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Multiresolution decomposition and wavelet analysis of urban aerosol fluxes in Italy and Austria

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Abstract

Observations of turbulent aerosol fluxes are fundamental to understand basic transport processes that govern changes of particle concentrations in the atmospheric boundary layer. The turbulent surface-atmosphere exchange of atmospheric particles can be quantified using several methods, including the eddy-covariance (EC) method and spectral flux estimation methods such as wavelet analysis and multiresolution decomposition. In this work, turbulent time series obtained by EC measurements in two different cities, Lecce (Italy) and Innsbruck (Austria), are spectrally analysed applying wavelet analysis and multiresolution decomposition, and the respective turbulent spectra are compared in these two European cities to quantify the contributions to turbulent fluxes in both the time and the frequency domains. As expected, particle emission is dominant in both cities following a similar diurnal cycle. Multiresolution decomposition reveals a similar cospectral peak of particle fluxes in both cities, with a median normalized frequency of $n = 0.087$ in Lecce and $n = 0.086$ in Innsbruck. Wavelet analysis shows that the 2 – 20 s time scales contribute very strongly to the particle flux in Lecce, while in Innsbruck the 20 – 200 s time scales are clearly dominant. In both cities, larger-sized eddies contribute only sporadically to turbulent aerosol fluxes. These results suggest that spectral similarity of urban particle number fluxes holds, to a large extent, even when comparing two very different urban environments and different meteorological conditions.

39 **Keywords:** Urban aerosol fluxes; spectral analysis; cospectral peak; spectral similarity, eddy-
40 covariance.
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Highlights

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- Atmospheric particle emission is dominant in Lecce and Innsbruck
- The diurnal cycle of particle number concentrations is related with traffic activities
- Spectral similarity of urban particle number fluxes in both cities

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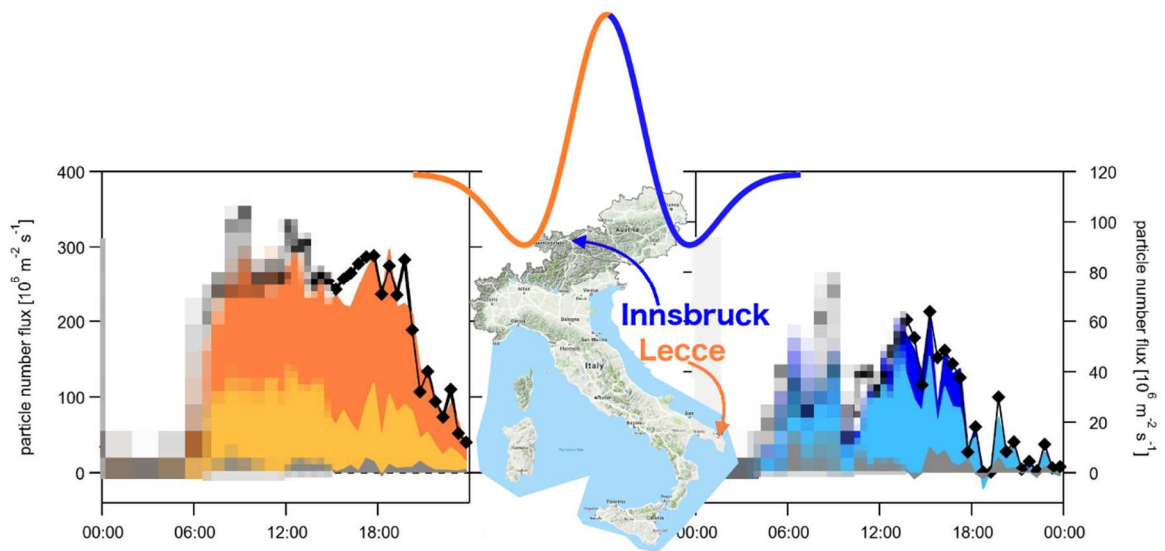
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Graphical abstract



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55 **1 Introduction**

56 Aerosols are continuously emitted from ground level sources and deposited back to the surface.
57 Aerosol surface/atmosphere exchange is dominated by turbulent transport in so-called “eddies” (e.g.
58 Reynolds, 1894; Foken, 2008). The net exchange is evaluated by measuring turbulent aerosol fluxes,
59 representing a key process governing the atmospheric aerosol burden and its impact on climate,
60 ecosystem health, and human health.

61 Our present knowledge of turbulent aerosol fluxes is incomplete, and thus, we lack a precise
62 understanding of the transport processes that govern changes in atmospheric particle concentrations.
63 The present paper aims to spectrally analyse turbulent time series and compare turbulent spectra in
64 two different European cities, Lecce (Italy) and Innsbruck (Austria) using eddy covariance (EC; e.g.
65 Aubinet et al., 2012) flux measurements. Wavelet analysis and multiresolution decomposition are
66 applied to the EC data to compare co-spectra of atmospheric turbulent time series and quantify the
67 spectral flux contributions. The importance of analysing data with wavelets lies in the time-frequency
68 localization that allows to study features of the signal both in the time and the frequency domains.

69 The turbulent surface-atmosphere exchange of atmospheric particles is usually quantified with
70 the EC method. This approach requires fast measurements of turbulent fluctuations of the vertical
71 wind speed and aerosol concentration, typically with a measurement frequency of 10 Hz, under steady
72 state conditions, i.e. negligible changes of the statistical properties of the vertical wind and aerosol
73 time series during the measurement interval (typically 30 min). Then, turbulent particle fluxes are
74 calculated as the covariance between the particle concentration and the vertical wind velocity. The
75 first studies in urban areas were performed starting from the 2000s (e.g. Nemitz et al., 2000; Dorsey
76 et al., 2002; Longley et al., 2004; Mårtensson et al., 2006; Martin et al., 2009).

77 Particle counters are typically not sufficiently fast for a measurement frequency of 10 Hz, and
78 sudden changes in particle concentration inherently violate the steady state assumption. Therefore,
79 due to limited instrument performance and non-ideal meteorological conditions, the post processing
80 of EC particle flux data must include the application of corrections, e.g. for slow response of particle
81 instrumentation, de-trending of time series, coordinate rotation of wind measurements, and rigorous
82 quality assurance/control. To date, there is no best practice procedure for measuring the surface-
83 atmosphere particle exchange, and no standard protocol for data processing and analysis. As an
84 alternative to standard EC data processing, wavelet analysis is a spectral flux estimation method that,
85 in contrast to EC, does not require the steady-state assumption. Even though spectral flux estimation
86 methods such as wavelet analysis have been used to calculate turbulent fluxes (e.g. Katul and
87 Parlange, 1995; Saito and Asanuma, 2008; van den Kroonenberg and Bange, 2007), these methods
88 have barely been used with respect to particle fluxes.

89 The present work compares the meteorological conditions and turbulent fluxes in Lecce (Italy)
90 and Innsbruck (Austria), and then analyse the source areas (footprints) of particle concentration and
91 flux measurements to estimate the contribution of different source areas to the measurement. After a
92 comparison of diurnal cycles of meteorological parameters and turbulent fluxes, multiresolution
93 decomposition is applied in order to compare normalized cospectra of particle number and buoyancy
94 fluxes and their diurnal evolution, and quantify the normalized frequency of the cospectral peak.
95 Finally, wavelet analysis is applied to extract the dominant time scales of the turbulent fluxes, and to
96 spectrally estimate 1 min particle number flux averages. The overall aim of this study is to
97 characterize the temporal dynamics of particle fluxes in two different cities with high resolution both
98 in the time and frequency domains.

99

100 **2 Sites and methods**

101 **2.1 Measurement sites and experimental setup**

102 Turbulent aerosol fluxes from two different cities, Innsbruck (Austria) and Lecce (Italy), were
103 compared. The measurements were carried out between 10 March and 24 April 2015 in Lecce, and
104 between 27 July and 21 August 2015 in Innsbruck. Both of the experimental sites, Innsbruck and
105 Lecce, were located on the rooftop of university buildings within the respective city centres.

106 Innsbruck is situated in the Alpine Inn Valley, with a population of about 130000 inhabitants.
107 The Innsbruck site (Karl et al., 2020; latitude 47°15'51.50" N; longitude 11°23'6.77" E) is located
108 south west of the city centre, and it is characterized by a mountain-valley wind system (Deventer et
109 al., 2018; von der Heyden et al., 2018). The prevailing wind direction during the measurement period
110 was ENE, and the maximum wind speed was 7.8 m s⁻¹ (Fig. 1a). The Innsbruck site is surrounded by
111 office and commercial areas, close to the Inn river and close to a busy state road (B171 "Innrain")
112 where vehicular traffic reached a maximum of 650 vehicles per hour. The measurements for this study
113 were carried out with a sonic anemometer (CSAT3, Campbell Scientific Inc., Logan, USA) and a
114 Condensation Particle Counter (CPC, Model 3772, TSI Inc., Shoreview, USA).

115 Lecce is situated in South-Eastern Italy close to the sea, with a population of about 95000
116 inhabitants. The Lecce site (latitude 40°21'22''N; longitude 18°10'02''E) is located in a street called
117 "student's street" for the high presence of schools and university buildings, and where traffic reached
118 a maximum of 2400 vehicles per hour. The main wind direction during the experimental campaign
119 was N and the maximum wind speed was 5.4 m s⁻¹ (Fig. 1b). The experimental setup for this study
120 comprised an ultrasonic anemometer (Gill R3-100, Gill Instruments, Lymington, UK) coupled with
121 a CPC (Grimm 5.403, Grimm Aerosoltechnik, Ainring, Germany).

122 The two sites exhibit several characteristic differences, including (1) the geographical and
123 topographical conditions, (2) the experimental setup, and (3) the surface properties. First, the street-
124 level elevation at the Innsbruck site is 570 m above sea level, i.e. much higher than at the Lecce site
125 at 55 m above sea level. Innsbruck is situated in the Inn valley that represents the transition zone from
126 Mediterranean to continental climate, while the Lecce site is about 12 km from the coastline of the
127 Adriatic Sea. Second, the eddy-covariance instrumentation was located on the roof of a university
128 building for the Innsbruck and Lecce sites at a height above ground level of 38.6 m and 14 m,
129 respectively. Furthermore, to minimize potential flow distortions due to the presence of high buildings
130 surrounding the sites, a wind sector between 140° and 190° was discarded from the Innsbruck data
131 (leaving 90 % of the total measurements), and a sector from 20° to 202° was eliminated from the
132 Lecce data (leaving 59 % of the total measurements). The first-order time constant of the CPC, used
133 in the Lecce site, was determined by estimating in laboratory the time response to a concentration
134 step. The results obtained in several repeated laboratory experiments was 1.3 ± 0.05 s (Conte et al.,
135 2018). The size range of the instrument is between $0.009 \mu\text{m}$ and $0.25 \mu\text{m}$. At the Innsbruck site, the
136 CPC has a nominal size range between $0.010 \mu\text{m}$ and $> 1 \mu\text{m}$, and the experimentally determined
137 response time was 0.46 s (von der Heyden et al., 2018). Third, to characterize the surface properties
138 of the urban canopies, the zero-plane displacement height (d_0) and the roughness length (z_0)
139 parameters were evaluated for both sites. These parameters were evaluated for Lecce using a fitting
140 of measured EC data under unstable conditions (Toda and Sugita, 2003). The value of d_0 was derived
141 from the standard deviation of temperature and of vertical wind speed; the value of z_0 , instead, was
142 derived from a wind profile equation including the derived d_0 value. These values were estimated to
143 be $d_0 = 7$ m and $z_0 = 1.3$ m (Conte et al., 2018). For Innsbruck, the parameters were evaluated with a
144 height-based morphometric method (Grimmond and Oke, 1999), a mean building height of $z_B = 20$
145 m, $d_0 = 0.7z_B$, and $z_0 = 0.1z_B$. The values were estimated to be about $d_0 = 14$ m and $z_0 = 2$ m (von der
146 Heyden et al., 2018).

147

148 **2.2 Source area analysis**

149 Source area (footprint) analysis is important to estimate the pollutants' relative weights at the
150 measurement point. A source area is defined as the fraction of the surface containing effective sources
151 and sinks contributing to the measurement point (e.g. Schmid and Oke, 1990; Kljun et al., 2002). In
152 this study, the source area for particle concentrations and fluxes was determined as the characteristic
153 dimension of the oval-shaped area representing 50 % of measured concentrations or fluxes. Following
154 Schmid (1994), the parameters defining the shape of the source areas were evaluated analytically by
155 two equations that represent the normalised source area dimensions, D_N . The source areas depend on

156 the measurement height (z_m), and surface-layer scaling parameters such as the surface roughness
 157 length (z_0), the Obukhov length (L), the standard deviation of lateral wind speed fluctuations (σ_v),
 158 and the friction velocity (u_*):

159

$$160 \quad D_N = \alpha_1 \left(\frac{z_m}{z_0}\right)^{\alpha_2} \exp \left[\alpha_3 \left(\frac{z_m}{L}\right)^{\alpha_4} \right] \left(\frac{\sigma_v}{u_*}\right)^{\alpha_5} \quad (\text{stable cases}) \quad (1)$$

$$161 \quad D_N = \alpha_1 \left(\frac{z_m}{z_0}\right)^{\alpha_2} \exp \left[1 - \alpha_3 \left(\frac{z_m}{L}\right)^{\alpha_4} \right] \left(\frac{\sigma_v}{u_*}\right)^{\alpha_5} \quad (\text{unstable cases}) \quad (2)$$

162

163 The parameters α_i ($i=1,5$) referring to two types of equation, are numerical inputs evaluated separately
 164 for stable (equation 1) and unstable (equation 2) conditions according to the tables given in Schmid
 165 (1994). The numerical values α_i differ for concentrations and fluxes and determine the elliptical shape
 166 of the source areas. In this work, the oval-shaped source areas were approximated by ellipses oriented
 167 along the wind direction. For every 30-min period, a function was defined, equal to one inside the
 168 elliptical domain and equal to zero outside the elliptical domain: the source area is finally the sum of
 169 all these functions normalized by the total number of measurements.

170

171 **2.3 Flux calculation and spectral analysis**

172 The EC data from Innsbruck and Lecce were studied applying wavelet analysis and multiresolution
 173 decomposition in order to investigate the temporal dynamics of particle fluxes in detail. The EC
 174 method is commonly used to calculate vertical turbulent fluxes between surfaces and the atmosphere.
 175 It requires fundamental assumptions such as steady state conditions and horizontal homogeneity (e.g.
 176 McMillen 1988; Foken & Wichura, 1996).

177 The EC method expresses a turbulent flux F_{EC} of a certain entity by the covariance between
 178 the concentration c of that entity and the vertical wind velocity w :

$$179 \quad F_{EC} = \overline{c'w'} \quad (3)$$

180

181 where the prime (') denotes the fluctuating part of that entity. Equation 3 represents the turbulent
 182 vertical transport of c by fluctuations of c and w around their mean values (Balducchi et al., 1988;
 183 Arya, 1999; Burba et al., 2012). The advantage of this method consists in the direct and in situ flux
 184 measurement without perturbation of the analysed ecosystem (Aubinet et al., 2012). An averaging
 185 period of 30 minutes is typically chosen. However, the 30-minute data obtained from EC need an
 186 accurate quality check and consistent post-processing. While the three-dimensional wind
 187 measurements using sonic anemometers are expected to be fully resolved at the sampling frequency
 188 of 10 Hz, the particle concentration measurements are dampened with experimentally determined

189 response times of 1.3 s for the CPC used in Lecce, and 0.46 s for the CPC used in Innsbruck. The
 190 expected reduction of the EC flux estimate due to imperfect sensor response can be taken into account
 191 as proposed by Horst (1997). To minimize the impact of the imperfect sensor response on the
 192 comparison of EC flux calculations and spectral analyses, all following calculations in this study are
 193 based on 1 Hz data sets. The quality of the final data is mainly influenced by sensor configuration,
 194 by meteorological conditions, and depending on the validity of the steady state assumption (Foken &
 195 Wichura, 1996; Moncrieff et al. 1996, Aubinet et al., 1996).

196 A method which does not rely on the steady state assumption to evaluate the turbulent flux is
 197 the wavelet approach (e.g. Torrence and Compo, 1998; Schaller et al., 2017). The total flux F_{Wavelet}
 198 is the integral of the cross-wavelet spectrum E_{cw} between the particle concentration c and the vertical
 199 wind velocity w , and can thus be computed as:

$$200$$

$$201 \quad F_{\text{Wavelet}} = \delta j \sum_{j=0}^J \frac{E_{cw}(j)}{s_j} \quad (4)$$

$$202$$

203 Equation (4) represents the sum over all scales s_j of the real part of the cross-wavelet spectrum E_{cw} ,
 204 and δj is the scale step size. The cross-wavelet spectrum is defined as:

$$205$$

$$206 \quad E_{cw}(j) = \frac{\delta t}{R} \frac{1}{N} \sum_{n=0}^{N-1} [W_n^c(s_j)(W_n^w(s_j))^*] \quad (5)$$

$$207$$

208 where $W_n^{c,w}(s_j)$ are the wavelet transforms of the time series of c and w , respectively, and the asterisk
 209 (*) indicates the complex conjugate of the wavelet transform, δt is the time step of the time series, R
 210 is a reconstruction factor specific to the applied wavelet (e.g. $R = 3.541$ for the Mexican Hat wavelet,
 211 cf. Torrence and Compo, 1998), and N is the number of data points of the time series in the averaging
 212 period. It is useful to recall that the wavelet transform of a generic discrete time series $x_n(t), n =$
 213 $0, \dots, N - 1$ is generally defined as the convolution of the series with a scaled “s” and time-shifted
 214 wavelet function $\psi(t)$:

$$215 \quad W_n(s) = \sum_{n'=0}^{N-1} x_{n'}(t) \times \Psi^* \left(\frac{(n - n')\delta t}{s_j} \right) \quad (6),$$

216 allowing to obtain information about variability both in time and frequency of the time series. For
 217 further details see Torrence and Compo (1998) and Schaller et al. (2017).

218 The scales “ s_j ” entering in the equations (4) - (6) are typically conveniently chosen as fractional
 219 powers of two:

220
$$s_j = s_0 2^{j\delta j} \quad (7)$$

221 with $j = 0$ to the maximum number of scales J , the smallest resolvable scale $s_0 = 2\delta t$, and the scale
 222 step size δj , in this study set to 0.25. The wavelet function is chosen according to the characteristics
 223 of the time series. In this study, the Mexican Hat wavelet was used to obtain the wavelet transforms
 224 of the two time series, $W_n^2(s)$, and information about the contribution of different frequency and time
 225 domains to the cross-wavelet spectrum. Thus, it is possible to evaluate turbulent fluxes over short
 226 averaging intervals, e.g. 1 min flux estimates.

227
 228 The data collected during both experimental campaigns were processed also using the
 229 multiresolution flux decomposition method (MRD). This analysis allows to decompose a time
 230 series $x_n(t)$ into averages on different time scales, obtained by M subsequent divisions of the whole
 231 time series starting from coarse (the whole series) to finer divisions in 2^M averaging points. This
 232 procedure of averaging using different window lengths corresponds to viewing the data at different
 233 resolutions, with each mode of the multiresolution decomposition corresponding to a certain moving
 234 average with a given length (Howell and Mahrt 1995):

235
 236
$$\bar{x}_n(2^m) = \frac{1}{2^m} \sum_{i=(n-1)2^{m+1}}^{n2^m} x_i \quad (8)$$

237
 238 where $n=1,2,\dots,2^{M-m}$ is the window and x_i denotes the “residual” of the time series, that is the
 239 difference between the series and its approximation of order n . Once the temporal series is
 240 transformed by MRD in a frequency series, considering again the concentration c and the vertical
 241 wind velocity w , it is possible to evaluate the turbulent flux F_{MRD} :

242
 243
$$F_{MRD} = \sum_{m=1}^M \frac{1}{2^{M-m}} \sum_{n=1}^{2^{M-m}} (2^{m-1} \overline{c_{2n}} - 2^m \overline{c_n}) (2^{m-1} \overline{w_{2n}} - 2^m \overline{w_n}) \quad (9)$$

244
 245
 246 The three methods, EC, wavelet and MRD, allow to evaluate turbulent fluxes of a certain quantity.
 247 The EC method allows to investigate a turbulent flux on the ecosystem scale but it does not allow to
 248 obtain information about short-time turbulent events due to a violation of the steady state
 249 assumptions, on which this method is based. To overcome the steady-state requirement of EC, the
 250 wavelet analysis provides a localization, and thus a detailed analysis, of the turbulent events in both
 251 time and frequency. Wavelet analysis can provide the variability of a specific turbulent event and
 252 how its variability changes in time (Torrence and Compo 1998). MRD can be defined as non-

253 overlapping spectral analysis that provides information about the time scales of atmospheric
254 fluctuations of a certain scalar. MRD typically has a smaller resolution in the frequency domain than
255 wavelet analysis. The MRD locally decomposes the turbulent flux trying to characterize the dominant
256 local events that correspond to the different averaging lengths considered in the analysis.

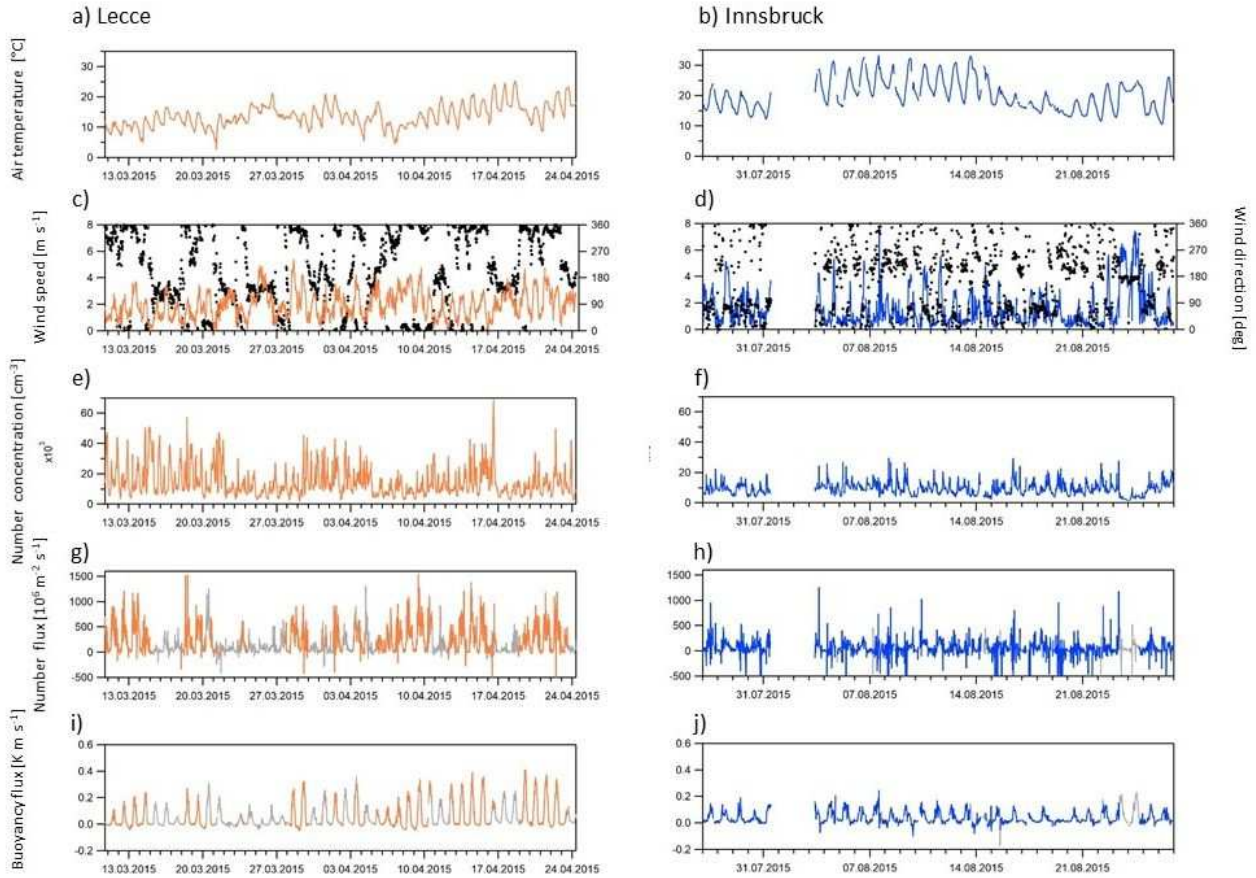
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258 **3 Results**

259 **3.1 Overview of campaigns in Lecce and Innsbruck**

260 **General overview**

261 Figure 1 and Table 1 summarize the time series, median values, and 5 % and 95 % percentiles of
262 meteorological data and particle concentrations and fluxes in the Lecce and Innsbruck measurement
263 campaigns. Due to the earlier time of year, the range of the air temperature (Fig. 1) and the median
264 air temperature is lower during the Lecce campaign in March and April 2015 than during the
265 Innsbruck campaign in July and August 2015. In both locations, wind speed exhibits a clear diurnal
266 cycle with calm nights and higher wind speeds during the day. The median wind speed of 1.8 m s^{-1} is
267 slightly larger in Lecce than in Innsbruck (median wind speed 1.1 m s^{-1}), however, a southerly Foehn
268 event in Innsbruck on 23 and 24 August is the reason for higher maximum wind speeds in Innsbruck.
269 The prevailing wind direction during the Lecce campaign is from the North, while the Innsbruck
270 campaign is characterized by a pronounced mountain-valley wind system with wind from the ENE
271 during the day and from the SW at night. This regular change of wind direction in the Inn valley
272 typically occurs between 09:30 and 11:00 in the morning, and between 17:30 and 19:00 in the
273 evening. While the baseline particle number concentration (5 % percentile) of roughly 4000 cm^{-3} is
274 similar in Lecce and Innsbruck, the median and the 95 % percentile values of particle number
275 concentration are considerably larger in Lecce (median: 11419 cm^{-3}) than in Innsbruck (median: 8990
276 cm^{-3}). Similarly, the particle number fluxes are larger in Lecce than in Innsbruck, and the fraction of
277 positive particle number fluxes, i.e. net particle emission, is considerably larger in Lecce (90.6 %)
278 than in Innsbruck (69.7 %) as well. The buoyancy flux reaches larger maximum values in Lecce than
279 in Innsbruck. The stability conditions are mostly near-neutral ($|z_e/L| < 0.1$, where $z_e=z-d_0$ is the
280 effective measurement height) or slightly unstable in Lecce, while unstable conditions ($z_e/L < -0.1$)
281 prevail in Innsbruck.



282

283 Fig. 1: Time series of a,b) air temperature, c,d) wind speed (continuous line) and wind direction (dots), e,f)
 284 particle number concentration, g,h) particle number flux, i,j) and buoyancy flux of the measurement campaigns
 285 in Lecce and Innsbruck; grey flux data are from discarded wind sectors.

286

287

288 Tab. 1: Median values and 5 % and 95 % percentiles of air temperature, wind speed, particle number
 289 concentration, particle number flux, friction velocity u^* , buoyancy flux, and stability parameter z_e/L (Obukhov

290 length $L = -\frac{u^{*2}}{k(\frac{g}{T_0})T^*}$, where k is von Kármán constant, g/T_0 is a buoyancy parameter and T^* is the friction

291

temperature) of the measurement campaigns in Lecce and Innsbruck.

	Lecce			Innsbruck		
	5 %	median	95 %	5 %	median	95 %
air temperature [°C]	8.3	13.5	20.5	13.3	19.6	29.8
wind speed [m s ⁻¹]	0.6	1.8	3.8	0.3	1.1	4.6
particle number concentration [cm ⁻³]	4336	11419	34983	3826	8990	17151
particle number flux [10 ⁶ m ⁻² s ⁻¹]	-14.9	210.8	789.8	-209.4	40.2	307.2
u^* [m s ⁻¹]	0.12	0.44	0.84	0.10	0.24	0.52
buoyancy flux [K m s ⁻¹]	-0.02	0.01	0.28	-0.01	0.02	0.12
z_e/L [-]	-0.2	0.0	0.1	-7.0	-0.7	0.2

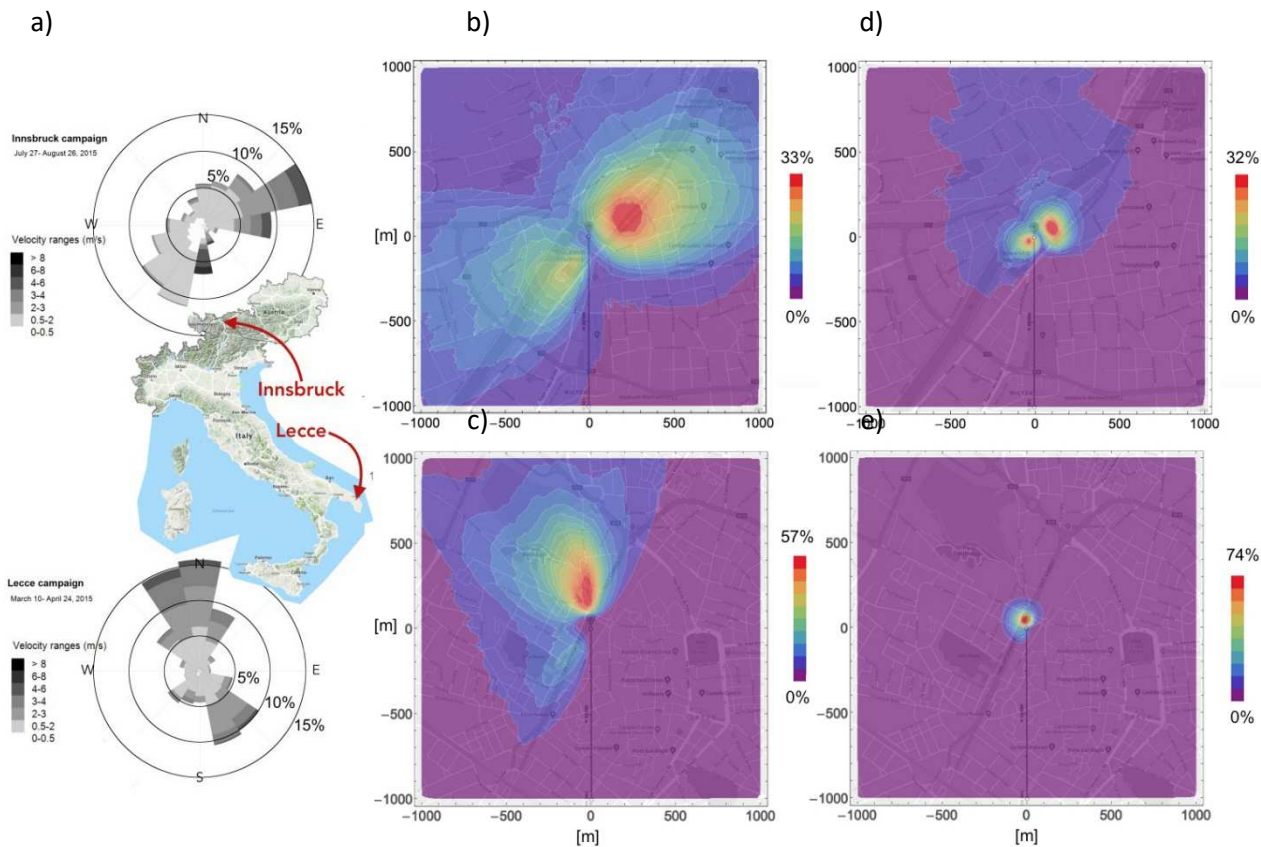
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295 Source area estimates

296



297

298 Figure 2: a) Location and wind roses of the measurement sites in Innsbruck and Lecce, as well as average
299 concentration source areas (b, c) and average flux source areas (d, e). The color scale represents the percentage
300 of measures that gives contribution to the source area in any point.

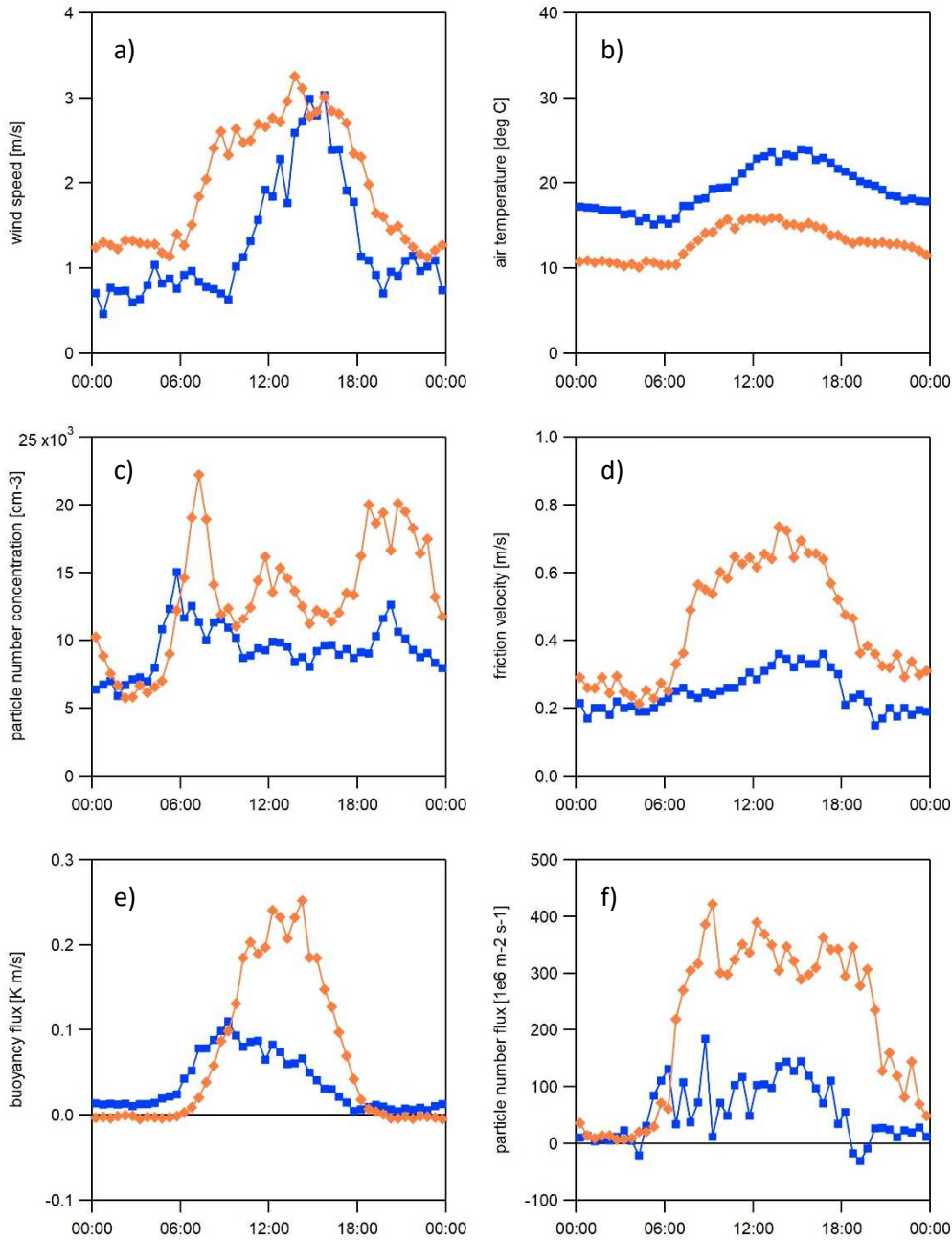
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302 The source area estimates calculated using Eqs. 1 and 2 (section 2.2) are reported in Figure
303 2b and 2c for concentrations, and in Figure 2d and 2e for fluxes, for Innsbruck and Lecce,
304 respectively. In Innsbruck, the dominant concentration source areas are to the N/NE with the densely
305 built old town including a state road and public transportation to the city centre, and to a lesser extent
306 to the SW with University buildings and residential areas. The Innsbruck flux footprint is mostly
307 affected by the nearby University buildings in the S/SW sector. In Lecce, the dominant concentration
308 and flux source areas are to the North of the measurement site with many buildings and roads with a
309 high traffic density. The colour scale ranges from red, corresponding to the most influential areas, to
310 blue for the areas minimally influencing the measurements. Normalized over all measurement
311 periods, the most influential areas with respect to concentrations (red colours in Fig. 2b,c) are about
312 24 % and 55 % of the average source area in Innsbruck and Lecce, respectively, while the areas of

313 minimum influence (blue colours) are about 2 % and 5 % of the average source area. For fluxes (Fig.
314 2d,e), the most influential areas are 10 % of the average flux source area in Innsbruck, and 12 % in
315 Lecce, while the minimum influential areas are about 1 % of the average flux source area in both
316 cities. In general, source areas for fluxes are smaller than those for concentrations (Contini et al.,
317 2012; Vesala et al., 2008; Conte et al., 2018) because their values are proportional to the difference
318 between particles crossing the measurement level in upward and downward direction (Vesala et al.,
319 2008), whereas particles crossing the measurement point always contribute positively to the
320 evaluation of concentration source areas. Furthermore, flux measurements consider only the turbulent
321 transport of a scalar, while concentration measurements consider also transport by the mean flow.

322

323 **Diurnal cycles**



324

325 Fig. 3: Median diurnal cycles of (a) wind speed, (b) air temperature, (c) particle number concentration, (d)
 326 friction velocity, (e) buoyancy flux, and (f) particle number flux in Lecce (orange) and in Innsbruck (blue);
 327 EC data after “bad” wind sectors removed.

328

329 Figure 3 shows median diurnal cycles of wind speed, air temperature, particle number concentration,
 330 friction velocity, buoyancy flux, and particle number flux obtained by EC method in Lecce (orange)
 331 and Innsbruck (blue). As mentioned before, there is a clear diurnal cycle of wind speed (Fig. 3a) with
 332 calm nights and higher wind speeds during the day, from about 06:00 to 21:00 in Lecce and from
 333 about 09:00 to 20:00 in Innsbruck. The peak wind speed of about 3 m s^{-1} is reached in the afternoon
 334 between 13:00 and 16:00 at both sites. The diurnal cycle of air temperature (Fig. 3b) is very similar
 335 but shifted to higher temperatures in Innsbruck during the summer season. Temperatures begin to rise

336 at 07:00 in the morning at both sites. Particle number concentrations (Fig. 3c) show several peaks
337 during the course of the day: First, at both sites there is a steep increase in particle number
338 concentration from 04:30 in the morning reaching the peak value of 15000 cm^{-3} in Innsbruck at 06:00,
339 and 22000 cm^{-3} at 07:00 in Lecce. A second, slightly lower peak can be observed at both sites in the
340 evening hours around 21:00. Only in Lecce, a clear third peak is found around noontime in
341 correspondence of the maximum in solar radiation. Since the central peak happens around midday at
342 noon, in correspondence of the maximum in solar radiation, it was possible that this increase in
343 concentrations was related to local nucleation. The absence of this central peak in the suburban
344 environment excluded a regional nucleation process pointing to an urban nucleation event. The
345 presence of this peak in the Lecce data set has already been discussed by Conte et al. (2015) and
346 Conte et al. (2018). This feature cannot be found in the Innsbruck data.

347 The friction velocity (Fig. 3d) is clearly enhanced during daytime both in Lecce and in
348 Innsbruck, which indicates well-developed turbulent conditions. The median friction velocity in
349 Lecce is between 0.6 and 0.8 m s^{-1} from 08:00 to 18:00, while the peak median friction velocity in
350 Innsbruck is much lower, around 0.3 m s^{-1} from 11:00 to 18:00. The buoyancy flux (Fig. 3e) is mostly
351 positive and exhibits a clear diurnal cycle with values close to zero during the night, and an increase
352 after sunrise, which occurs earlier in the Innsbruck summer data (July/August) compared to the Lecce
353 data from March/April. Again, the peak median buoyancy flux in Lecce is more than double
354 compared to the Innsbruck peak value. Both in Lecce and Innsbruck, the median buoyancy flux drops
355 back to values close to zero in the evening around 18:00. The particle number fluxes (Fig. 3f) are
356 mostly positive both in Lecce and in Innsbruck, indicating net particle emission especially during the
357 day. From 00:00 to 04:00 the net particle number fluxes are close to zero, and start to increase at the
358 same time when particle number concentrations increase. During daytime, median particle number
359 fluxes in Lecce are between 300 and $400 \cdot 10^6 \text{ m}^{-2} \text{ s}^{-1}$, while the median particle number fluxes in
360 Innsbruck are between 50 and $150 \cdot 10^6 \text{ m}^{-2} \text{ s}^{-1}$. Interestingly, in Innsbruck there is a brief period of net
361 particle deposition in the evening from 18:30 to 20:00, which cannot be observed in Lecce.

362 Adding data from the “bad” wind sectors does not considerably change the diurnal cycles in
363 Innsbruck with respect to wind speed, air temperature, particle number concentration, and friction
364 velocity but it slightly decreases the median estimates of the buoyancy flux and the particle number
365 flux in Lecce. When taking into account only near-neutral and unstable conditions ($z_e/L < 0.1$), there
366 is no considerable change of the diurnal cycles of wind speed, air temperature, particle number
367 concentration, friction velocity and buoyancy fluxes both in Lecce and in Innsbruck. However,
368 particle number fluxes are slightly enhanced in Lecce compared to Fig. 3.

369

370

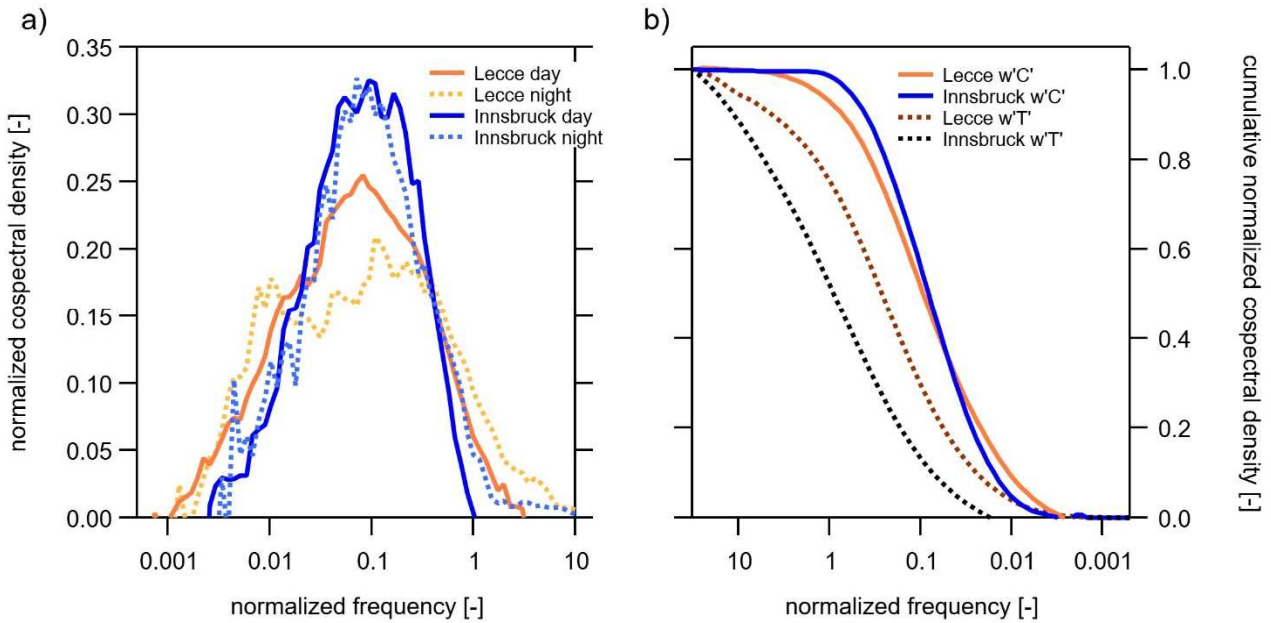
371 3.2 Spectral analysis by multiresolution decomposition

372 We use multiresolution decomposition (MRD) as a simple tool to calculate cospectra of the vertical
373 wind speed w , temperature T , and particle number concentration C , respectively. The integrals of the
374 MRD cospectra, as given by Eq. 8, represent the total covariances $w'T'$ and $w'C'$, i.e. the buoyancy
375 and particle number fluxes. Due to the fact that the typical length scales of eddies increase with
376 increasing height in the surface layer, it is common practice to investigate cospectra as a function of
377 the normalized frequency $n = f z_e/u$, where f is the frequency [s^{-1}], z_e is the effective measurement
378 height [m], and u is the horizontal wind speed [$m s^{-1}$] (e.g. Foken, 2008). To compare the contributions
379 of different frequencies to the total flux, individual cospectra are normalized by the corresponding
380 flux to yield the unit area normalized cospectral density. The normalized cospectra indicate the
381 contributions of different frequencies to the buoyancy and particle number fluxes.

382

383 Comparing MRD cospectra in Lecce and Innsbruck

384



385

386 Fig. 4: a) Median normalized cospectra of particle number fluxes in Lecce (orange) and Innsbruck (blue)
387 during daytime (11:00-17:00, solid) and night-time (23:00 – 05:00, dotted) conditions; b) cumulative median
388 normalized cospectra of particle number fluxes (solid) and buoyancy fluxes (dotted) in Lecce and Innsbruck.
389

390 Figure 4a compares the median normalized cospectra of particle number fluxes measured in Lecce
391 and Innsbruck during daytime conditions from 11:00 to 17:00, and during the night from 23:00 to
392 05:00. The location of the cospectral peak n_m is close to a value of $n_m = 0.085$, which is expected

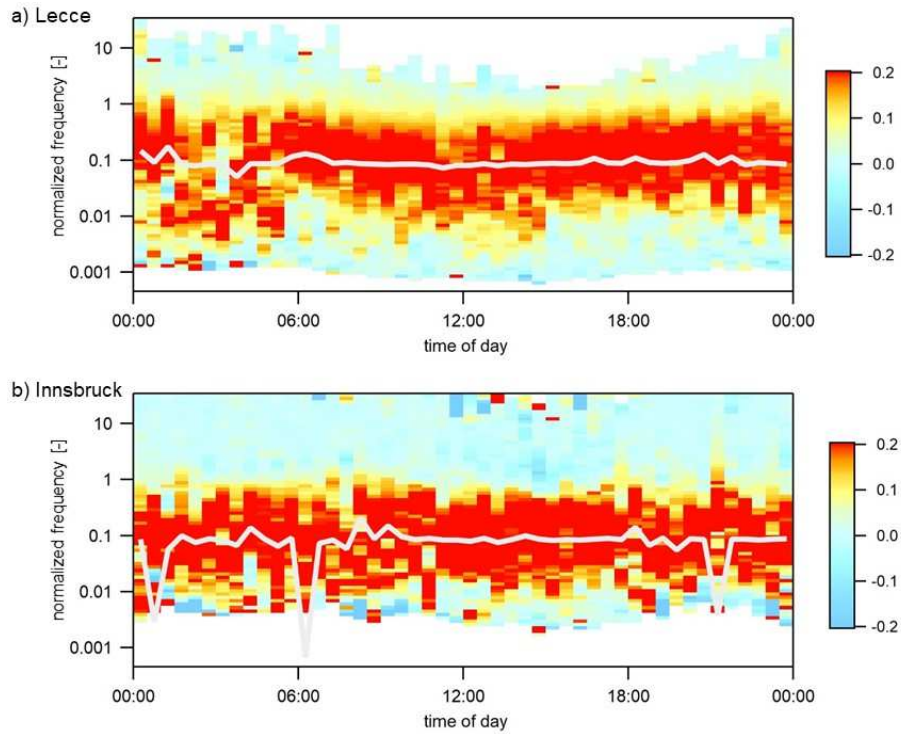
393 under neutral and unstable conditions (Kaimal et al., 1972). In Lecce, the median cospectrum at night
394 is broader than during the day, with equal flux contributions in the normalized frequency range
395 between 0.01 and 1. In Innsbruck, the cospectrum is narrower compared to Lecce, and very similar
396 during daytime and night-time, indicating a narrower distribution of eddy sizes compared to Lecce.
397 This indicates that the typical eddy size associated with the normalized frequency of the cospectral
398 peak is more dominant in Innsbruck compared to Lecce, possibly due to a larger heterogeneity of the
399 urban canopy in the flux footprint of Lecce compared to Innsbruck.

400 When comparing median cumulative cospectra in Lecce and Innsbruck (Fig. 4b), high frequency
401 contributions are evidently higher for buoyancy fluxes than for particle number fluxes. There are no
402 considerable contributions to particle number fluxes at normalized frequencies greater than 1,
403 especially in Innsbruck. This illustrates the imperfect response of the particle counters to
404 concentration fluctuations, while the sonic anemometer measurement fully resolves the acoustic
405 temperature fluctuations at the measurement frequency. In contrast, low frequency contributions to
406 buoyancy fluxes are much smaller than to particle number fluxes, indicating that scalar similarity of
407 temperature and particle number concentration is limited. In summary, the particle flux cospectra are
408 shifted to lower frequencies compared to the buoyancy flux cospectra, and the particle and buoyancy
409 flux cospectra are more similar in Lecce than in Innsbruck. The same observations are found when
410 separating daytime and night-time data.

411

412

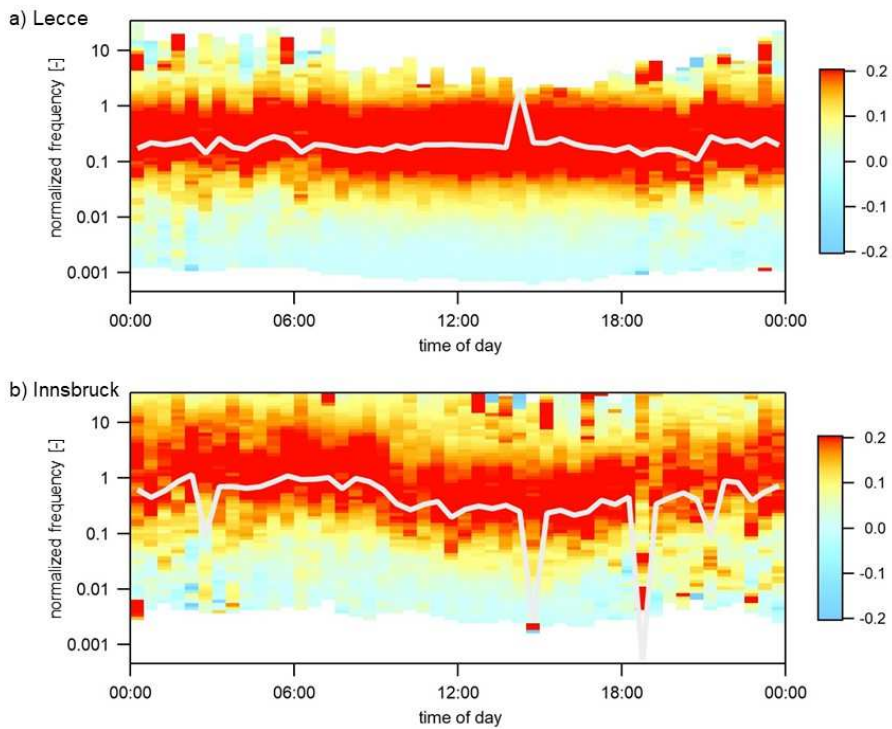
413 **Diurnal evolution of MRD cospectra**



414

415 Figure 5: Diurnal cycle of median normalized particle number flux cospectra in a) Lecce and b) Innsbruck;

416 the white line indicates the cospectral peak fitted to the median cospectra of all 30 min intervals.



417

418 Figure 6: Diurnal cycle of median normalized buoyancy flux cospectra in a) Lecce and b) Innsbruck; the white

419 line indicates the cospectral peak fitted to the median cospectra of all 30 min intervals.

420

421 Figures 5 and 6 show the diurnal evolution of the median normalized cospectra with respect to particle
 422 number fluxes and buoyancy fluxes in Lecce and Innsbruck obtained from multiresolution
 423 decomposition. The white line indicates the cospectral peak at the normalized frequency n_m fitted to
 424 the median cospectra of all 30 min intervals assuming a model for the normalized cospectrum C_0
 425 according to Eq. 10 (e.g. Horst, 1997)

$$426 \quad C_0 = \frac{2}{\pi} \frac{n/n_m}{1+(n/n_m)^2} \quad (10)$$

427 Thus, the cospectral peak indicates the normalized frequency n with the highest contribution to the
 428 flux, i.e. the inverse of the dominant time scale of turbulent transport. The red colours indicate the
 429 normalized frequency ranges with the highest contribution to the total fluxes at different times of the
 430 day.

431 For particle number fluxes (Fig. 5), the main contributions are in the normalized frequency
 432 range from 0.01 to 1, both in Lecce and in Innsbruck. In Lecce (Fig. 5a), there is a lot of scatter from
 433 00:00 to 06:00, when fluxes are generally small. In the morning after 06:00, there seems to be a slight
 434 shift from higher frequencies to lower frequencies around noon, and then again a slight increase to
 435 higher frequencies contributing to the particle number flux in the evening. In contrast, there is no
 436 obvious diurnal pattern in the median normalized cospectra of buoyancy fluxes (Fig. 6a), and the
 437 contributing frequencies are clearly shifted to higher frequencies from 0.1 to 1. In Innsbruck, there is
 438 no obvious diurnal pattern in the median normalized cospectra of particle number fluxes (Fig. 5b) but
 439 a shift from higher frequency contributions during night-time to lower frequency contributions to the
 440 buoyancy flux during the day (Fig. 6b). This indicates that the dominant time scales of turbulent
 441 transport are different for the two studied scalars, particle number concentration and temperature. The
 442 peak frequencies of the particle flux cospectra are similar in Lecce (median of $n_m = 0.087$) and in
 443 Innsbruck (median of $n_m = 0.086$), close to the expected value of $n_m = 0.085$, and relatively constant
 444 throughout the day (Lecce interquartile range 0.085-0.091; Innsbruck interquartile range 0.079-
 445 0.089). In contrast, the buoyancy flux cospectra show clearly higher peak frequencies (Lecce median
 446 of $n_m = 0.193$; Innsbruck median of $n_m = 0.445$) compared to the particle flux cospectra. The higher
 447 median value in Innsbruck is mostly due to much higher peak frequencies during the night-time. In
 448 the well-mixed convective boundary layer, the Eulerian integral scale Λ_i associated with the scalar
 449 flux is related to the normalized frequency of the cospectral peak, n_m , and the effective measurement
 450 height z_e according to Eq. 11 (Mann and Lenschow, 1994):

$$451 \quad \Lambda_i = \frac{z_e}{2\pi n_m} \quad (11)$$

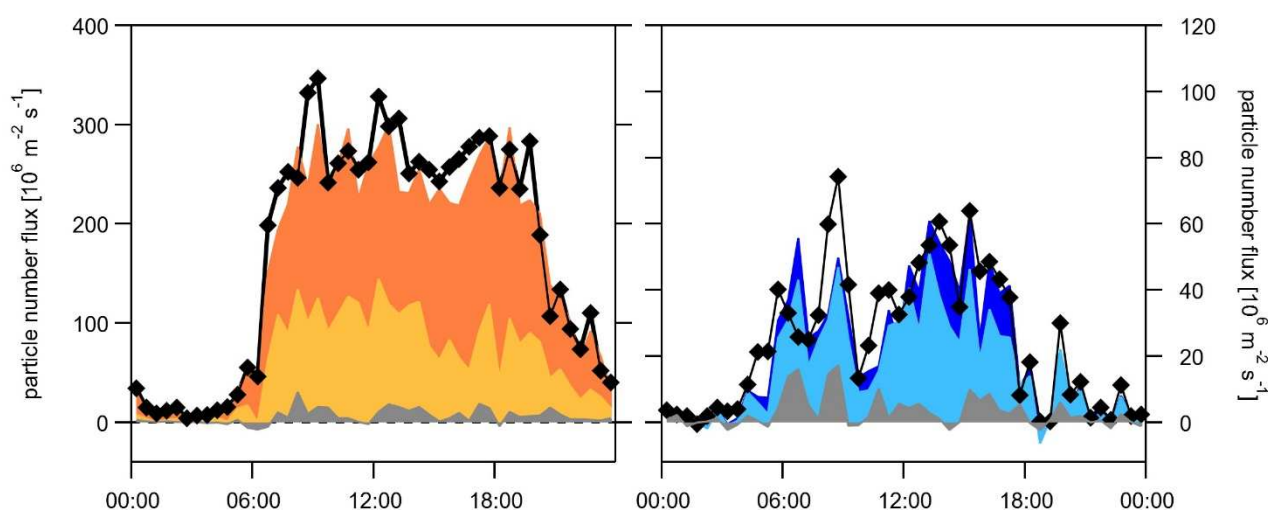
452 For the particle fluxes, this length scale is relatively stable with a median value of $\Lambda_i = 13$ m in Lecce
 453 and $\Lambda_i = 45$ m in Innsbruck. In contrast, the median Eulerian integral scale of the buoyancy fluxes is

454 $\Lambda_i = 6$ m in Lecce, and exhibits a diurnal cycle in Innsbruck with values around $\Lambda_i = 6$ m in the night-
 455 time and Λ_i up to 20 m during daytime. It should be noted that the effective measurement height in
 456 Innsbruck is $z_e = 24.6$ m, and $z_e = 7$ m in Lecce. Thus, both the ratio of the Eulerian integral scales of
 457 particle flux cospectra and the ratio of the effective measurement heights in Innsbruck and Lecce are
 458 approximately 3.5, and the different values of the Eulerian integral scale of particle flux cospectra are
 459 proportional to the different effective measurement heights above the two urban canopies.

460

461 3.3 Wavelet analysis

462 Wavelet analysis yields cross-wavelet spectra of particle number concentration and wind speed with
 463 information about the contributions of different frequency ranges of the turbulent time series to the
 464 total fluxes in both the time and frequency domains. This allows to analyse the typical contributions
 465 of different time scales, i.e. eddy sizes, in more detail.



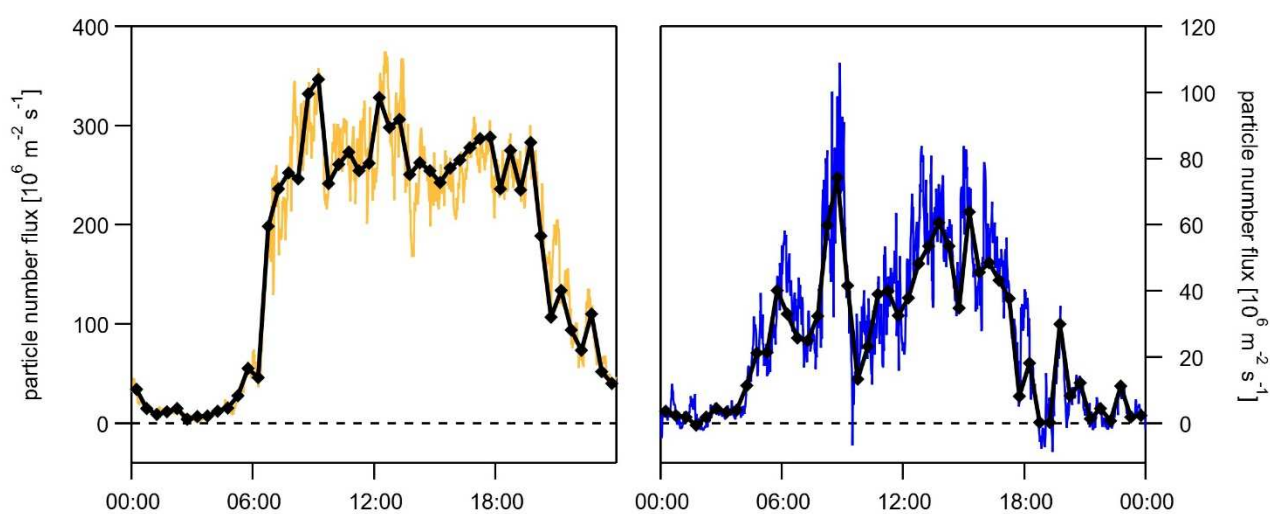
466

467 Figure 7: Median diurnal cycle of particle number fluxes reconstructed from wavelet analysis for 30 min
 468 intervals (black) in Lecce (left) and Innsbruck (right); flux contributions of 2 – 20 s scales in orange and dark
 469 blue, 20 – 200 s scales in yellow and light blue, and 200 – 2000 s scales in grey.

470 Figure 7 shows the median diurnal cycle of 30 min particle number fluxes in Lecce and Innsbruck
 471 reconstructed from wavelet analysis, and the contributions of three different frequency ranges
 472 corresponding to time scales of 2 – 20 s (small eddy sizes), 20 – 200 s (medium eddy sizes), and 200
 473 – 2000 s (larger eddy sizes). Note that the contributions of scales can be positive or negative, and
 474 therefore, the sum of the three frequency ranges may be slightly larger or smaller at any time interval.
 475 Interestingly, the 2 – 20 s time scales (orange) contribute very strongly to the particle flux in Lecce,
 476 while in Innsbruck the 20 – 200 s time scales (light blue) are clearly dominant. This is consistent with

477 the smaller Eulerian integral scale of the particle fluxes in Lecce compared to Innsbruck (cf. section
 478 3.2). Both in Lecce and Innsbruck, the contribution of the 200 – 2000 s time scales to the total median
 479 particle flux is very small, i.e. less than $20 \cdot 10^6 \text{ m}^{-2} \text{ s}^{-1}$. However, in many cases on individual days,
 480 quick changes in fluxes are due to larger-than-usual contributions of the 200 – 2000 s time scales.
 481 Thus, sporadically, larger eddies can contribute significantly to turbulent fluxes and may lead to
 482 sudden changes of the particle number concentration.

483 Taking advantage of the localization in the time domain, an estimate of the particle number flux can
 484 be calculated from cross-wavelet spectra not only for 30 min intervals (like in EC) but also for shorter
 485 time intervals, e.g. 1 min.



486

487 Figure 8: Median diurnal cycle of particle number fluxes reconstructed from wavelet analysis for 30 min
 488 intervals (black) and 1 min intervals (smoothed with a 30 min rolling average filter) in Lecce (left, orange) and
 489 Innsbruck (right, blue).

490 Figure 8 shows the median diurnal cycle of particle number fluxes in Lecce and Innsbruck,
 491 reconstructed from wavelet analysis based on 30 min intervals and 1 min intervals. Apparently, the 1
 492 min flux estimates are scattered around the 30 min flux estimates with deviations of about 30 %
 493 during daytime. In Lecce, the deviations are quite small during the night and constant during the day.
 494 In Innsbruck, the strongest deviations occur in the transition periods of wind direction reversal, i.e.
 495 around 09:00 in the morning, and between 17:30 and 19:00 in the evening. During periods when the
 496 1 min flux estimates are close to the 30 min flux estimates, steady state conditions are expected.
 497 However, this interpretation should not be applied to median diurnal cycles of the 1 min and 30 min
 498 wavelet flux estimates but to individual time periods of the Lecce and Innsbruck data sets.

499 **Conclusions**

500 On average, particle emission is dominant in both cities. The particle emission fluxes follow a similar
501 diurnal cycle but are much stronger in Lecce compared to Innsbruck. The diurnal cycle of the particle
502 number concentration shows three maxima in Lecce: one in the morning, one at noon and one in the
503 late afternoon. In Innsbruck, only the morning and late afternoon maxima are visible. These two peaks
504 are probably associated with high traffic density in the morning and evening rush hours, and the
505 noontime peak in Lecce is likely due to local particle nucleation. The source areas for fluxes were
506 estimated to be much smaller than those for concentrations. The fact that the flux source areas include
507 the respective main roads located near the experimental sites, and that the diurnal cycle of particle
508 emission fluxes is consistent with vehicular traffic activity, traffic is identified as a major emission
509 source of particles both in Lecce and in Innsbruck.

510 Multiresolution decomposition allowed to compare the normalized cospectra of vertical wind speed
511 and particle number concentration, and thus, the contribution of different frequencies to the total
512 particle number flux. On the one hand, the shape of the median normalized cospectra of particle fluxes
513 was slightly different in Lecce and Innsbruck, with a broader distribution of frequencies contributing
514 to the flux in Lecce compared to Innsbruck, especially at night. On the other hand, the cospectral
515 peak, i.e. the frequency with the largest contribution to the particle number flux, was similar in Lecce
516 and in Innsbruck. The cospectral peak of particle fluxes was relatively constant throughout the day
517 both in Lecce and Innsbruck, with a median value of 0.087 in Lecce and 0.086 in Innsbruck. The
518 cospectral maximum is expected at $n_m = 0.085$ under neutral conditions. These findings support the
519 assumption of spectral similarity for cospectra of urban particle number fluxes in two cities with
520 different micrometeorological conditions. At the same time, scalar similarity with respect to
521 temperature and particle number concentration is limited. However, when comparing cospectra of
522 particle number fluxes and buoyancy fluxes, the scalars are more similar in Lecce than in Innsbruck;
523 in both cities, the cospectral peak of the buoyancy fluxes is shifted to higher frequencies, and the
524 contribution of high frequencies is larger to the buoyancy flux than to the particle flux. This may
525 indicate that high frequency particle flux contributions were dampened due to insufficient time
526 resolution of the condensation particle counters, even when using 1 Hz data in this study.

527 Considering the results of the wavelet analysis, full resolution of the high frequency
528 contributions may be particularly important in Lecce, where the 2 – 20 s time scales contribute very
529 strongly to the particle flux in Lecce, while in Innsbruck the 20 – 200 s time scales are clearly
530 dominant. Both in Lecce and Innsbruck, the contribution of the 200 – 2000 s time scales to the total
531 median particle flux is very small. Thus, in order to quantify particle emission fluxes in these two
532 cities, the averaging interval used for the EC calculations may be changed to shorter periods of 5 or

533 10 min, and steady-state assumptions may not be required for full 30 min intervals. However, larger-
534 sized eddies might sporadically contribute in a very strong way to particle number fluxes. Even if the
535 steady state assumption is not fulfilled, reconstruction of flux estimates from wavelet analysis can
536 still be applied to shorter time intervals. For the studied periods in Lecce and Innsbruck, the 1 min
537 wavelet estimates of particle number fluxes deviated by less than 30 % from the 30 min EC flux
538 estimates on average during daytime.

539 In summary, the presented results clearly suggest that spectral similarity of urban particle
540 number fluxes holds to a large extent, even when comparing two very different urban environments
541 such as Lecce close to the Adriatic Sea and Innsbruck in the Alpine Inn valley. The results also
542 illustrate the great potential of spectral flux estimation methods such as wavelet analysis for
543 application in future studies of turbulent particle fluxes.

544

545 **Acknowledgements**

546 The authors thank the University of Salento (in Lecce) for hosting the instruments in its premise. A
547 special thanks to Dr. Antonio Donateo and Dr. F.M. Grasso (ISAC-CNR) for their help in the setup
548 of the measurement campaign and in data collection in Lecce. The authors are grateful for the
549 hospitality and support provided by T. Karl and M. Graus (University of Innsbruck), and the help by
550 L. von der Heyden and M. Deventer in carrying out the particle number flux measurements during
551 the Innsbruck Air Quality Study 2015. Part of this study was financially supported by the TU Berlin
552 Office of International Affairs.

553

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Graphical abstract

