

Impact of chirality on the aggregation modes of L-phenylalanine- and D-glucose-decorated phenylene–thiophene oligomers†

Omar Hassan Omar,^{*a} Marta Falcone,^{‡b} Alessandra Operamolla^{§b} and Gianluigi Albano^{§*b}

The aggregation modes of three L-phenylalanine- or D-glucose-functionalized phenylene–thiophene oligomers were investigated using UV-Vis absorption and electronic circular dichroism (ECD) spectroscopies in both solution aggregation and thin film states. The ECD measurements revealed superior capability to provide information about their first level of supramolecular organization and to detect the possible co-existence of multiple aggregation pathways, very difficult to identify by using only UV-Vis absorption. The impact on the (chiro)optical response of the nature of the chiral pendant group, especially in terms of intermolecular hydrogen bonding and steric hindrance, as well as that of the oligothiophene π -conjugated length, will be elucidated. The ability to recognize and control different aggregation modes could be very useful for the preparation of thin films of π -conjugated oligomers with highly tunable chiroptical features in view of optical sensing and other innovative opto-electronic applications.

Received 30th April 2021,
Accepted 31st May 2021

DOI: 10.1039/d1nj02125g

rsc.li/njc

Introduction

Organic π -conjugated materials play a prominent role in optoelectronic and nanotechnology applications, such as an active layer^{1,2} in organic light-emitting diodes (OLEDs),^{3,4} organic photovoltaic (OPV) cells,⁵ organic field-effect transistors (OFETs)⁶ and sensors.⁷ The intrinsic properties of organic π -conjugated systems (such as light absorption, emission, related energy gap, *etc.*) can be finely modulated for the desired application not only by chemical structure modifications, but also by controlling their organization in the solid state at different levels of hierarchy, from the first-order supramolecular arrangement^{8–10} to the nano/mesoscale.^{11,12} A crucial aspect is the possibility of different local domains; the same molecular system may undergo competing aggregation pathways during film fabrication, leading to polymorphic aggregated phases (kinetic *vs.* thermodynamic states), each featuring different

optoelectronic properties.^{13,14} The ability to recognize and control these different aggregation modes, especially in the case of thin films, would be useful for obtaining layers of chiral π -conjugated systems with highly tunable properties for innovative optoelectronic applications.¹⁵

Chirality also is a further valuable tool for controlling the structural organization of π -conjugated molecules at all hierarchical levels;¹⁵ stereodefined chiral elements in the molecular structure can drive their self-organization towards chiral supramolecular architectures, such as self-assembled helices, spirals or chiral sheets;^{16–19} this first-order supramolecular chirality could be transferred to larger scale chiral morphologies, generating twisted ribbons, helical fibers and other supramolecular aggregation patterns.^{17,20} Chiral functionalities tend to favour regularly twisted π -stacked arrangements, affecting in a reproducible way the fundamental electro-optical properties such as exciton migration and light emission, which in some cases may be hampered by too intense co-facial interactions.^{16,21} In this regard, the introduction of chiral appendages in π -conjugated systems as organic semiconductors in photonics and electronics opens the way to highly innovative technological applications such as transistors based on enantioselective sensing layers for chiral analytes in the vapour phase,²² enantiopure chiral magnets,²³ chirality-induced spin selectivity (CISS),²⁴ enantioselective electrochemical sensors,²⁵ detection/production of circularly polarized (CP) light.²⁶ In these devices, the active layers made of a chiral π -

^a CNR-ICCOM Istituto di Chimica dei Composti Organometallici, Via Edoardo Orabona 4, Bari 70126, Italy. E-mail: hassan@ba.iccom.cnr.it

^b Dipartimento di Chimica e Chimica Industriale, Università di Pisa, Via Giuseppe Moruzzi 13, Pisa 56124, Italy

† Electronic supplementary information (ESI) available: Experimental details and supplementary figures. See DOI: 10.1039/d1nj02125g

‡ Present address: Valsynthese SA, Fabrikstrasse 48, CH-3900 Brig, Switzerland.

§ Present address: Dipartimento di Chimica, Università degli Studi di Bari "Aldo Moro", Via Edoardo Orabona 4, 70126 Bari, Italy. E-mail: gianluigi.albano@uniba.it

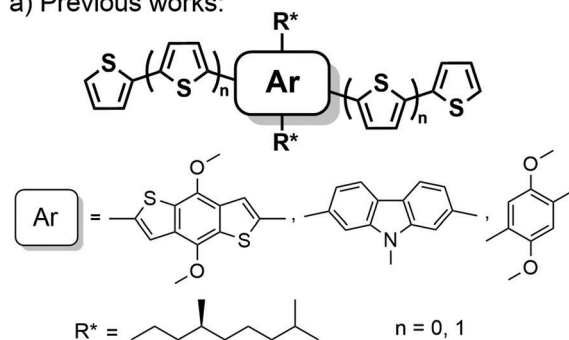
1 conjugated system (in particular small molecules) as neat
2 materials are still relatively uncommon, although this repre-
3 sents a clear advantage in terms of simplicity, reproducibility
4 and ease of preparation.

5 Instrumental techniques able to characterize aggregation
6 states at multiple hierarchical levels are fundamental for clar-
7 ifying the structure–property relationships of such
8 systems.^{11,12,27} Chiroptical spectroscopies, and in particular
9 electronic circular dichroism (ECD), play a central role in the
10 elucidation of the aggregation pattern of chiral π -conjugated
11 systems, providing valuable information about the first level of
12 supramolecular organization at the nanometre scale for both
13 solution and solid states, often complementary to that provided
14 by microscopy techniques.^{28–31}

15 Thiophene-based polymers represent one of the most inter-
16 esting families of conductive polymers for optoelectronic
17 applications.^{32,33} It is not surprising that chiral oligo- and
18 poly-thiophene derivatives are frequently investigated by ECD
19 spectroscopy. An interesting motif for the ECD study is of
20 oligo(arylenethienylene)s, whose general structure is shown in
21 Fig. 1a. In these structures, the substituent introduced on the
22 central arylene core tunes their supramolecular organization
23 depending of the chemical nature of the chiral substituent
24 group, as demonstrated in our previous studies. We studied in
25 detail the chiroptical properties of a series of structurally
26 related (arylenethienylene)s with benzo[1,2-*b*:4,5-*b'*]
27 dithiophene, 9*H*-carbazole and 1,4-phenylene cores bearing
28 the 3,7-dimethyl-1-octyl alkyl chain as chiral element
29 (Fig. 1a),^{34–38} this chiral appendage was easily prepared in
30 enantiopure form from a natural terpene (*i.e.*, β -citronellol),³⁹
31 and was devoid of functionalities and transfers its molecular
32 asymmetry (*i.e.*, a chiral centre at position 3 of the aliphatic
33 chain) to the entire system (*i.e.*, chirality at supramolecular
34 and mesoscopic scales) exclusively through dispersion forces and
35 steric effects. Another approach involves the use of proteino-
36 genic α -amino acids or monosaccharides as chiral pendant
37 groups; on one hand, the α -amino acid units are able to form
38 multiple and highly directional non-covalent interactions
39 through hydrogen-bonding, and therefore may impact in dif-
40 ferent ways the chiral supramolecular organization of the
41 oligo(arylenethienylene) backbone; on the other, the monosac-
42 charides units are chiral and polar bulky appendages which
43 may lead to reduced torsional freedom.

44 In this work, we investigated the (chiro)optical properties of
45 a family of disubstituted phenylene–thiophene oligomers **1–3**,
46 decorated on the phenyl ring by enantiopure α -amino acids or
47 monosaccharide units (Fig. 1b). We selected the *t*BOC-
48 protected *L*-phenylalanine moiety, linked to the phenylene core
49 through a flexible spacer (a linear six-carbon-atom oxaalkyl
50 chain), in order to favour intermolecular π -stacking and hydro-
51 gen bonding interactions. Conversely, we used peracetylated
52 glucose units directly bound to the 1,4-phenylene ring through
53 a β -glucosidic linkage, *i.e.*, without a linear spacer, in order
54 to evaluate the effect of steric hindrance. ECD measurements on
55 chiral conjugated structures **1–3** were performed in both
56 solution aggregation and thin film states to provide unique

a) Previous works:



b) This work:

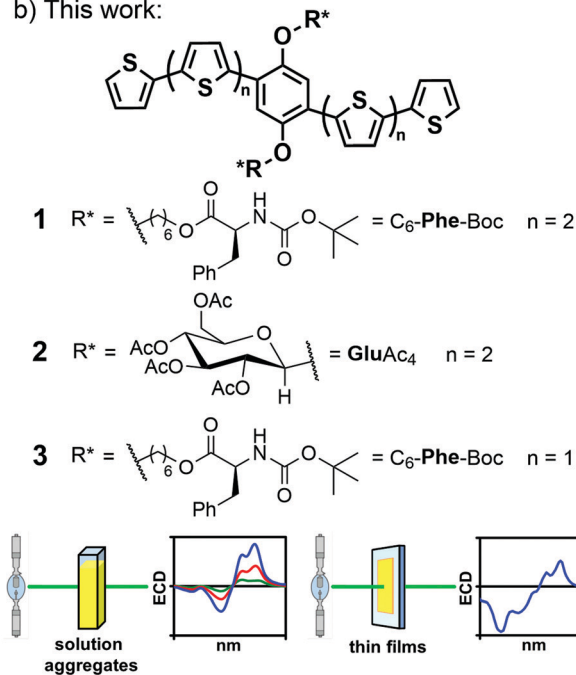


Fig. 1 Scope of the work: three *L*-phenylalanine- or *D*-glucose-
functionalized phenylene–thiophene oligomers **1–3** have been investi-
gated by electronic circular dichroism (ECD) spectroscopy in both solution
aggregation and thin film states, in order to study the impact of chirality on
their aggregation modes.

information about their first level of supramolecular organiza-
tion, discovering in some cases the co-existence of multiple
aggregation pathways which would not have been identified by
using UV-Vis absorption spectroscopy alone.

Results and discussion

The three chiral oligo(*p*-phenylenethienylene)s **1–3** were synthe-
sized in good yields according to our previously reported
procedure,⁴⁰ a versatile synthetic protocol based on the
Suzuki–Miyaura cross-coupling reaction.^{41,42}

The aggregation modes of the phenylene–thiophene oligo-
mers **1–3** were first studied in solution. In particular, their
behaviour in different solvent mixtures, consisting of a good
solvent (*i.e.*, chloroform or THF) and a poor solvent or non-
solvent (*i.e.*, methanol or water) was investigated. The relevant

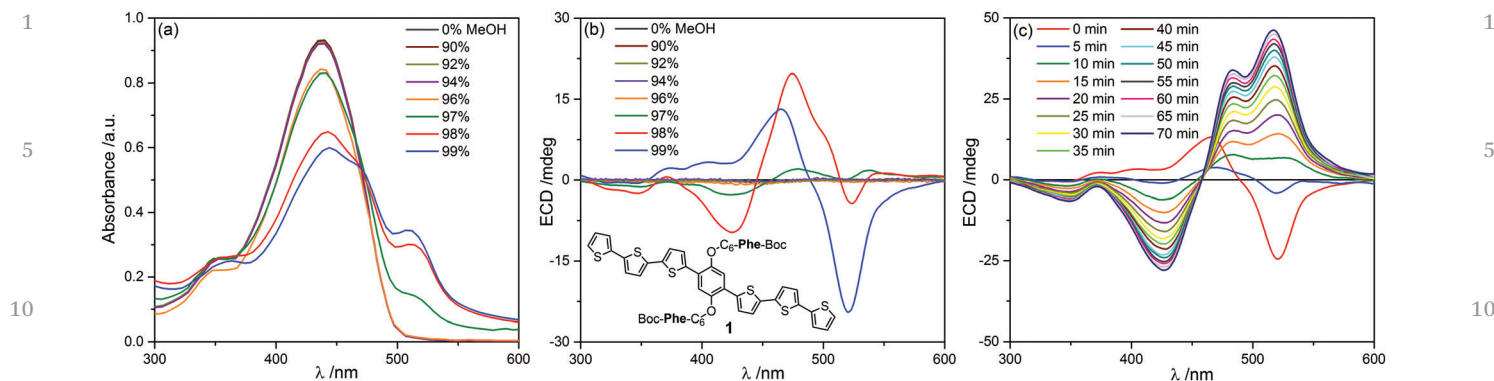


Fig. 2 (Chiro)optical characterization of *L*-phenylalanine-functionalized phenylene–thiophene oligomer **1** in solution; (a) UV-Vis absorption and (b) ECD spectra in $\text{CHCl}_3/\text{MeOH}$ mixtures with increasing amounts of MeOH; (c) evolution of the ECD spectrum in 99% MeOH as a function of time. Sample concentration: 6×10^{-5} M; cell length: 0.2 cm.

solutions or suspensions were characterized by UV-Vis absorption and ECD spectroscopies. The spectral profiles obtained by the progressive addition of a non-solvent represent the step-by-step picture of the aggregation process. The results obtained by this experimental approach are expected to be in good agreement with the material's solid state behaviour, including thin films.

The (chiro)optical characterization of *L*-phenylalanine-functionalized oligomer **1** was performed in $\text{CHCl}_3/\text{CH}_3\text{OH}$ mixtures (Fig. 2). As it is soluble in chlorinated solvents, only molecularly dispersed non-aggregated species can be observed in CHCl_3 solution. The UV-Vis absorption spectrum in CHCl_3 (Fig. 2a, grey line) showed a broad band with the maximum centered at 440 nm, attributable to the $\pi \rightarrow \pi^*$ transition of the *p*-phenylene–thiophene conjugated backbone. The corresponding ECD spectrum (Fig. 2b, grey line) was very weak with dissymmetry factors g_{abs} (defined as $\Delta\epsilon/\epsilon$ or $\Delta A/A$) of about 10^{-5} , indicating that the perturbation exerted on the intrinsically achiral π -conjugated chromophore of **1** by individual chiral centres of the two *L*-phenylalanine moieties is very small in dissolved molecules. The formation of solution supramolecular aggregates can be favoured by the addition of significant amounts of methanol as a poor solvent. In the UV-Vis absorption spectra recorded from a mixed solvent (Fig. 2a), we observed solvatochromism starting from 97% MeOH; the main band was slightly red-shifted and included a shoulder at 481 nm, while a new narrow band (the so-called “aggregation band”) appeared at 512 nm, which increased rapidly in intensity by further addition of methanol. The changes in the absorption spectrum offered an unequivocal proof for the occurrence of solution aggregates; on one hand, the aggregation induces an increased planarization of the oligothiophene backbone which causes the slight red-shift of the pristine absorption band;^{43–47} on the other hand, the new band at 512 nm is a typical manifestation of self-assembly promoted by π -stacking between the aromatic backbones as well as by ancillary interactions between the side-chains^{48,49} (in the current case, hydrogen bonding of *L*-phenylalanine moieties). However, ECD spectra (Fig. 2b) provided us more insight into

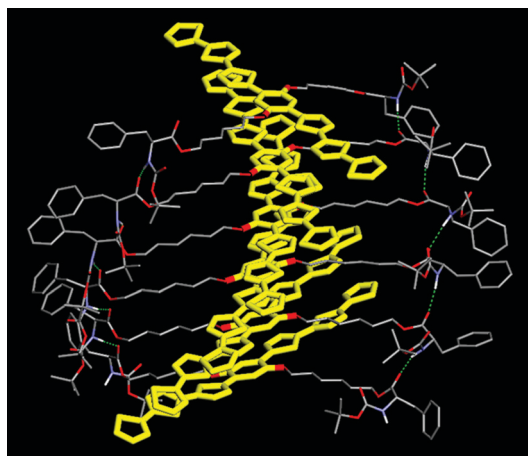
the aggregation modes of chiral oligomer **1**. By progressive addition of methanol, significantly different ECD profiles were observed; in 97% MeOH, a weakly structured ECD signal appeared which may be described as a positive exciton couplet, with maximum dissymmetry factor $g_{\text{abs}} = 1.3 \times 10^{-4}$ (Fig. 2b, green line); in 99% MeOH, a more pronounced negative exciton couplet was found, which attained maximum $g_{\text{abs}} = -2.5 \times 10^{-3}$ (Fig. 2b, blue line), almost 250-fold larger than that of molecularly isolated species in pure chloroform. In both cases, a pronounced long-wavelength tailing appeared as a typical distortion, due to the scattering by suspended particles, confirming the presence of solution aggregates. A further aspect came from the study of the ECD spectrum in 99% MeOH as a function of time (Fig. 2c); we clearly observed a progressive evolution from the above-mentioned negative exciton couplet ($t = 0$ min) to an asymmetric positive exciton couplet ($t = 70$ min), having a similar shape (but higher intensity) to the spectrum in 97% MeOH. In order to confirm the genuinity of these ECD results, we have also ruled out the occurrence of any kind of artifact due to linear dichroism (LD) by simultaneously recording the corresponding LD spectra; as expected, in all cases we found negligible signals.

All these results suggested that the chiral oligomer **1** could adopt two concomitant aggregation pathways, leading to different supramolecular structures held together by the π -stacking of conjugated backbones and the hydrogen bonding of α -amino acid moieties in the side chains, with a well-defined helicity between π -conjugated chains packed on the top of each other; a kinetic or metastable aggregation state, associated with the negative ECD couplet and characterized by a left-handed supramolecular helicity, or a thermodynamic aggregation state, responsible for the positive ECD couplet due to a right-handed supramolecular helicity. In this context, the amount of methanol plays an important role in determining which aggregation pathway will be assumed. As the good solvent (*i.e.*, CHCl_3) may help in supramolecular rearrangements, by increasing the fraction of the non-solvent (*i.e.*, MeOH) the evolution from the kinetic to thermodynamic aggregation state will be slower. Therefore, in 97% MeOH, chiral oligomer **1** self-assembles

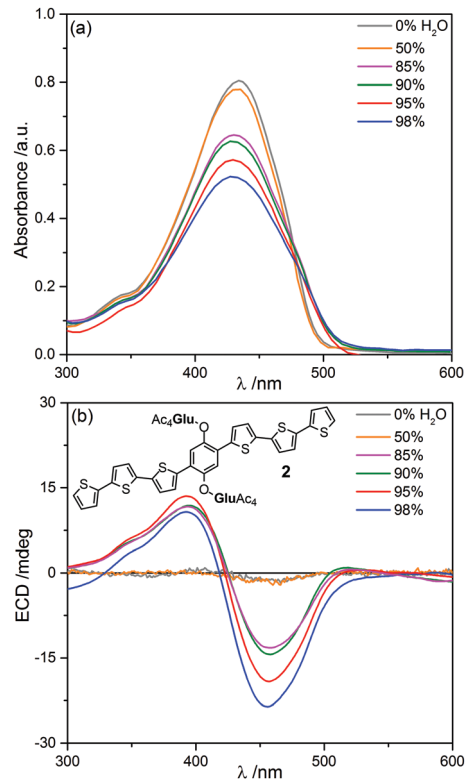
1 directly into the most thermodynamically stable aggregation
 2 state, while in 99% MeOH, it is first trapped into the kinetic
 3 state, from which progressively evolves through a simple re-
 4 ordering process towards the thermodynamic one within 70
 5 minutes (Fig. 2c). Conversely, the evolution of UV-Vis absorp-
 6 tion in 99% MeOH as a function of time (see Fig. S1 in ESI†)
 7 showed no changes, confirming that ECD spectroscopy is able
 8 to reveal the multiple aggregation pathways of chiral π -
 9 conjugated materials, a field of investigation which is limitedly
 10 accessible by other spectroscopies.

11 In order to describe which intermolecular interactions are
 12 responsible for the solution aggregates, we also developed a
 13 computational model for a hexamer of aggregated chains of the
 14 phenylene–thiophene oligomer **1** (Fig. 3), showing π -conjugated
 15 backbones arranged into a right-handed supramolecular helicity
 16 (*i.e.*, in agreement with the above mentioned thermody-
 17 namic aggregation state). The backbones are at a reciprocal
 18 distance of 4 Å and present a twist angle $\tau = 15^\circ$. The network
 19 of hydrogen bonds is represented by dotted green lines. Although
 20 the model was obtained with a simplified molecular modelling
 21 procedure (see ESI†), it confirmed that hydrogen bonding is
 22 largely responsible for the formation of stable supramolecular
 23 aggregates, inducing the helical twist between adjacent π -
 24 conjugated backbones responsible for the exciton couplet in
 25 the ECD spectrum. Moreover, similar to some α -amino acid-
 26 functionalized poly(*p*-phenylene ethynylene)s previously
 27 reported in the literature,⁵⁰ the hydrogen bonding network
 28 resembles that of polypeptide parallel β -sheets, with both
 29 amide groups of each molecule acting as hydrogen-bond
 30 donors for the ester moieties of the adjacent one.

31 With peracetylated *D*-glucopyranose units directly bound to
 32 the π -conjugated backbone without any linear spacer, chiral
 33 phenylene–thiophene **2** is devoid of hydrogen bond donors and
 34 is more sterically hindered. These differences with respect to
 35 oligomer **1** clearly influenced its optical and chiroptical



36 Fig. 3 Computational model for a hexamer of aggregated chains of
 37 phenylene–thiophene oligomer **1**, showing the right-handed supramole-
 38 cular helicity of π -conjugated backbones (in yellow), the relative dispo-
 39 sition of the L-phenylalanine-functionalized side chains (in grey), and the
 40 hydrogen bond network (dotted green lines).



41 Fig. 4 (Chiro)optical characterization of *D*-glucose-functionalized phe-
 42 nylene–thiophene oligomer **2** in solution; (a) UV-Vis absorption and (b)
 43 ECD spectra in THF/H₂O mixtures with increasing amounts of H₂O.
 44 Sample concentration: 6×10^{-5} M; cell length: 0.2 cm.

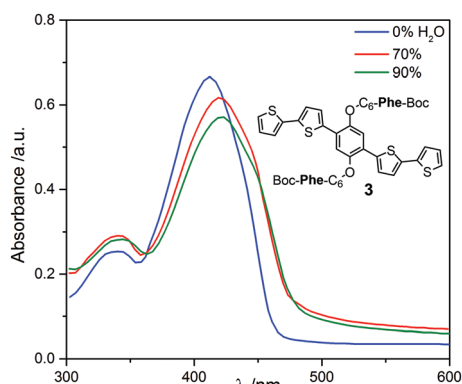
45 characteristics. In CHCl₃/CH₃OH mixtures, the absorption
 46 spectra of **2** (Fig. S2a in ESI†) showed only negligible differ-
 47 ences with that recorded in pure CHCl₃ and no evident relation-
 48 ship between the spectral profile and the increasing amount of
 49 non-solvent could be found. Likewise, the ECD spectra (Fig.
 50 S2b, ESI†) were very weak and with dissymmetry factors g_{abs}
 51 always $< 10^{-4}$. Such results suggested the absence of supramole-
 52 cular aggregates in CHCl₃/CH₃OH solvent mixtures; the
 53 increased steric hindrance, together with the absence of sites
 54 acting as hydrogen bonding donors, makes the self-assembly of
 55 π -conjugated backbones of **2** into chiral supramolecular aggre-
 56 gates more difficult with respect to α -amino acid-functionalized
 57 oligomer **1**. However, a different behaviour was found by
 58 changing the solvent mixture, using THF as the “good” solvent
 59 and H₂O as the “poor” solvent (Fig. 4). The UV-Vis absorption
 60 spectrum of **2** in pure THF (Fig. 4a, grey line) showed a broad
 61 band centered at 435 nm, due to the $\pi \rightarrow \pi^*$ transition of the *p*-
 62 phenylene–thiophene backbone. With the addition of consider-
 63 able amounts of H₂O, we observed the gradual appearance of a
 64 small shoulder at 479 nm attributable to the formation of
 65 solution aggregates. ECD spectra (Fig. 4b) again provided us
 66 more information on the self-assembly behaviour of the chiral
 67 oligomer **2**; starting from 85% content in H₂O, the occurrence
 68 of a negative exciton couplet was found, with an increasing
 69 intensity at higher amounts of non-solvent, attaining the max-
 70 imum g_{abs} of -1.7×10^{-3} in 98% H₂O (Fig. 4b, blue line).

1 Moreover, the ECD spectra did not show changes as a function of
2 time.

3 These findings pointed at the occurrence of a single, ther-
4 modynamically stable aggregation mode in oligomer **2**; a
5 supramolecular structure with π -conjugated chains arranged
6 into a right-handed helicity, responsible for the negative ECD
7 couplet. In particular, the steric hindrance of monosaccharide
8 units directly linked to the phenylene moiety increased the
9 distance between adjacent π -conjugated chains, making the
10 exciton coupling less effective with respect to oligomer **1**, and
11 thus justifying the lower values of dissymmetry factor attained
12 ($g_{\text{abs}}^{\text{max}} = -2.5 \times 10^{-3}$ in 99% MeOH for **1** vs. $g_{\text{abs}}^{\text{max}} = -1.7 \times 10^{-3}$
13 in 98% H₂O for **2**).

14 The chiral phenylene–thiophene **3** was successfully used as
15 an active layer in a sensor based on organic thin film transistors
16 (OTFTs) for the discrimination of chiral analytes in the vapour
17 phase.²² However, here for the first time we report its spectro-
18 scopic and chiroptical features. This material structurally dif-
19 fers from **1** for the shorter length of π -conjugated
20 oligothiophene segments (two bithiophene units instead of
21 two terthiophene ones). Interestingly, **3** showed a dramatically
22 different behaviour from **1** in solution aggregation processes.
23 We first investigated its (chiro)optical properties in CHCl₃/
24 CH₃OH mixtures (see Fig. S3 in the ESI†). Surprisingly, the
25 absence of supramolecular aggregates was observed, with very
26 weak ECD signals (g_{abs} of about 10^{-5}) regardless of the amount
27 of MeOH added as the non-solvent. This system is less prone to
28 self-assembly than **1** under the same conditions, probably
29 because π – π interactions are weaker due to the shorter length
30 of π -conjugated segments.

31 By switching to other solvent/non-solvent mixtures, a par-
32 tially different behaviour was observed in THF/H₂O (Fig. 5). The
33 UV-Vis absorption spectrum of **3** in pure THF (Fig. 5, blue line)
34 showed the main broad band with maximum at 413 nm,
35 associated with the $\pi \rightarrow \pi^*$ transition. By increasing the
36 amount of non-solvent, we observed clear solvatochromism
37 starting from 70% H₂O (Fig. 5, red line), with a 10 nm red-
38 shifted main absorption band and the appearance of a



39 Fig. 5 Optical characterization of L-phenylalanine-functionalized pheny-
40 lene–thiophene oligomer **3** in solution; UV-Vis absorption spectra in THF/
41 H₂O mixtures with increasing amounts of H₂O. Sample concentration: $6 \times$
42 10^{-5} M; cell length: 0.2 cm.

43 shoulder at 457 nm, attesting the occurrence of supramolecular
44 aggregation. The corresponding ECD spectra (see Fig. S4 in
45 ESI†) were very weak, with dissymmetry factors $g_{\text{abs}} \sim 10^{-5}$.
46 Such results could be explained by hypothesizing the self-
47 assembly of oligomer **3** chains into ordered achiral supramo-
48 lecular structures, where the perfect cofacial stacking of π -
49 conjugated chains is preferred compared to a helical twisted
50 one. This result could be influenced by a hydrophobic effect
51 caused by the use of H₂O as a non-solvent. A different possible
52 explanation involves the participation of water in supramole-
53 cular aggregation, by the formation of bridge bonds between
54 two adjacent oligomer molecules thanks to hydrogen bonds;
55 this would create less torsion between adjacent π -conjugated
56 chains, favouring an achiral ordered stacking.

57 As mentioned above, the investigation of aggregation modes
58 in solution is often preparatory to their study in the solid state.
59 In particular, the ability to recognize and control different
60 aggregation modes in the case of thin films would be useful
61 for obtaining layers of chiral π -conjugated systems with highly
62 tunable properties for innovative optoelectronic applications.²²
63 For this reason, the chiro(optical) properties of our chiral
64 oligomers **1**–**3** were also investigated in thin films.

65 ECD measurements in thin films require more attention
66 than in isotropic solution, as they may result in more complex
67 spectra due to the interference of macroscopic anisotropies
68 often present in the solid state, *i.e.* linear dichroism (LD) and
69 linear birefringence (LB), which provide a significant contribu-
70 tion to the measured ECD spectrum.^{51–53} In particular, the
71 experimental ECD signal of a thin film is the sum of two main
72 contributions: (a) the intrinsic isotropic component of circular
73 dichroism, termed CD_{iso} , which is invariant upon sample
74 orientation and (b) the LDLB, arising from the combined effect
75 of macroscopic anisotropies LD and LB, which inverts the sign
76 by sample flipping.¹⁵ CD_{iso} and LDLB terms can be considered
77 as the chiroptical response of two different hierarchical scales
78 of chirality; CD_{iso} is related to molecular chirality and supra-
79 molecular chiral aggregates, while, in certain cases, their
80 further organization into mesoscopic chiral domains can gener-
81 ate a LDLB contribution.¹⁵ Therefore, recognizing CD_{iso} and
82 LDLB in the ECD spectra of thin films can provide useful
83 information on their solid state organization at different hier-
84 archy levels.

85 We prepared thin films of chiral phenylene–thiophene oli-
86 gomers **1**–**3** by drop casting (DC) technique; $\sim 100 \mu\text{L}$ of a
87 chloroform solution of each compound (concentration: $1.0 \times$
88 10^{-3} M) were deposited on quartz plates, followed by slow
89 evaporation of the solvent in a closed chamber saturated with
90 CHCl₃ vapours. All the obtained samples, with a thickness of
91 about 100 nm, appeared macroscopically homogeneous and
92 semi-transparent.

93 To establish the intrinsic anisotropy of the investigated
94 samples, for each thin film sample the ECD was recorded for
95 both front and back side. In all the cases we found same
96 spectral profiles. Therefore, in the following section, we can
97 safely neglect any contribution due to LDLB and attribute all
98 ECD signals to a pure CD_{iso} arising from the occurrence of plain
99

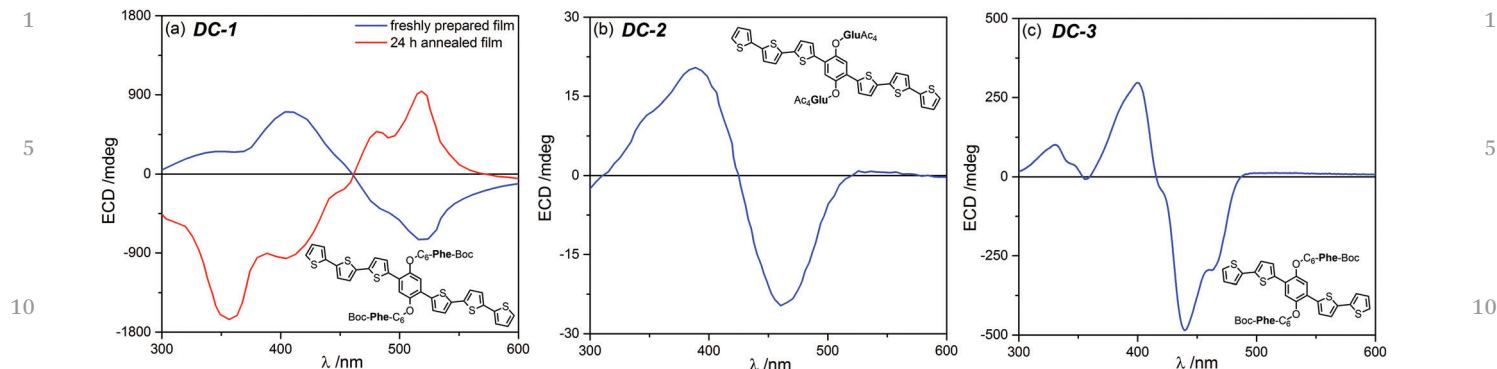


Fig. 6 ECD spectra (normalized with respect to maximum absorbance) recorded for chiral phenylene–thiophene oligomers **1–3** as thin films prepared by drop casting from a CHCl_3 solution; (a) drop-casted L-phenylalanine-functionalized oligomer **1** (**DC-1**), as a freshly prepared sample (blue line) and after 24 h of solvent annealing (red line); (b) drop-casted peracetylated D-glucose-functionalized oligomer **2** (**DC-2**); (c) drop-casted L-phenylalanine-functionalized oligomer **3** (**DC-3**).

supramolecular chiral aggregates. We could hypothesize that the steric hindrance of these chiral supramolecular aggregates (responsible for the CD_{iso} contribution) did not allow their further organization into highly anisotropic domains (*i.e.*, characterized by high LD and LB), thus resulting in the absence of any LDLB contribution.

Drop-cast samples of L-phenylalanine-functionalized phenylene–thiophene oligomer **1** (**DC-1**) showed an ECD spectrum with a strong asymmetric negative exciton couplet, with maximum g_{abs} values of -3.9×10^{-2} at 516 nm and $+4.0 \times 10^{-2}$ at 403 nm (Fig. 6a, blue line and Table 1, entry 1). However, a further interesting aspect came from solvent annealing, a typical procedure used to obtain more homogeneous thin films. **DC-1** samples left for 24 h in a chamber saturated with vapours of a “good” solvent (*i.e.*, CHCl_3) yielded dramatic changes in the ECD spectrum; a complex, asymmetric positive exciton couplet appeared, with maximum g_{abs} of $+5.4 \times 10^{-2}$ at 519 nm and -9.7×10^{-2} at 357 nm (Fig. 6a, red line and Table 1, entry 2). Interestingly, the ECD spectral shape of freshly prepared and annealed **DC-1** samples resembled that observed for **1** in 99% MeOH at, respectively, $t = 0$ min and $t = 70$ min (Fig. 6a vs. 2c). Therefore, the two competing aggregation pathways previously observed for oligomer **1** in solution also occurred in thin films; a kinetic supramolecular order, associated with the negative ECD couplet (Fig. 6a, blue line)

and characterized by a left-handed supramolecular helicity, is reached in the short time required for solvent evaporation in the drop-casting process; during the subsequent solvent annealing process, π -conjugated chains rearrange into a thermodynamically stable aggregation state, responsible for the positive ECD couplet (Fig. 6a, red line), due to a right-handed supramolecular helicity. We also recorded the IR spectra of **DC-1** as a freshly prepared sample and after 24 h of solvent annealing (Fig. S6, ESI[†]); although both provided evidence of hydrogen bonding interactions, thus confirming their role in the supramolecular organization of the oligomer **1**, no significant differences were actually found between them. This is a further confirmation of the importance of ECD spectroscopy in detecting the possible co-existence of multiple aggregation pathways, succeeding where other techniques failed.

Concerning oligomer **2**, the steric hindrance of monosaccharide units exerted a clear effect on its chiroptical properties in thin films. In fact, the ECD spectrum of **DC-2** showed much less intense signals with respect to **DC-1** (almost two orders of magnitude smaller); a negative exciton couplet with dissymmetry factor g_{abs} values of -9.4×10^{-4} at 460 nm and $+8.1 \times 10^{-4}$ at 389 nm (Fig. 6b and Table 1, entry 3), suggesting the occurrence of a supramolecular architecture with a right-handed helicity, although with a less effective exciton coupling of π -conjugated chains with respect to **1**. In this case, we observed no evolution of the chiroptical signals upon solvent annealing, confirming the hypothesis of formation of a single, thermodynamically stable aggregation state.

However, the most interesting chiroptical properties were observed for drop-cast thin films of chiral phenylene–thiophene oligomer **3** (**DC-3**), since it did not show any significant ECD signal in a variety of solvent mixtures. The ECD spectrum of **DC-3** showed a complex and strong, asymmetric negative couplet, which attained a maximum dissymmetry factor g_{abs} value of -1.7×10^{-2} at 440 nm and $+1.2 \times 10^{-2}$ at 400 nm. Such results could be explained by hypothesizing the occurrence of a chiral supramolecular architecture with a strong left-handed helical twist between adjacent π -conjugated backbones. **DC-3** samples retained the same ECD spectrum after solvent

Table 1 Chiroptical properties of chiral phenylene–thiophene oligomers **1–3** as thin films prepared by drop casting from a CHCl_3 solution and dissymmetry factor g_{abs} values recorded at two representative wavelengths

Entry	Thin film sample	λ (nm)	g_{abs}
1	DC-1 (freshly prepared)	403	$+4.0 \times 10^{-2}$
		516	-3.9×10^{-2}
2	DC-1 (24 h annealed)	357	-9.7×10^{-2}
		519	$+5.4 \times 10^{-2}$
3	DC-2	389	$+8.1 \times 10^{-4}$
		460	-9.4×10^{-4}
4	DC-3	400	$+1.2 \times 10^{-2}$
		440	-1.7×10^{-2}

annealing (24 h under CHCl₃ vapours), suggesting that the observed aggregation mode corresponded to an energetically stable situation. Moreover, since chiral oligomer **3** is the only compound whose chiroptical properties in solution and films were not similar, we ruled out the occurrence of any kind of artifact by repeating the same measurements on thin films of the opposite enantiomer of **3**, *i.e.*, D-phenylalanine-decorated phenylene–thiophene oligomer (*ent*-**3**); a specular ECD spectrum with a positive couplet was obtained (Fig. S7, ESI[†]), confirming the total reproducibility of our results. However, this difference between solution and thin films ECD properties is actually not surprising; if oligomer **3** is less prone to self-assembly in solution because π – π interactions are less effective due to the shorter π -conjugated length, then with the removal of CHCl₃ by slow evaporation in drop-cast thin films, an efficient π -packing into chiral supramolecular architectures is still possible.

It is worth emphasizing that phenylene–thiophene oligomers **1** and **3** exhibited intense ECD signals in thin films, with g_{abs} values in the order of 10^{-2} (and up to about 0.1 in the case of annealed DC-1 sample). This ability of displaying high (or very high) discrimination in the absorption of CP light would be highly desirable in thin films of chiral organic semiconductors, as it could open the way to their application as active layers in innovative optoelectronic devices, for example CP light-detecting organic field-effect transistors.^{54,55}

30 Conclusions

In conclusion, we studied the (chiro)optical properties of three chiral phenylene–thiophene oligomers functionalized with enantiopure L-phenylalanine or D-glucopiranoside moieties, in both solution aggregation and thin film states. We investigated the impact of chirality on their aggregation modes, evaluating the effect of both intermolecular hydrogen bonding and dispersion forces, as well as the role played by steric hindrance and π -conjugated length in promoting different supramolecular architectures. We demonstrated that ECD provides unique and valuable information about the first level of supramolecular organization, discovering in some cases the co-existence of multiple aggregation pathways which would not have been identified by using only UV-Vis absorption spectroscopy. The ability to recognize and finely control these different aggregation modes, enabled by our study, is very useful for obtaining layers of chiral organic semiconductors with highly tunable properties (in particular with strong dissymmetry factor g_{abs} values) especially in the case of thin films for optoelectronic applications.

55 Conflicts of interest

There are no conflicts to declare.

Acknowledgements

The authors are indebted to Prof. Gennaro Pescitelli for his help in the computational model calculations and to Prof. Lorenzo Di Bari for the fruitful discussions.

Notes and references

- O. Ostroverkhova, *Chem. Rev.*, 2016, **116**, 13279–13412.
- P. Friederich, A. Fediai, S. Kaiser, M. Konrad, N. Jung and W. Wenzel, *Adv. Mater.*, 2019, **31**, 1808256.
- U. Mitschke and P. Bäuerle, *J. Mater. Chem.*, 2000, **10**, 1471–1507.
- G. M. Farinola and R. Ragni, *Chem. Soc. Rev.*, 2011, **40**, 3467–3482.
- A. Facchetti, *Chem. Mater.*, 2011, **23**, 733–758.
- C. Wang, H. Dong, W. Hu, Y. Liu and D. Zhu, *Chem. Rev.*, 2012, **112**, 2208–2267.
- M. A. Rahman, P. Kumar, D.-S. Park and Y.-B. Shim, *Sensors*, 2008, **8**, 118–141.
- J. Cornil, D. Beljonne, J.-P. Calbert and J.-L. Brédas, *Adv. Mater.*, 2001, **13**, 1053–1067.
- F. J. M. Hoeben, P. Jonkheijm, E. W. Meijer and A. P. H. J. Schenning, *Chem. Rev.*, 2005, **105**, 1491–1546.
- T. L. Andrew and T. M. Swager, *J. Polym. Sci., Part B: Polym. Phys.*, 2011, **49**, 476–498.
- A. Salleo, R. J. Kline, D. M. DeLongchamp and M. L. Chabinyc, *Adv. Mater.*, 2010, **22**, 3812–3838.
- J. Rivnay, S. C. B. Mannsfeld, C. E. Miller, A. Salleo and M. F. Toney, *Chem. Rev.*, 2012, **112**, 5488–5519.
- P. A. Korevaar, T. F. A. de Greef and E. W. Meijer, *Chem. Mater.*, 2014, **26**, 576–586.
- Y. Dorca, E. E. Greciano, J. S. Valera, R. Gómez and L. Sánchez, *Chem. – Eur. J.*, 2019, **25**, 5848–5864.
- G. Albano, G. Pescitelli and L. Di Bari, *Chem. Rev.*, 2020, **120**, 10145–10243.
- M. Verswyvel and G. Koeckelberghs, *Polym. Chem.*, 2012, **3**, 3203–3216.
- Y. Yang, Y. Zhang and Z. Wei, *Adv. Mater.*, 2013, **25**, 6039–6049.
- M. Liu, L. Zhang and T. Wang, *Chem. Rev.*, 2015, **115**, 7304–7397.
- C. Resta, S. Di Pietro, M. Majerić Elenkov, Z. Hameršak, G. Pescitelli and L. Di Bari, *Macromolecules*, 2014, **47**, 4847–4850.
- C. C. Lee, C. Grenier, E. W. Meijer and A. P. H. J. Schenning, *Chem. Soc. Rev.*, 2009, **38**, 671–683.
- L. A. P. Kane-Maguire and G. G. Wallace, *Chem. Soc. Rev.*, 2010, **39**, 2545–2576.
- L. Torsi, G. M. Farinola, F. Marinelli, M. C. Tanese, O. H. Omar, L. Valli, F. Babudri, F. Palmisano, P. G. Zambonin and F. Naso, *Nat. Mater.*, 2008, **7**, 412–417.
- C. Train, M. Gruselle and M. Verdager, *Chem. Soc. Rev.*, 2011, **40**, 3297–3312.
- P. C. Mondal, C. Fontanesi, D. H. Waldeck and R. Naaman, *Acc. Chem. Res.*, 2016, **49**, 2560–2568.

- 1 25 J. G. Ibanez, M. E. Rincón, S. Gutierrez-Granados, M. H. Chahma, O. A. Jaramillo-Quintero and B. A. Frontana-Urbe, *Chem. Rev.*, 2018, **118**, 4731–4816.
- 26 J. R. Brandt, F. Salerno and M. J. Fuchter, *Nat. Rev. Chem.*, 2017, **1**, 0045.
- 27 V. Palermo and P. Samorì, *Angew. Chem., Int. Ed.*, 2007, **46**, 4428–4432.
- 28 N. Berova, L. D. Bari and G. Pescitelli, *Chem. Soc. Rev.*, 2007, **36**, 914–931.
- 10 29 G. Pescitelli, L. Di Bari and N. Berova, *Chem. Soc. Rev.*, 2011, **40**, 4603–4625.
- 30 G. Pescitelli, L. Di Bari and N. Berova, *Chem. Soc. Rev.*, 2014, **43**, 5211–5233.
- 31 M. Górecki, M. A. M. Capozzi, G. Albano, C. Cardellicchio, L. Di Bari and G. Pescitelli, *Chirality*, 2018, **30**, 29–42.
- 15 32 R. D. McCullough, *Adv. Mater.*, 1998, **10**, 93–116.
- 33 I. F. Perepichka, D. F. Perepichka, H. Meng and F. Wudl, *Adv. Mater.*, 2005, **17**, 2281–2305.
- 34 G. Albano, M. Lissia, G. Pescitelli, L. A. Aronica and L. Di Bari, *Mater. Chem. Front.*, 2017, **1**, 2047–2056.
- 20 35 G. Albano, F. Salerno, L. Portus, W. Porzio, L. A. Aronica and L. Di Bari, *ChemNanoMat*, 2018, **4**, 1059–1070.
- 36 G. Albano, M. Górecki, G. Pescitelli, L. Di Bari, T. Javorfi, R. Hussain and G. Siligardi, *New J. Chem.*, 2019, **43**, 14584–14593.
- 25 37 G. Albano, L. A. Aronica, A. Minotto, F. Cacialli and L. Di Bari, *Chem. – Eur. J.*, 2020, **26**, 16622–16627.
- 38 F. Zinna, G. Albano, A. Taddeucci, T. Colli, L. A. Aronica, G. Pescitelli and L. Di Bari, *Adv. Mater.*, 2020, **32**, 2002575.
- 30 39 G. Albano, T. Colli, T. Biver, L. A. Aronica and A. Pucci, *Dyes Pigm.*, 2020, **178**, 108368.
- 40 O. Hassan Omar, F. Babudri, G. M. Farinola, F. Naso, A. Operamolla and A. Pedone, *Tetrahedron*, 2011, **67**, 486–494.
- 35 41 A. Operamolla, S. Colella, R. Musio, A. Loiudice, O. Hassan Omar, G. Melcarne, M. Mazzeo, G. Gigli, G. M. Farinola and F. Babudri, *Sol. Energy Mater. Sol. Cells*, 2011, **95**, 3490–3503.
- 42 A. Ruiz-Carretero, Y. Atoini, T. Han, A. Operamolla, S. Ippolito, C. Valentini, S. Carrara, S. Sinn, E. A. Prasetyanto, T. Heiser, P. Samorì, G. Farinola and L. De Cola, *J. Mater. Chem. A*, 2019, **7**, 16777–16784.
- 43 D. Iarossi, A. Mucci, F. Parenti, L. Schenetti, R. Seeber, C. Zanardi, A. Forni and M. Tonelli, *Chem. – Eur. J.*, 2001, **7**, 676–685.
- 44 F. Brustolin, F. Goldoni, E. W. Meijer and N. A. J. M. Sommerdijk, *Macromolecules*, 2002, **35**, 1054–1059.
- 45 H. Goto, Y. Okamoto and E. Yashima, *Macromolecules*, 2002, **35**, 4590–4601.
- 46 A. Mucci, F. Parenti, R. Cagnoli, R. Benassi, A. Passalacqua, L. Preti and L. Schenetti, *Macromolecules*, 2006, **39**, 8293–8302.
- 47 G. Albano, T. Colli, L. Nucci, R. Charaf, T. Biver, A. Pucci and L. A. Aronica, *Dyes Pigm.*, 2020, **174**, 108100.
- 48 F. Babudri, D. Colangiuli, L. Di Bari, G. M. Farinola, O. Hassan Omar, F. Naso and G. Pescitelli, *Macromolecules*, 2006, **39**, 5206–5212.
- 49 G. Pescitelli, O. Hassan Omar, A. Operamolla, G. M. Farinola and L. Di Bari, *Macromolecules*, 2012, **45**, 9626–9630.
- 25 50 C. Resta, G. Pescitelli and L. Di Bari, *Macromolecules*, 2014, **47**, 7052–7059.
- 51 Y. Shindo and M. Nakagawa, *Rev. Sci. Instrum.*, 1985, **56**, 32–39.
- 52 Y. Shindo, M. Nakagawa and Y. Ohmi, *Appl. Spectrosc.*, 1985, **39**, 860–868.
- 53 R. Kuroda, T. Harada and Y. Shindo, *Rev. Sci. Instrum.*, 2001, **72**, 3802–3810.
- 54 Y. Yang, R. C. da Costa, M. J. Fuchter and A. J. Campbell, *Nat. Photonics*, 2013, **7**, 634–638.
- 35 55 L. Zhang, I. Song, J. Ahn, M. Han, M. Linares, M. Surin, H.-J. Zhang, J. H. Oh and J. Lin, *Nat. Commun.*, 2021, **12**, 142.

40

45

50

55