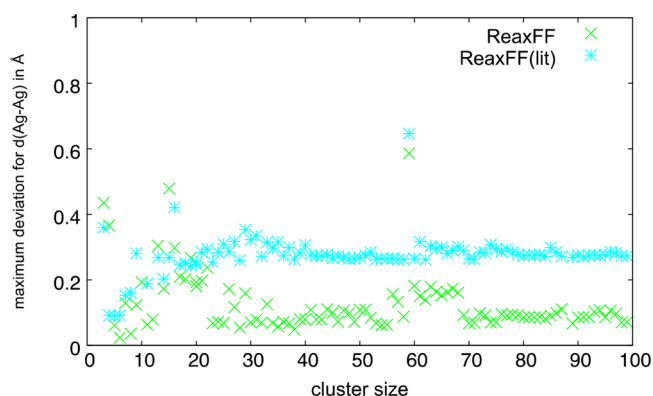
The EEM parameters of the literature have been used in the  
246 following.<sup>17,23</sup> Since the charges on metallic clusters are very  
247 small, these parameters have only a slight influence on the  
248 parametrization. The FF optimization was focused on bonds,  
249 angles, and binding energies. Fifteen parameters, correspond-  
250 ing to atomic and interatomic terms, were optimized.  
251

3.1.2. *Performance of the New Silver FF in Relation to the*  
252 *Previous Descriptions.* Starting from the silver force field  
253 proposed in the literature,<sup>24</sup> the optimization produced a new  
254 force field with 11% of the initial error to reproduce the data  
255 contained in the training set. The silver force field reported in  
256 the literature<sup>24</sup> will be denoted as ReaxFF(lit) in the following.  
257 In [Tables S1–S2](#), the new force field parameters are compared  
258 to the previous ones.  
259

[Figure 1](#) shows the evolution of the average first neighbor  
260 distance as a function of the cluster size for the training data  
261 (DFT structures), for the literature (ReaxFF(lit)),<sup>24</sup> and the  
262 new force fields (ReaxFF). In [Figure S1](#) in [Supporting](#)  
263 [Information](#), the standard deviations are shown for the  
264 distances of each cluster. It shows that the new model better  
265 describes the data for clusters larger than 20 atoms.  
266

267 Nevertheless the new model systematically underestimates the  
268 spread of the distances.

269 The evolution of the average distance of the clusters  
270 obtained with the new ReaxFF presents the same trend as the  
271 training set data. In contrast, the average distances obtained  
272 with ReaxFF(lit) are too long compared to the DFT data  
273 except for the smallest clusters. The average distance converges  
274 to 2.86 Å for large clusters (for DFT) and ReaxFF in contrast  
275 to 3.03 Å, which is the value obtained for ReaxFF(lit). This  
276 improvement is also evident in the maximum deviations of the  
277 Ag–Ag distances (Figure 2).



**Figure 2.** Maximum deviation of the first-neighbor distances for each cluster relative to the training set for the published reactive force field (ReaxFF(lit))<sup>24</sup> and the one developed here (ReaxFF). It corresponds to the maximum absolute value of the differences between the distances obtained by ReaxFF and DFT.

278 It was noticed that the maximum deviation decreases from  
279 0.3 Å (ReaxFF(lit)) to about 0.1 Å (new ReaxFF) for the  
280 larger clusters. The three exceptions, clearly visible in the plots  
281 (Figure 2), represent the Ag<sub>3</sub>, Ag<sub>15</sub>, and Ag<sub>59</sub> clusters. A deeper  
282 inspection of these cases and the comparison between the  
283 DFT and ReaxFF geometries revealed that they appeared for  
284 the largest bond distances predicted by DFT for surface atoms.  
285 The Ag<sub>3</sub> ReaxFF description predicts a perfect equilateral  
286 triangle (bond length 2.63 Å), while at the DFT level the  
287 distance between two atoms (3.06 Å) is larger than the other  
288 ones (2.65 Å). For Ag<sub>15</sub>, four distances around 3.4 Å obtained  
289 by DFT were reduced by the ReaxFF to values around 2.9 Å.  
290 In the case of Ag<sub>59</sub> two atoms with a distance of 3.2 Å in DFT,  
291 were separated by ReaxFF (new distance: 3.67 Å). We  
292 anticipate that these differences may be explained by the

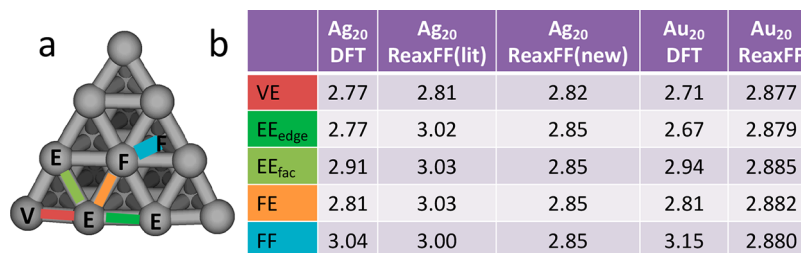
smaller range of the attraction between silver atoms predicted  
by ReaxFF.

In Figure 3, the DFT and ReaxFF geometries of the Ag<sub>20</sub> pyramid are compared. The structure can be completely defined by the reported distances between vertex (V), edge (E), and facet (F) atoms, which have 3, 6, and 9 first neighbors, respectively. It is apparent that the new ReaxFF does not correctly reproduce the variations of the distance within the silver pyramid, describing it as quasi-crystalline. However, also ReaxFF(lit) had problems with atom distances and overestimated the distances between the edge and facet atoms (3.00 and 3.03 Å), which disagrees with the reference DFT data. It is evident that ReaxFF cannot describe correctly the distances between the atoms with low coordination on edges and vertices which should be usually shorter. In the case of gold (ReaxFF by Järvi et al., 2008<sup>17</sup>) the interatomic distances are all very close to each other (2.88 Å). Similar to Ag<sub>20</sub>, the pyramid obtained with ReaxFF is a perfect crystal. It should be noted that all published force fields<sup>18–21</sup> for gold systems are based on this model.

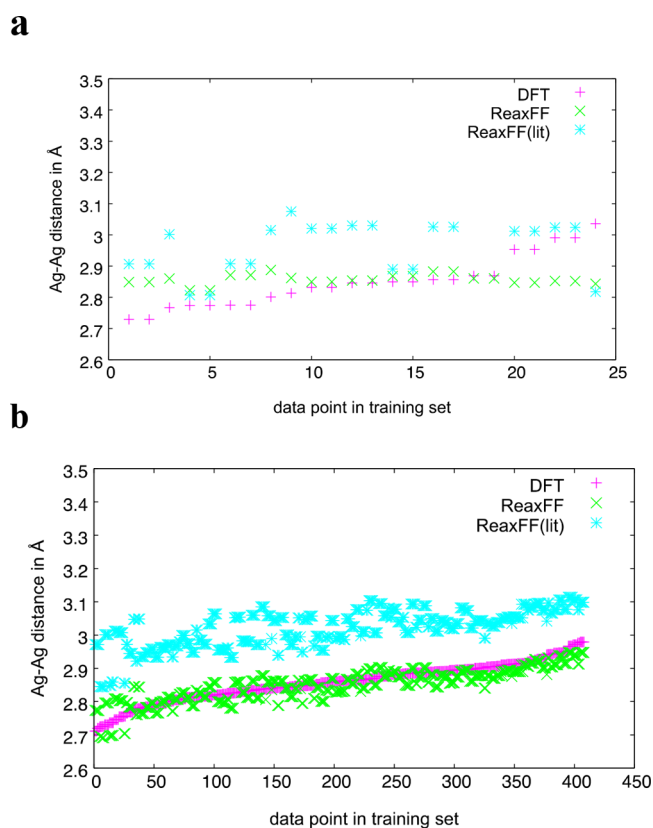
For all the clusters in the training set, the Ag–Ag distances obtained with the two force fields were compared to the DFT data. Figure 4 shows two typical results for a smaller (Ag<sub>10</sub>) and a large cluster (Ag<sub>98</sub>). The geometry of both clusters is better reproduced by the new ReaxFF. As already observed for the pyramid, the variation of distances for the small cluster obtained by the ReaxFF is more limited compared to DFT. For the larger clusters, the hierarchy of distance lengths is better respected.

To study the reliability of new ReaxFF, we prepared a validation set made of all Ag–Ag–Ag angles of large clusters (~128000 points) which were not included in the training set. The total error obtained for this set is 30% lower than the one obtained with the published ReaxFF with a maximum angular deviation of 7°.

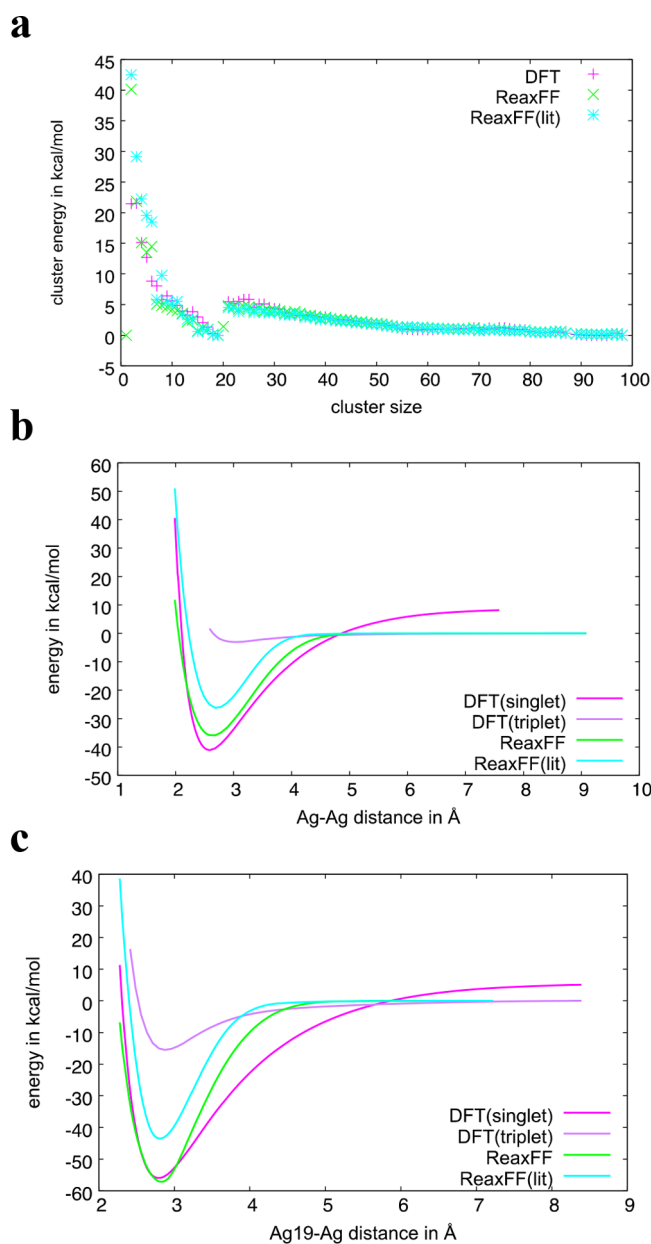
With regard to the reproduction of the energy data of the training set, the cluster energies and the potential dissociation energy curves of the dimer and the pyramid vertex atom are shown in Figure 5. It should be noted that the DFT curves for the dissociation of the dimer and the pyramid vertex atom were calculated taking the two possible electronic states, singlet and triplet, into account as done in Järvi et al.'s work on gold.<sup>17</sup> The dissociation energy curve should be close to the minimum of the singlet DFT curve (actually here 1.2 kcal/mol for Ag<sub>19</sub>–Ag) and tends to the results for the triplet at a larger distance. This is due to the fact that the triplet is more stable than the singlet at a larger distance. To obtain the agreement, we used in the training set the DFT data of the singlet curve close to 340



**Figure 3.** (a) Optimized geometry of the Ag<sub>20</sub> pyramid computed from the ReaxFF published in the literature.<sup>24</sup> The vertex (V), edge (E), and facet (F) atoms are shown. (b) Comparison of DFT and ReaxFF results for silver and gold. The Ag–Ag distances are given in Å. Due to symmetry the distances in part b are nonaveraged values which completely define the geometry of the pyramid. The standard deviations for the distances are 0.12 Å (Ag<sub>20</sub>, DFT), 0.095 Å (Ag<sub>20</sub>, ReaxFF(lit)), 0.013 Å (Ag<sub>20</sub>, ReaxFF), 0.19 Å (Au<sub>20</sub>, DFT), and 0.003 Å (Ag<sub>20</sub>, ReaxFF).



**Figure 4.** First-neighbor distances between silver atoms for two clusters as a function of the data point obtained by DFT, the published reactive force field (ReaxFF(lit)),<sup>24</sup> and the one developed here (ReaxFF): (a) Ag<sub>10</sub> and (b) Ag<sub>98</sub>. For the Ag<sub>10</sub> (Ag<sub>20</sub>) cluster the root-mean-square deviations of the ReaxFF(lit) and ReaxFF results with respect to the DFT data are 0.423 (0.170) and 0.085 (0.029) Å, respectively.



**Figure 5.** (a) Cluster energies (in kcal/mol) for all clusters of 2–99 atoms, (b) potential energy surfaces of the dissociation of a silver dimer, and (c) the dissociation surfaces of a vertex atom of the pyramid as a function of the distance between the vertex atom and its closest neighbor on the edge. Results obtained by DFT, the published reactive force field (ReaxFF(lit))<sup>24</sup> and the one developed here (ReaxFF). The root-mean-square deviations of the ReaxFF(lit) and ReaxFF with respect to the DFT data are 0.171 and 0.023 kcal/mol, respectively.

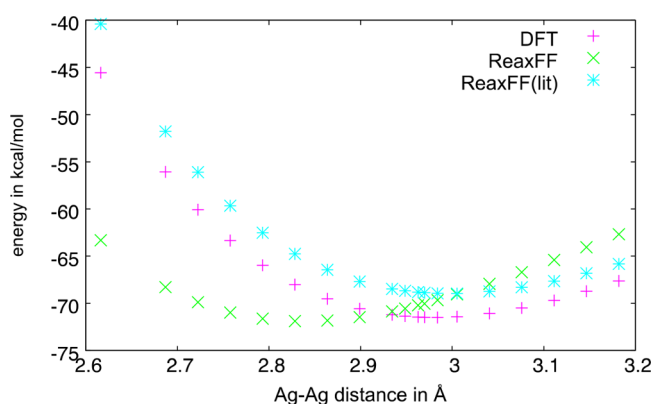
341 the minimum and those for the triplet at larger distances. For  
342 Ag–Ag, the deviation for the dissociation energy is quite larger  
343 (5.3 kcal/mol).

344 The new ReaxFF respects the hierarchy of cluster energies,  
345 slightly improving the absolute values of the energies of small  
346 silver clusters of 2 to 20 atoms relative to ReaxFF(lit). The  
347 dissociation energies of the dimer and the pyramid vertex atom  
348 are significantly improved at the minimum equilibrium energy  
349 level, with an error of 1.2 kcal/mol in the case of the pyramid  
350 compared to 15 kcal/mol for the previous model. Nevertheless,  
351 due to the exponential function modeling the dissociation of  
352 the bonds in the ReaxFF ( $E_{bond} = -D_e \cdot BO_{ij} \exp [p_{be,1} \cdot (1 -$   
353  $BO_{ij}^{p_{be,2}})]$ ), the range of the attraction peak is shorter compared  
354 with the DFT data. (As shown in Table S1 in the [Supporting](#)  
355 [Information](#), both parameters  $p_{be,1}$  and  $p_{be,2}$  have been  
356 optimized.) This could be a possible explanation of the  
357 problems observed in [Figure 1](#) for Ag<sub>3</sub>, Ag<sub>16</sub>, and Ag<sub>59</sub> as  
358 discussed above. The ReaxFF tends to connect or separate  
359 atoms that are at intermediate distances larger than 3 Å.

360 For the EOS energies of body-centered cubic, face-centered  
361 cubic, and hexagonal closed packed solids, the newReaxFF is  
362 quite correct. [Figure 6](#) illustrates the case of face-centered  
363 cubic solid. The Ag–Ag equilibrium position is 2.85 Å instead  
364 of 3.01 Å according to the data of the training set given by  
365 Lloyd et al.<sup>24</sup> (which we have integrated into our set, as  
366 discussed above). It should be emphasized that the DFT  
367 calculations by Lloyd et al. do not correctly describe the Ag–

Ag distance with respect to the experimental values (2.87 Å).<sup>28</sup>  
This explains the disagreement between DFT and ReaxFF(lit)  
observed for example for the Ag–Ag distance in the clusters  
(see [Figure 1](#)). The new Ag–Ag equilibrium position of 2.85 Å  
better agrees with the experiment and confirms the validity of  
the new ReaxFF.

**3.2. Ag–S Potential Parametrization. 3.2.1. Training**  
*Data and Parameter Selection.* To establish a training set for  
the adsorption of thiolate molecules on silver clusters and  
nanoparticles, we studied the MeS–Ag<sub>20</sub> isomers. The silver



**Figure 6.** Energy of expansion of a cubic-face centered solid as a function of the Ag–Ag distance obtained by DFT, the published reactive force field (ReaxFF(lit)),<sup>24</sup> and the one developed here (ReaxFF). The Ag–Ag equilibrium distance experimentally observed is 2.87 Å.<sup>28</sup>

pyramid was chosen since it meets the following criteria important for the parametrization:

- Its geometry is close to a face-centered cubic structure typical of many silver clusters and crystalline nanoparticles.
- It possesses specific sites: facets, edges, and vertices.
- It is large enough to behave as large clusters.
- It is sufficiently small to allow DFT calculations at a reasonable computational cost.

As illustrated in Figure 3, the pyramid has three distinguishable types of atoms: vertex atoms, edge atoms, and facet atoms. The geometry optimizations in the most stable electronic state (open-shell doublet) only led to the occupation of five adsorption sites by MeS, common to gold and silver: one on top (V) and four on bridge sites (EE<sub>edge</sub>, VE, FE, and EE). Here, the sites are named according to the atom types involved. For the FE and EE sites, a distinction between two isomers at the same site can be made according to the orientation of the methyl (up or down). This leads to seven different isomers (V, EE<sub>edge</sub>, VE, FE<sub>up</sub>, EE<sub>up</sub>, FE<sub>down</sub>, and EE<sub>down</sub>).

The geometric data considered in the training set were the Ag–S interatomic distances of the seven MeS–Ag<sub>20</sub> isomers, and the Ag–S–Ag and Ag–Ag–S angles, which defined the correct adsorption geometry of the thiolates on the metal part. The interatomic Ag–Ag distances were also retained in the training set to include the impact of the Ag–S interaction on the deformation of the silver pyramid in the error function.

For the energetic terms, the dissociation curves of a silver atom with a thiolate, the potential energy surface of the variation of the Ag–S–Ag angle on a silver dimer, and the relative energies of our MeS–Ag<sub>20</sub> isomers were included. These choices were motivated by the study of the training set used by Järvi et al.<sup>17,19</sup>

EEM parameters were not optimized.<sup>17,21</sup> From our quantum calculations, it was found that the Mulliken charges for the isomers were not appropriate. Indeed, for the isomer EE<sub>edge</sub>, for example, sulfur was positively charged (+0.6) and the metal atoms interacting with S negatively charged (−0.3), which is the opposite in the case of the Bader and NPA descriptions. Therefore, we decided to avoid the optimization of EEM parameters with Mulliken charges.

In total, the training set contained 1566 points. The weights were increased to focus the optimization on distances and angles. For the Ag–S distances and the Ag–S–Ag and Ag–Ag–S angles, an accuracy of 0.01 Å and 1°, respectively, was selected. On the contrary, the accuracy of the Ag–Ag and Ag–Ag–Ag angles within the pyramids were kept at 0.1 Å and 5.0° except for the bonds between two Ag atoms in direct contact with S (0.01 Å).

The initial set of parameters was extracted from the literature.<sup>21</sup> The Au–H and Au–C parameters obtained by Monti et al. were used for the Ag–H and Ag–C interaction. The training set considered in this work cannot be employed for optimizing these parameters. Only the interatomic parameters acting on the Ag–S bonds (distance and strength), without  $\pi$  and  $\pi$ – $\pi$  bonds, and the parameters describing the two types of Ag–S–Ag and Ag–Ag–S angles were improved. In total, 25 parameters were optimized. The training set and its geometric file is joined in the Supporting Information. In Tables S3–S6, the new force field parameters are shown.

### 3.2.2. Results and Comparison for the Ag–S Force Fields.

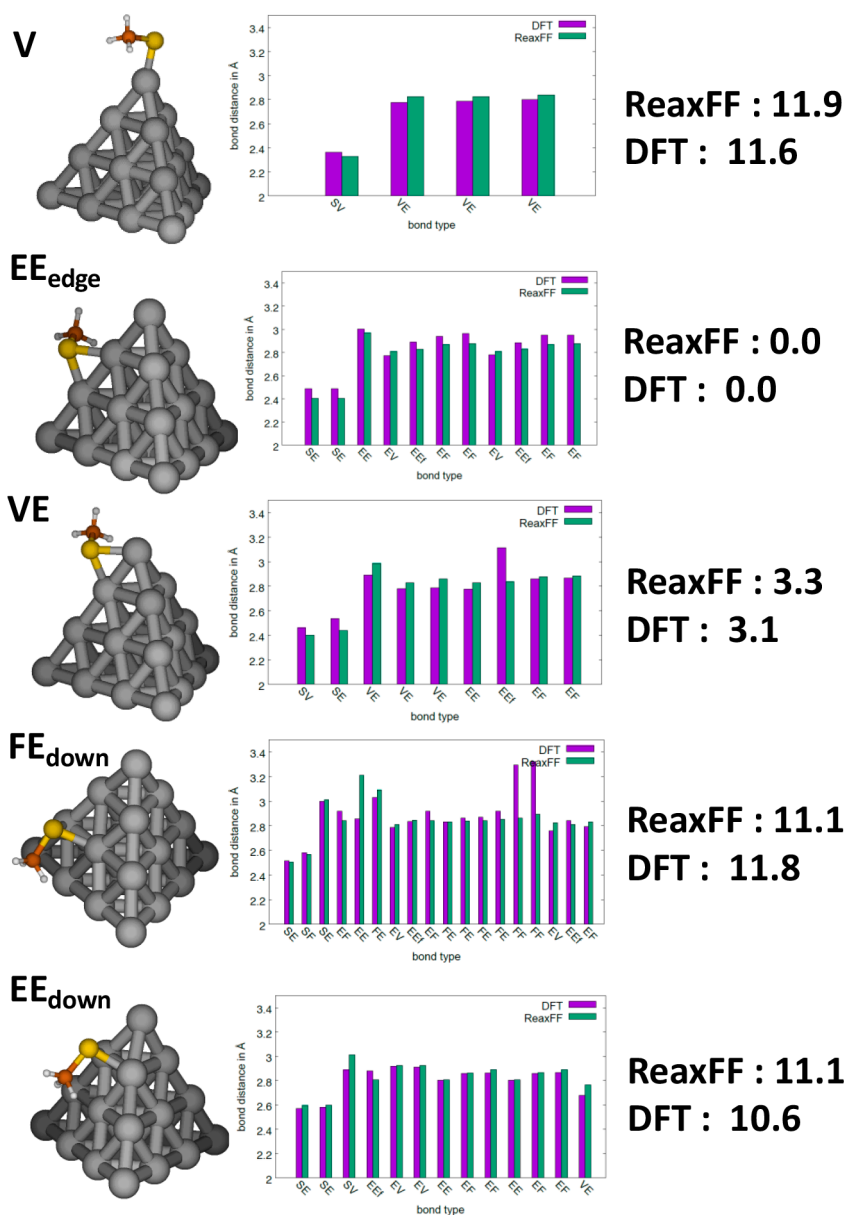
The best force field obtained corresponded to a reduction of the error around 15% (with respect to the parameters of Au–S). Figure 7 illustrates the geometric comparisons of the MeS–Ag<sub>20</sub> isomers and their relative energies between the ReaxFF and DFT data. Only the down isomers are shown for FE and EE, since the results for the up counterparts are very similar. The histograms from left to right represent the Ag–S distances and all the Ag–Ag interatomic distances of the silver atoms in contact with sulfur. The following results were obtained:

- The energy hierarchy is very well reproduced by ReaxFF with  $EE_{edge} < VE < FE_{down} \approx V \approx EE_{down}$ .
- According to the histograms of the interatomic distances, the Ag–S distances are well reproduced for all isomers (difference less than 0.1 Å).
- In the case of the bridged isomers, the elongations of the distance between the both Ag atoms in contact with sulfur are in good agreement with the DFT results.
- The global reorganization of the pyramid suffers from the tendency to favor crystalline descriptions. Thus, ReaxFF does not reproduce the extraction of facet atoms, predicting FF distances similar to the other Ag–Ag distances, see the two distances FF for the FE<sub>down</sub> isomer. For the same isomer, an elongation of the EE<sub>edge</sub> distance is predicted by ReaxFF which is not observed by DFT.

The potential energy curves of MeS–Ag dissociation obtained with the new force field and DFT are shown in Figure 8. In the case of AgS, the singlet is more stable than the triplet, and the results can be directly compared with this curve. The dissociation curve is well reproduced. The equilibrium of the Ag–S–Ag angle is around 71°, which corresponds to the angles found for the MeS–Ag<sub>20</sub> isomers by DFT.

To see if ReaxFF was capable of reproducing the differences between the adsorption on gold and silver as observed by DFT, we carried out the same comparison made in the case of the MeS–Au<sub>20</sub> isomers. All isomers obtained by DFT calculations were optimized with the three ReaxFF published<sup>19–21</sup> and the structural and energy data were collected in Table S1 in the Supporting Information. The three ReaxFF are denoted by the first author of the corresponding articles: Järvi, Bae and Monti.<sup>19–21</sup> Please note that for the ReaxFF of Bae and Aikins,

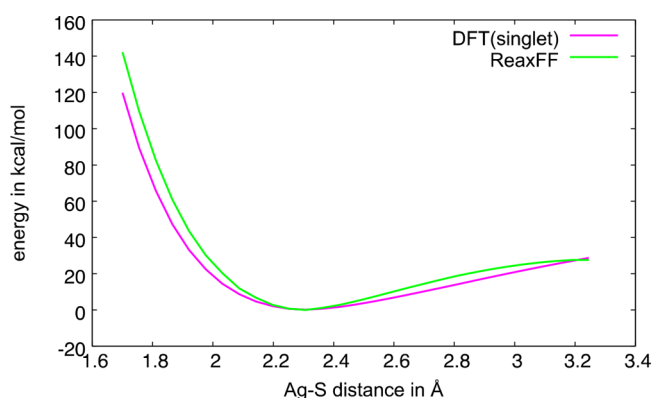
## Relative energies (kcal/mol)



**Figure 7.** Optimized geometries, relevant interatomic distances and relative energies (kcal/mol) obtained by ReaxFF and DFT for each MeS-Ag<sub>20</sub> isomers.

482 the parameters published in Table SI1 of this article<sup>20</sup>  
 483 (denoted by NP specific) have been used. These parameters  
 484 were chosen, since in this article it has been stated that they  
 485 should better work for nanoparticles. Figure S2 shows the final  
 486 structures obtained by the different methods. Usually, ReaxFF  
 487 does not reproduce exactly the geometry obtained with DFT,  
 488 while for silver a good agreement was found. This is not  
 489 surprising since the models were not optimized using thiolates  
 490 adsorbed on clusters but were optimized using thiolates  
 491 adsorbed on metallic surfaces. To clarify, we will discuss the  
 492 EE<sub>edge</sub> isomer as an example. In the case of gold, the Au–S  
 493 distance of this isomer should be 0.12 Å shorter than for silver,  
 494 while the bond between the Au atoms in contact with S should  
 495 be 0.25 Å larger. These differences are not reproduced by  
 496 ReaxFF, which obtains a similar distance for both metals. It is

interesting to note that the three ReaxFF can predict  
 qualitatively different structures. Thus, for the V isomer in  
 the case of the Monti potential, the Au–S interaction is  
 identical with that of our DFT reference. The Au–S bond is  
 vertical and parallel to the central rotation axis of the pyramid.  
 In the case of the other two potentials, the Au–S distance is  
 too large, and the axis of the Au–S bond is almost parallel to  
 an edge (S–Au(V)–Au(E) = 167°). Two different geometries  
 also appear during the optimization of the EE<sub>down</sub> isomer:  
 the first one, obtained from the Bae potential, does not have a  
 metal–ligand interaction, the ligand being separated from the  
 pyramid. In the case of the Järvi and Monti potentials, the  
 geometry obtained has an interaction site that is closer to a  
 hollow site than to the bridged site obtained in DFT.



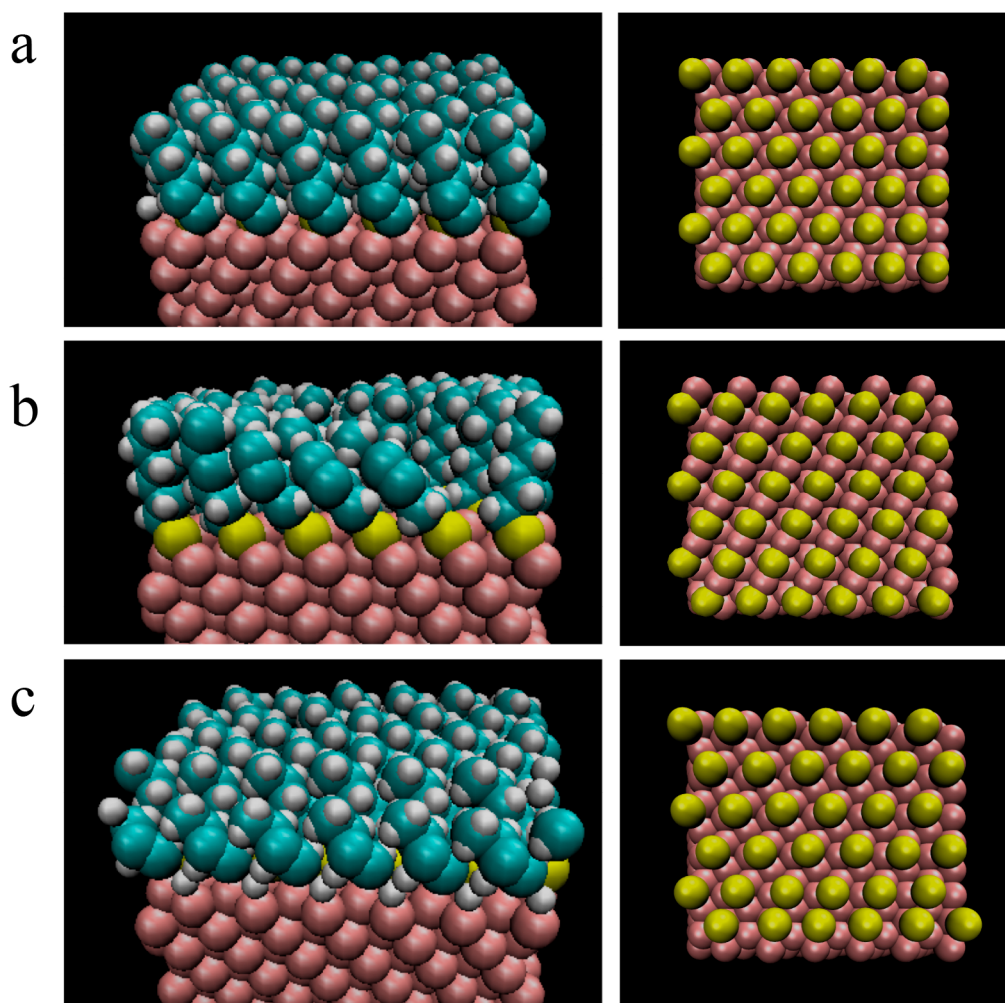
**Figure 8.** Potential energy surface of the dissociation of MeS-Ag (kcal/mol).

511 With respect to the energy data, the hierarchy of isomers  
512 calculated with Järvi potential is closest to the DFT hierarchy.  
513 That of Bae clearly shows a very bad reproduction of energies,  
514 in particular, the V isomer is the most stable one. Monti's  
515 potential seems to favor the bridge sites on the edges of the  
516 clusters. The energetic difference between ReaxFF and DFT  
517 are usually larger than those observed between silver and gold.

**3.3. Application of the Ag–Thiolate ReaxFF to Self-** 518  
**assembled Monolayers.** To validate and test the perform- 519  
ance of the new force field, molecular dynamics simulations 520  
were carried out to characterize the adsorption of thiulates 521  
assembled on the (111) facets of silver and gold slabs. The 522  
initial structures corresponded to one published by Groenbeck 523  
et al.,<sup>49</sup> which is a hexagonal assembly of four thiulates on the 524  
Au(111) surface. The size of this model was increased by 525  
replication (duplication in all directions) giving a final number 526  
of 36 thiulates. For comparison, the same configuration was 527  
employed for simulating the adsorption on silver, even though 528  
from experiments it was clear that other SAM arrangements 529  
could be more stable than the selected case. The aim of these 530  
calculations was just to see if the new ReaxFF was capable of 531  
predicting different structures starting from a similar SAM 532  
arrangement on two different metals. The simulations were 533  
carried out for methane and hexane thiulates. In the case of 534  
gold, the three published force fields were used, while for silver 535  
only the new ReaxFF was employed. 536

The initial structures were first minimized and then 537  
equilibrated at  $T = 300$  K, slowly increasing the temperature 538  
(Berendsen thermostat), in the NVT ensemble. 539

Figure 9 shows three snapshots at 100 K in the case of gold 540  
(Järvi and Monti FFs, Figure 9a,b) and silver surfaces (Figure 541



**Figure 9.** Snapshots of simulations of a SAM of hexanethiolate on (a) gold (ReaxFF proposed by Järvi et al.), (b) gold (ReaxFF proposed by Monti et al.), and (c) silver (new ReaxFF), at 100 K. On the lower right edge, the alkane chains have been removed to see the assembly of sulfur atoms.



542 9c) in the case of hexanethiolates. In the lower right edge, the  
543 SAM is shown without the hexane above. All snapshots for  
544 methane and hexane thiolate at 100 and 300 K are shown in  
545 the Supporting Information (Figures S3–S9). In the case of  
546 the NP-specific ReaxFF proposed by Bae and Atkins,<sup>20</sup> the  
547 thiolates desorbed when the temperature was increased,  
548 leading to the formation of S–S bonds (see Figure S9 in the  
549 Supporting Information). To check the validity of these  
550 simulations, they were carried out with LAMMPS and ADF  
551 using a slow increase of temperature and they both resulted in  
552 a similar desorption phenomenon. This suggests that this  
553 specific force field cannot be used for simulating thiolates on  
554 gold surfaces. It should be noted that Bae and Aikins<sup>20</sup> had  
555 already written that this force field should not be used for this  
556 kind of systems. We also verified if the same problem exists for  
557 the force field of Table 1 in the article by Bae and Aikins.<sup>20</sup> In  
558 the final SAM configuration, obtained by hundreds of  
559 picosecond NVE simulations, the thiolates were stably  
560 adsorbed on the surface, which shows that desorption only  
561 occurs with the NP-specific model (see Figure S10 in the  
562 Supporting Information). The snapshots obtained for gold–  
563 thiolates is markedly different from the Järvi and Monti  
564 ReaxFF representations. In the first case (and also for the force  
565 field by Bae published in Table 1 of their article), a well-  
566 ordered SAM is observed, but the gold atoms remain within  
567 the surface layer. On the contrary, the Monti potential allows  
568 reconstructions, some sulfur atoms leave the top layer and  
569 form chains of Au–S–Au staples. This is in agreement with  
570 experimental results indicating the formation of staples.<sup>2,5</sup>  
571 Another, reproduced feature was that the structure of SAMs  
572 did not change when the temperature was increased to 300 K.  
573 In the case of silver, a hexagonal SAM was observed but no  
574 indication of the formation of staples was apparent. As the  
575 temperature was increased from 100 to 300 K, the assembly  
576 became more irregular (see figures in the Supporting  
577 Information).

#### 4. CONCLUSIONS

578 In this long and laborious investigation, a new reactive force  
579 field, tuned to simulate silver- and silver–thiolate-based  
580 materials, is developed. The Ag–Ag potential has been  
581 improved with respect to the one published in the literature.  
582 using a large training set consisting of DFT data of clusters of  
583 different sizes. Several hundred tests were performed, to  
584 identify the atomic and interatomic parameters essential for the  
585 reproduction of the data. ReaxFF is capable of predicting  
586 correctly the evolution of the average distance between the  
587 silver atoms of all the clusters and average energetic properties  
588 even though it fails to reproduce the smaller distances of  
589 under-coordinated atoms located on the edges of clusters. We  
590 could speculate that only with the inclusion of other terms in  
591 the equations could these tendencies be remodulated.

592 Our new Ag–S potential is very efficient for exploring the  
593 adsorption of the SAMs, characterizing the dynamics of the  
594 thiolate–silver systems, and reproducing realistic structural and  
595 energetic properties of the isomers. Regarding the published  
596 ReaxFF for Au–S interactions, each potential correctly  
597 reproduces specific characteristics of the isomers: the Järvi  
598 potential gives a suitable energy hierarchy, the potential by Bae  
599 best reproduces the Au–S–Au angles, and the potential by  
600 Monti provides the best description of the Au–S distances.

601 To enable an easy use of the new force field, we added the  
602 text file of the force field and the directory of the simulation of

a SAM of hexanethiolates on Ag(111) (corresponding to 603  
Figure 9c) along with the input and output files for lammeps 604  
simulations to the Supporting Information. Even though the 605  
developed force field is not perfect, the results of the 606  
simulations are promising, in agreement with quantum 607  
chemistry calculations and experimental data. In future, it 608  
will be interesting to integrate other structures such as the 609  
Ag<sub>44</sub>(SR)<sub>30</sub> and Ag<sub>25</sub>(SR)<sub>18</sub> clusters or surfaces into the 610  
training sets to improve the proposed force field. 611

#### ■ ASSOCIATED CONTENT

##### Supporting Information

The Supporting Information is available free of charge at 614  
<https://pubs.acs.org/doi/10.1021/acs.jctc.0c00480>. 615

Tables S1–S6, the parameters of the new force field; 616  
Tables S7, results for MeS–Au<sub>20</sub> cluster obtained by 617  
ReaxFF and DFT; Figure S1, standard deviations of the 618  
distances for the clusters obtained by ReaxFF and DFT; 619  
Figures S2, structures of MeS–Au<sub>20</sub> obtained by ReaxFF 620  
and DFT; Figures S3–S9, snapshots of simulations of a 621  
SAM of methanethiolate on silver and gold for different 622  
force fields (PDF) 623

Training set and geometry files for Ag–Ag and Ag–S 624  
force field development: trainsetAgAg.in (TXT) 625

Training set and geometry files for Ag–Ag and Ag–S 626  
force field development: geoAgAg (TXT) 627

Training set and geometry files for Ag–Ag and Ag–S 628  
force field development: trainsetAgS.in and geoAgS 629  
(TXT) 630

Force field file: ffield.AgSCH (TXT) 631

Input and output files for Lammeps simulations of SAMs 632  
on silver: data.surf (TXT) 633

Input and output files for Lammeps simulations of SAMs 634  
on silver: in.surfGrad (TXT) 635

Input and output files for Lammeps simulations of SAMs 636  
on silver: log.lammeps (TXT) 637

Input and output files for Lammeps simulations of SAMs 638  
on silver: toto\_reax.xyz (TXT) 639

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##### Notes

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659

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## 665 ■ REFERENCES

666 (1) Woodruff, D. P. The interface structure of n-alkylthiolate self-  
667 assembled monolayers on coinage metal surfaces. *Phys. Chem. Chem.*  
668 *Phys.* **2008**, *10*, 7211–7221.  
669 (2) Häkkinen, H. The gold-sulfur interface at the nanoscale. *Nat.*  
670 *Chem.* **2012**, *4*, 443–455.  
671 (3) Daniel, M. C.; Astruc, D. Gold nanoparticles: assembly,  
672 supramolecular chemistry, quantum-size-related properties, and  
673 applications toward biology, catalysis, and nanotechnology. *Chem.*  
674 *Rev.* **2004**, *104*, 293–346.  
675 (4) Jadzinsky, P. D.; Calero, G.; Ackerson, C. J.; Bushnell, D. A.;  
676 Kornberg, R. D. Structure of a thiol monolayer-protected gold  
677 nanoparticle at 1.1 Å resolution. *Science* **2007**, *318*, 430–433.  
678 (5) Cossaro, A.; Mazzarello, R.; Rousseau, R.; Casalis, L.; Verdini,  
679 A.; Kohlmeier, A.; Scoles, G.; et al. X-ray diffraction and computation  
680 yield the structure of alkanethiols on gold (111). *Science* **2008**, *321*,  
681 943–946.  
682 (6) Torres, E.; Blumenau, A. T.; Biedermann, P. U. Mechanism for  
683 phase transitions and vacancy island formation in alkylthiol/Au (111)  
684 self-assembled monolayers based on adatom and vacancy-induced  
685 reconstructions. *Phys. Rev. B: Condens. Matter Mater. Phys.* **2009**, *79*,  
686 075440.  
687 (7) Grönbeck, H.; Häkkinen, H.; Whetten, R. L. Gold-thiolate  
688 complexes form a unique  $c(4 \times 2)$  structure on Au(111). *J. Phys.*  
689 *Chem. C* **2008**, *112*, 15940–15942.  
690 (8) Li, A.; Piquemal, J. P.; Richardi, J.; Calatayud, M. Butanethiol  
691 adsorption and dissociation on Ag (111): A periodic DFT study. *Surf.*  
692 *Sci.* **2016**, *646*, 247–252.  
693 (9) Luedtke, W. D.; Landman, U. Structure and thermodynamics of  
694 self-assembled monolayers on gold nanocrystallites. *J. Phys. Chem. B*  
695 **1998**, *102*, 6566–6572.  
696 (10) Rapino, S.; Zerbetto, F. Dynamics of thiolate chains on a gold  
697 nanoparticles. *Small* **2007**, *3*, 386–388.  
698 (11) Pool, R.; Schapotschnikow, P.; Vlugt, T. J. Solvent effects in the  
699 adsorption of alkyl thiols on gold structures: A molecular simulation  
700 study. *J. Phys. Chem. C* **2007**, *111*, 10201–10212.  
701 (12) Lane, J. M. D.; Grest, G. S. Spontaneous asymmetry of coated  
702 spherical nanoparticles in solution and at liquid-vapor interfaces. *Phys.*  
703 *Rev. Lett.* **2010**, *104*, 235501.  
704 (13) Djebaili, T.; Richardi, J.; Abel, S.; Marchi, M. Atomistic  
705 simulations of the surface coverage of large gold nanocrystals. *J. Phys.*  
706 *Chem. C* **2013**, *117*, 17791–17800.  
707 (14) Mäkinen, V.; Koskinen, P.; Häkkinen, H. Modeling thiolate-  
708 protected gold clusters with density-functional tight-binding. *Eur.*  
709 *Phys. J. D* **2013**, *67*, 38.  
710 (15) Fihey, A.; Hettich, C.; Touzeau, J.; Maurel, F.; Perrier, A.;  
711 Köhler, C.; Frauenheim, T.; et al. SCC-DFTB parameters for  
712 simulating hybrid gold-thiolates compounds. *J. Comput. Chem.*  
713 **2015**, *36*, 2075–2087.  
714 (16) Tarrat, N.; Rapacioli, M.; Cuny, J.; Morillo, J.; Heully, J. L.;  
715 Spiegelman, F. Global optimization of neutral and charged 20- and 55-  
716 atom silver and gold clusters at the DFTB level. *Comp. Comput.*  
717 *Theor. Chem.* **2017**, *1107*, 102–114.  
718 (17) Järvi, T. T.; Kuronen, A.; Hakala, M.; Nordlund, K.; Van Duin,  
719 A. C. T.; Goddard, W. A.; Jacob, T. Development of a ReaxFF  
720 description for gold. *Eur. Phys. J. B* **2008**, *66*, 75–79.  
721 (18) Keith, J. A.; Fantauzzi, D.; Jacob, T.; Van Duin, A. C. Reactive  
722 forcefield for simulating gold surfaces and nanoparticles. *Phys. Rev. B:*  
723 *Condens. Matter Mater. Phys.* **2010**, *81*, 235404.  
724 (19) Järvi, T. T.; Van Duin, A. C.; Nordlund, K.; Goddard, W. A., III  
725 Development of interatomic reaxff potentials for Au–S–C–H  
726 systems. *J. Phys. Chem. A* **2011**, *115*, 10315–10322.

(20) Bae, G. T.; Aikens, C. M. Improved ReaxFF Force Field  
Parameters for Au–S–C–H Systems. *J. Phys. Chem. A* **2013**, *117*,  
10438–10446.  
729  
(21) Monti, S.; Carravetta, V.; Ågren, H. Simulation of gold  
730 functionalization with cysteine by reactive molecular dynamics. *J.*  
731 *Phys. Chem. Lett.* **2016**, *7*, 272–276.  
(22) Monti, S.; Barcaro, G.; Sementa, L.; Carravetta, V.; Ågren, H.  
733 Dynamics and self-assembly of bio-functionalized gold nanoparticles  
734 in solution: Reactive molecular dynamics simulations. *Nano Res.*  
735 **2018**, *11*, 1757–1767.  
(23) Yeon, J.; Adams, H. L.; Junkermeier, C. E.; van Duin, A. C.;  
737 Tysoe, W. T.; Martini, A. Development of a ReaxFF force field for  
738 Cu/S/C/H and reactive MD simulations of methyl thiolate  
739 decomposition on Cu (100). *J. Phys. Chem. B* **2018**, *122*, 888–896.  
(24) Lloyd, A.; Cornil, D.; van Duin, A. C. T.; van Duin, D.; Smith,  
741 R.; Kenny, S. D.; Beljonne, D.; et al. Development of a ReaxFF  
742 potential for Ag/Zn/O and application to Ag deposition on ZnO.  
743 *Surf. Sci.* **2016**, *645*, 67–73.  
(25) Zhang, X.-Q.; Iype, E.; Nedeia, S. V.; Jansen, A. P. J.; Szyja, B.  
745 M.; Hensen, E. J. M.; van Santen, R. A. Site Stability on Cobalt  
746 Nanoparticles: A Molecular Dynamics ReaxFF Reactive Force Field  
747 Study. *J. Phys. Chem. C* **2014**, *118*, 6882–6886.  
(26) Morones, J. R.; Elechiguerra, J. L.; Camacho, A.; Holt, K.;  
749 Kouri, J. B.; Ramirez, J. T.; Yacaman, M. J. The bactericidal effect of  
750 silver Nanoparticles. *Nanotechnology* **2005**, *16*, 2346–2353.  
(27) Jain, P. K.; Huang, X.; El-Sayed, I. H.; El-Sayed, M. A. Noble  
752 Metals on the Nanoscale: Optical and Photothermal Properties and  
753 Some Applications in Imaging, Sensing, Biology, and Medicine. *Acc.*  
754 *Chem. Res.* **2008**, *41*, 1578–1586.  
(28) Weast, R. C.; Astle, M. J.; Beyer, W. H. *CRC handbook of*  
756 *chemistry and physics*, 69th ed.; CRC Press: Boca Raton, FL, 1988.  
(29) Greenwood, N. N.; Earnshaw, A. *Chemistry of the Elements*, 2nd  
758 ed.; 1997.  
(30) Van Duin, A. C.; Dasgupta, S.; Lorant, F.; Goddard, W. A.  
760 ReaxFF: a reactive force field for hydrocarbons. *J. Phys. Chem. A* **2001**,  
761 *105*, 9396–9409.  
(31) Chenoweth, K.; Van Duin, A. C.; Goddard, W. A. ReaxFF  
763 reactive force field for molecular dynamics simulations of hydro-  
764 carbon oxidation. *J. Phys. Chem. A* **2008**, *112*, 1040–1053.  
(32) Mortier, W. J.; Ghosh, S. K.; Shankar, S. Electronegativity-  
766 equalization method for the calculation of atomic charges in  
767 molecules. *J. Am. Chem. Soc.* **1986**, *108*, 4315–4320.  
(33) Janssens, G. O.; Baekelandt, B. G.; Toufar, H.; Mortier, W. J.;  
769 Schoonheydt, R. A. Comparison of cluster and infinite crystal  
770 calculations on zeolites with the electronegativity equalization method  
771 (EEM). *J. Phys. Chem.* **1995**, *99*, 3251–3258.  
(34) Verstraelen, T.; Ayers, P. W.; Van Speybroeck, V.; Waroquier,  
773 M. ACKS2: Atom-condensed Kohn-Sham DFT approximated to  
774 second order. *J. Chem. Phys.* **2013**, *138*, 074108.  
(35) van Duin, A. C. T.; Goddard, W. A.; Islam, M. M.; van Schoot,  
776 H.; Trnka, T.; Yakovlev, A. L. *ReaxFF 2019.3*; SCM, Theoretical  
777 Chemistry, Vrije Universiteit:Amsterdam, The Netherlands; [http://](http://www.scm.com)  
778 [www.scm.com](http://www.scm.com).  
(36) Plimpton, S. Fast Parallel Algorithms for Short-Range  
780 Molecular Dynamics. *J. Comput. Phys.* **1995**, *117*, 1–19.  
(37) Aktulga, H. M.; Fogarty, J. C.; Pandit, S. A.; Grama, A. Y.  
782 Parallel reactive molecular dynamics: Numerical methods and  
783 algorithmic techniques. *Parallel Computing* **2012**, *38*, 245–259.  
(38) van Duin, A. C.; Baas, J. M.; Van De Graaf, B. Delft molecular  
785 mechanics: a new approach to hydrocarbon force fields. Inclusion of a  
786 geometry-dependent charge calculation. *J. Chem. Soc., Faraday Trans.*  
787 **1994**, *90*, 2881–2895.  
(39) Iype, E.; Hütter, M.; Jansen, A. P. J.; Nedeia, S. V.; Rindt, C. C.  
789 M. Parameterization of a reactive force field using a Monte Carlo  
790 algorithm. *J. Comput. Chem.* **2013**, *34*, 1143–1154.  
(40) Trnka, T.; Tvaroska, I.; Koca, J. Automated Training of ReaxFF  
792 Reactive Force Fields for Energetics of Enzymatic Reactions. *J. Chem.*  
793 *Theory Comput.* **2018**, *14*, 291–302.  
794

- 795 (41) Shchygol, G.; Yakovlev, A.; Trnka, T.; van Duin, A. C. T.;  
796 Verstraelen, T. ReaxFF Parameter Optimization with Monte-Carlo  
797 and Evolutionary Algorithms: Guidelines and Insights. *J. Chem.*  
798 *Theory Comput.* **2019**, *15*, 6799–6812.
- 799 (42) Metropolis, N.; Rosenbluth, A. W.; Rosenbluth, M. N.; Teller,  
800 A. H.; Teller, E. Equation of state calculations by fast computing  
801 machines. *J. Chem. Phys.* **1953**, *21*, 1087–1092.
- 802 (43) Kirkpatrick, S.; Gelatt, C. D.; Vecchi, M. P. Optimization by  
803 simulated annealing. *Science* **1983**, *220*, 671–680.
- 804 (44) Perdew, J. P.; Burke, K.; Ernzerhof, M. Generalized gradient  
805 approximation made simple. *Phys. Rev. Lett.* **1996**, *77*, 3865.
- 806 (45) Couty, M.; Hall, M. B. Basis sets for transition metals:  
807 Optimized outer p functions. *J. Comput. Chem.* **1996**, *17*, 1359–1370.
- 808 (46) Frisch, M. J.; Pople, J. A.; Binkley, J. S. Self-consistent  
809 molecular orbital methods 2S. Supplementary functions for Gaussian  
810 basis sets. *J. Chem. Phys.* **1984**, *80*, 3265–3269.
- 811 (47) Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.;  
812 Robb, M. A.; Cheeseman, J. R.; Scalmani, G.; Barone, V.; Petersson,  
813 G. A.; Nakatsuji, H.; Li, X.; Caricato, M.; Marenich, A.; Bloino, J.;  
814 Janesko, B. G.; Gomperts, R.; Mennucci, B.; Hratchian, H. P.; Ortiz, J.  
815 V.; Izmaylov, A. F.; Sonnenberg, J. L.; Williams-Young, D.; Ding, F.;  
816 Lipparini, F.; Egidi, F.; Goings, J.; Peng, B.; Petrone, A.; Henderson,  
817 T.; Ranasinghe, D.; Zakrzewski, V. G.; Gao, J.; Rega, N.; Zheng, G.;  
818 Liang, W.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.;  
819 Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Vreven, T.;  
820 Throssell, K.; Montgomery, J. A., Jr.; Peralta, J. E.; Ogliaro, F.;  
821 Bearpark, M.; Heyd, J. J.; Brothers, E.; Kudin, K. N.; Staroverov, V.  
822 N.; Keith, T.; Kobayashi, R.; Normand, J.; Raghavachari, K.; Rendell,  
823 A.; Burant, J. C.; Iyengar, S. S.; Tomasi, J.; Cossi, M.; Millam, J. M.;  
824 Klene, M.; Adamo, C.; Cammi, R.; Ochterski, J. W.; Martin, R. L.;  
825 Morokuma, K.; Farkas, O.; Foresman, J. B.; Fox, D. J. *Gaussian 09*,  
826 Rev. A.02; Gaussian, Inc.: Wallingford, CT, 2016.
- 827 (48) Muniz-Miranda, F.; Menziani, M. C.; Pedone, A. Assessment of  
828 Exchange-Correlation Functionals in Reproducing the Structure and  
829 Optical Gap of Organic-Protected Gold Nanoclusters. *J. Phys. Chem.*  
830 *C* **2014**, *118*, 7532.
- 831 (49) Grönbeck, H.; Häkkinen, H.; Whetten, R. L. Gold-Thiolate  
832 Complexes Form a Unique  $c(4 \times 2)$  Structure on Au(111). *J. Phys.*  
833 *Chem. C* **2008**, *112*, 15940–15942.
- 834 (50) Chen, M.; Dyer, J. E.; Li, K.; Dixon, D. A. Prediction of  
835 structures and atomization energies of small silver clusters,  $(Ag)_n$ ,  $n <$   
836  $100$ . *J. Phys. Chem. A* **2013**, *117*, 8298–8313.
- 837 (51) Kroon-Batenburg, L. M. J.; Van Duijneveldt, F. B. The use of a  
838 moment-optimized DZP basis set for describing the interaction in the  
839 water dimer. *J. Mol. Struct.: THEOCHEM* **1985**, *121*, 185–199.