

Analog FM free-space optical communication based on a mid-infrared quantum cascade laser frequency comb

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Abstract: Quantum cascade laser frequency combs are nowadays well-appreciated sources for infrared spectroscopy. Here their applicability for free-space optical communication is demonstrated. The spontaneously-generated intermodal beat note of the frequency comb is used as carrier for transferring the analog signal via frequency modulation. Exploiting the atmospheric transparency window at 4 μ m, an optical communication with a signal-to-noise ratio up to 65 dB is realized, with a modulation bandwidth of 300 kHz. The system tolerates a maximum optical attenuation exceeding 35 dB. The possibility of parallel transmission of an independent digital signal via amplitude modulation at 5 Mbit/s is also demonstrated.

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1. Introduction

The mid-infrared (MIR) region of the electromagnetic spectrum ($\lambda > 3\mu$ m) is of particular interest both from a fundamental and an applied research point of view. Here fall in fact the strongest absorptions related to the ro-vibrational transitions of light molecules of environmental/atmospheric interest such as water, carbon dioxide, methane and nitrogen oxides. From a fundamental viewpoint, the study of the absorption spectra with ultrahigh precision and sensitivity can give insights into fundamental laws of physics [1–7]. From an applied perspective, the high-sensitivity detection of such gases through advanced spectroscopy techniques [8,9] is important for environmental and health monitoring [10]. The MIR region is not completely covered by atmospheric gases absorptions: Two main transparency windows (one around 4μ m and one around 10μ m) are present. Here the atmospheric transmittance (per km) is > 80% (with peaks of 99.965%) and they can be conveniently exploited for free-space optical communication (FSOC) [11,12]. In addition, the attenuation due to the aerosol is also significantly lower compared to visible and near-infrared wavelengths [13].

Since its demonstration in 1994 at Bell Labs [14], the quantum cascade laser (QCL) has been largely used and appreciated as coherent source [15]. QCLs are chip-scale room-temperature current-driven semiconductor-heterostructure lasers where the laser transition takes place between specifically-tailored sublevels in quantum wells within the conduction band. This sublevels structure is given by the semiconductor layers heterostructure that creates quantum wells for the carriers. Since the structure and thickness of the layers are highly engineerable through molecular-beam epitaxy, the levels' structure can be arranged to address a wide range of wavelengths using the same material system. In combination with the possibility of using different families of semiconductor materials, the emission wavelength can be engineered in a wide range, spanning from the MIR ($\lambda > 3.15\mu$ m [16,17]) to the THz [18]. Due to the very short laser transition lifetime (<1 ps), QCLs can be modulated at high speed (GHz and above) [19,20]. Their high-power emission (~ W) [21] and current driving make them ideal candidates for infrared spectroscopy

[22–26]. QCLs can operate both in single longitudinal mode [27] and in frequency-comb regime [28–31]. Indeed, broadband QCLs [32,33] can emit frequency combs (QCL-combs) due to the high third-order nonlinearity characterizing the active region which enables a four-wave-mixing (FWM) parametric process within the waveguide [34]. FWM correlates the generated modes [35–39] establishing a fixed phase relation among them and ensuring a coherent frequency-comb emission [40–44], which proved to be successfully exploitable for dual-comb spectroscopy [45,46]. An intermodal beat-note signal (IBN) is associated with the frequency-comb emission, given by the beatings between first-neighbor modes. This signal, falling in the range 5-20 GHz, is determined by the waveguide length and effective refractive index, and it can be as narrow as few hundreds of Hz or less due to the high overall coherence characterizing QCLs and QCL-combs [41,47,48]. IBNs can be easily modulated by modulating the laser's driving current [44] or by other means (e.g. optical modulation [49] or modulating the current passing through an integrated resistive heater [50]).

The single-mode emission of QCLs has already been exploited for FSOC relying on amplitude modulation (AM) [51,52]. Regarding QCL-combs, the possibility of exploiting their IBN for transmitting an analog signal in the microwaves in a wireless-broadcasting fashion has been recently demonstrated [53]. This configuration is mainly useful for indoor short-range communication.

In this work we demonstrate the possibility of exploiting the IBN of a QCL-comb for carrying the signal with frequency-modulation (FM) over the optical channel corresponding to the atmospheric transparency window around 4 μ m and we characterize the performance of the communication system. Moreover, we demonstrate that an independent digital signal can be encoded on the same optical channel with the same device via AM. This configuration is particularly well-suited for long-distance point-to-point transmission in the atmosphere.

2. Methods and discussion

2.1. Experimental setup

The experimental setup used for demonstrating and characterizing the optical communication system is sketched in Fig. 1. The laser source utilized for the transmission is a QCL-comb designed and fabricated at ETH Zürich. The device is a buried-heterostructure Fabry-Pérot QCL able to operate in continuous wave at room temperature [40,43]. Its emission is centered around 4.70 μ m and the mode spacing (f_s) corresponding to the intermodal beat note (IBN) is about 7.06 GHz. Further details about the laser source (LIV curves and laser spectrum) are given in Supplement 1. The working temperature of the laser for this experiment varied between 12°C and 20°C. The laser radiation propagates in the free space for 2.0 m and it is sent through a variable attenuator to a fast thermoelectric-cooled photovoltaic MCT detector (PVI-4TE-8-0.5x0.5-T08-wBaF2-35 by Vigo System), operating in the spectral range from 3 μ m to 9 μ m. The detector features a two-stage amplification system: the first stage is a high-speed transimpedance, the second stage is an AC-coupled preamplifier. The bandwidth, as defined by its -3-dB cutoff, is 700 MHz.

The electronic part of the setup consists mainly of two subsystems. The first one, dedicated to signal generation, modulation and transmission, is composed of the signal generation system and the laser current driver. The laser is powered by an ultra-low-noise current driver (QubeCL, by ppqSense) featuring a module for current modulation that is fed with the signal to be transmitted. The module adds to the DC bias current a current modulation proportional to the modulation signal. The current modulation bandwidth is 2 MHz and the modulation amplitude is up to 10 mA. When a larger modulation bandwidth is required, a different current modulation scheme is used. A variable resistor is in fact placed in parallel to the laser chip with the aim of subtracting a fraction of the driving current to the chip itself. The variable resistor is a junction-gate field-effect transistor (JFET) and the modulation signal is applied to its gate pin. The current modulation

bandwidth is 50 MHz and the modulation amplitude is up to 5 mA. For demonstration purposes, the laser current is modulated with sine waves, square waves and analog audio signals generated via a function generator or a computer system (CS1). The signals to be sent to the current modulators are selected via a multiplexer (MUX). The modulation of the laser current induces a frequency modulation of the IBN of the laser. The IBN is out of the nominal bandwidth of the detector but can still be well detected anyway (see Fig. 2). At the IBN frequency, the signal-to-noise ratio (S/N) loss is about 22 dB compared to the in-bandwidth response (see Supplement 1 for a characterization). The loss of detected signal amplitude can be effectively compensated with a low-noise RF amplifier. The second subsystem of the electronic part is devoted to the extraction of the original signal from the detected IBN via a demodulation process. At first, the IBN is down-mixed from 7.06 GHz to 54.6 MHz, then it is band-pass-filtered (BPF2, band 5 MHz) and finally amplified (green box in Fig. 1). Afterwards, a delay-line frequency discriminator (DLD) - optimized for operating at frequencies around the down-mixed carrier performs the demodulation of the IBN [54] extracting the original signal (see Supplement 1 for a description of the DLD). After the DLD, a low-pass filter removes the down-mixed carrier (blue box in Fig. 1). Finally, the demodulated signal is sent to an oscilloscope to be compared with the original signal and to a second computer system (CS2) to be further analyzed.



Fig. 1. Experimental setup of the system for the optical free-space communication. The source is a MIR QCL-comb (top left). The laser's driving system and the signals encoding system is highlighted in the yellow box. On the right, the components devoted to the detection, demodulation and analysis of the signals is represented. BPF: band-pass filter. LPF: low-pass filter. CS: computer system.

2.2. Intermodal beat-note characterization

In order to perform the best FM analog transmission it was crucial to determine the working conditions of the laser (temperature and driving current) where the IBN is intense and narrow (with the smallest frequency noise) and the FM is optimized, i.e. the working condition where the frequency tuning of the IBN with bias current is large. The analysis is performed at several



Fig. 2. Spectrum of the QCL-comb IBN in free-running operation at $T = 12^{\circ}$ C and I = 730 mA detected by the MCT detector and acquired with a spectrum analyzer.

temperatures in the range from $T = 20^{\circ}$ C to $T = 12^{\circ}$ C. For each temperature, a scan of the driving current, acquiring the IBN, is performed. As an example, the IBN response at two different temperature values is reported in Fig. 3. At $T = 16^{\circ}$ C the IBN is pretty intense and narrow (the error bars representing the width are small) but the tuning with current of the IBN frequency is very small and it is almost zero around I = 740 mA. On the other hand, at $T = 12^{\circ}$ C the tuning is generally large and there are intervals where the IBN is intense and narrow. Finally, we chose $T = 12^{\circ}$ C and I = 730 mA as working conditions, since here all the considered parameters are optimized (see Fig. 2 representing the spectrum of the IBN in free-running operation). We remark that these working conditions are definitely compatible with standard operating conditions, i.e. room temperature and moderate driving current.



Fig. 3. Frequency tuning with bias current of the IBN for two different temperatures of operation of the QCL. The bars represent the width of the IBN, where the values have been multiplied by a factor of 1000 in order to have visible bars in the plots. The peak values of the IBN are color coded.

Afterwards, a characterization of the modulated IBN has been performed. With the MUX, a sinusoidal signal has been sent to the current modulator and the amplitude of the signal has been chosen as the maximum value tolerated by the laser for preserving the frequency-comb emission. In terms of driving current, the current modulation amplitude is about 1% of its DC value. The frequency of the signal has been varied to characterize the response of the system. In Fig. 4 the IBN modulated at different frequencies is shown. From the plots, a peak frequency deviation of about 200 kHz can be inferred as half width of the spectra (more evident at low modulation).



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Fig. 4. Spectra of the detected IBN with different sinusoidal frequency modulations acquired with a spectrum analyzer. The QCL-comb operates at $T = 12^{\circ}$ C and I = 730 mA. From the plots, a peak frequency deviation of about 200 kHz can be inferred as half width of the spectra (more evident at low modulation frequency, i.e. 1 kHz and 10 kHz).

2.3. Characterization of the received analog FM signal

For quantitatively estimating the quality of the transmission, the signal demodulated by the DLD has been acquired with an oscilloscope. Fig. 5-left shows the received signal obtained with a modulation frequency of 1 kHz. For performing a complete Fourier analysis of the response of the system, the acquisition has been repeated for several modulation frequencies from 100 Hz to 1 MHz. For each acquisition the corresponding fast-Fourier transform (FFT) has been computed (see Fig. 5-right for the FFT of the 1-kHz signal). From the FFTs we extracted the modulation response of the system (Fig. 6), the S/N and the harmonic distortion characterizing the received signal (Fig. 7). From the modulation response analysis, a modulation bandwidth of ~ 300 kHz defined by its -3-dB cutoff can be retrieved. This value is perfectly in agreement with the well-known single-mode QCLs' frequency modulation bandwidth [55]. The S/N has been computed as the ratio between the power of the peak corresponding to the signal and the power level corresponding to the noise floor. On the other hand, the harmonic distortion has been computed as the ratio between the power of the second harmonic to the power of the fundamental frequency peaks of the received signal. These last two parameters have been computed within a bandwidth of 20 kHz, covering a range of values absolutely compatible with a good quality audio transmission. Moreover, a substantial flatness of their values with the modulation frequency clearly emerges, confirming the good performance of the system.

2.4. Attenuation effects

For a concrete assessment of the performance of the FSOC system, an evaluation of the impact of optical losses is crucial. Optical losses can be split into two groups: static (or quasi-static)



Fig. 5. Left: Demodulated signal at 1 kHz. Right: FFT of the demodulated signal with the fundamental frequency and second harmonic peaks marked with red dots.



Fig. 6. Modulation response of the system. The bandwidth, as defined by its -3-dB cutoff, is ~ 300 kHz, perfectly in agreement with the well-known single-mode QCLs' frequency modulation bandwidth [55].



Fig. 7. Left: S/N as a function of modulation frequency. Right: Harmonic distortion as a function of modulation frequency. The harmonic distortion is computed as the ratio between the power of the second harmonic and the power of the fundamental frequency peaks of the received signal.

and dynamical. As an example, losses due to optical elements, molecular absorption, aerosol scattering, geometrical losses due to the divergence of the optical beam, belong to the first group. Losses due to atmospheric turbulence belong to the second group. The dynamical losses importantly impact the performance of FSOC systems and are hard to be modeled and simulated. Here we concentrated on the simulation of quasi-static losses effects by placing a variable optical attenuator on the beam path. In particular, we studied how the S/N changes with optical attenuation and we investigated the possibility of recovering it by adapting the electronic gain of the RF amplifier placed at the end of the IBN down-mixing stage (see Fig. 1). The chosen modulation frequency is 1 kHz, being representative for audio transmission. In Fig. 8 the dots lying on the blue curve show a S/N decreasing with increasing optical attenuation (case of fixed electronic gain). On the other hand, the brighter dots show that, when the electronic gain is increased, the S/N can be significantly recovered, up to a value of 14 dB (see the values at 34.0 dB and 35.2 dB of optical attenuation), resulting in a maximum S/N loss of 10 dB over the explored range of optical attenuation. This recovery capability indicates that in the described case the S/N is mainly limited by the acquisition electronics noise floor.



Fig. 8. S/N as a function of optical attenuation for a modulation frequency of 1 kHz. The variable electronic gain used for recovering the S/N is color coded.

For giving a practical reference of the scale of optical attenuation values we have investigated, we discuss here the equivalent conditions for an outdoor FSOC. We assume a communication distance of 5 km and an optical system optimized for minimizing the geometrical losses, i.e. for collecting all the optical power at the receiver location. We also neglect the possible effects given by atmospheric turbulence. The residual attenuation is generally given by the scattering due to molecules and particles (aerosol). In the spectral region covered by the QCL-comb (around 4.70 μ m, see Supplement 1), the average molecular absorption value is about 1.22 dB/km [11,12,56], 6.1 dB over 5 km. This means that an additional attenuation of at least 29.0 dB can be tolerated by the system, which corresponds to a visibility range [57–60] of 0.88 km, given e.g. by light-fog conditions. We remark here that, by using a QCL operating at a more favorable wavelength still within the 4- μ m transparency window, a negligible molecular absorption can be found, allowing an effective communication even with a lower visibility range.

2.5. Fidelity of the received signal

The communication system has been tested also for transmission of real audio streams. The MUX has been switched for transmitting audio signals generated by a computer system (CS1 in Fig. 1). The sent and the received signal have been acquired simultaneously with an oscilloscope for comparison, with two different levels of optical attenuation (see Fig. 9-left). The received signal looks very similar to the sent one, especially with the lower optical attenuation. For

quantitatively estimating the fidelity of the received audio signal, we chose as estimator the first-order correlation between the received and the sent signal (g_{12}) . The normalized first-order correlation is defined as

$$g_{12}(\tau) = \frac{\langle S_1(t) \cdot S_2(t-\tau) \rangle}{\sqrt{\langle S_1^2 \rangle \langle S_2^2 \rangle}},\tag{1}$$

where S_1 and S_2 are the sent and the received AC signal, respectively, t is time and τ is the time delay introduced between the two streams. The symbol $\langle \rangle$ denotes the averaging process over the arrays and the product \cdot is computed element-by-element. For the computation of the product between the two signals, the second array is shifted of an amount corresponding to τ and then the two arrays are conveniently cut in order to fit the right number of elements. The arrays are cut as less as possible and the dimension is kept constant over the whole computation. In Fig. 9-right the g_{12} is represented for two different levels of optical attenuation. The significant value is the peak of the curve which is 0.87 for the lower attenuation case and 0.75 for the higher attenuation, confirming that, providing an additional electronic amplification, for optical attenuation values within 35 dB the quality of the transmission is not significantly degraded. On the other hand, the width of the curve is related to the time response of the system. The linewidth of ~ 0.2 ms corresponds to a low-pass frequency of 5 kHz. This cutoff is probably due to the audio signal generator and is compatible with audio stream reproduction, since voice and music spectra generally fall within this frequency. Some recordings of audio tracks transmitted with the system have been performed for demonstration purposes: See Supplement 1 for a description and Visualization 1, Visualization 2, Visualization 3, and Visualization 4 for a demonstration.



Fig. 9. Left: Sent and received audio signal with two different values of optical attenuation. Right: g_{12} related to the sent and the received signal at the two different levels of optical attenuation.

2.6. Simultaneous analog FM/digital AM transmission

The realized communication system is also capable of simultaneously transmitting on the same optical channel an independent AM digital signal together with the FM analog signal. The MUX has been set to send in parallel a 100-kHz sine wave to the driver current modulation module (FM of the IBN) and a 5-MHz square wave to the JFET current modulator (AM modulation of the QCL optical power). The results are represented in Fig. 10. In this case, the optical attenuation is 21.5 dB. The thresholds set for digital communication are highlighted as red horizontal lines. The two thresholds have been selected as the maximum and the minimum signal levels where the intervals between the wavefronts of the square-wave signal does differ less than 30% with respect to the mean value (computed at the zero crossing) over the whole acquisition. The two thresholds define a window of 24 mV, which is more than sufficient for guaranteeing an error-free digital

communication. Moreover, from the same figure a good shape of the analog FM signal can be inferred. This configuration has been tested for several values of optical attenuation up to 28.7 dB. The system still maintains a good capability of transmission with a digital window of 4 mV and a good shape of the demodulated analog FM signal. This demonstrates the possibility of exploiting the full FM bandwidth of the system for analog communication (300 kHz) together with an independent digital communication at a speed of 5 Mbit/s, essentially limited by the current modulator and not by the QCL internal dynamics.



Fig. 10. Simultaneous digital AM + analog FM transmission with an optical attenuation of 21.5 dB. Top left: Received square wave. Top right: Square wave zoom. The thresholds set for digital communication are highlighted as red horizontal lines. Bottom left: Demodulated sine wave transmitted in parallel with the square wave.

3. Conclusion

In this work, we have demonstrated the possibility of exploiting the intermodal beat-note signal (IBN) spontaneously-generated by QCL-combs for performing point-to-point free-space optical communication in the MIR transparency window around 4μ m. A S/N up to 65 dB has been achieved, with a modulation bandwidth of 300 kHz. With an optical attenuation up to 35 dB the loss in terms of S/N is less than 10 dB. With this system, audio streams of very good quality have been transmitted, with a fidelity expressed in terms of first-order correlation between the sent and received signal (g_{12}) up to 87%. The transmission has been performed on a single channel (monophonic) for simplicity, but the wide bandwidth characterizing the system is potentially able to host e.g. a 38 kHz stereo subcarrier allowing the transmission of a composite stereophonic signal typically adopted for FM radio broadcasting.

The system also allows a parallel transmission of an independent digital signal via AM at a speed of 5 Mbit/s.

The presented proof-of-principle system could be improved in many aspects. First of all, the use of QCLs with optimized growth, fabrication and contacting would strongly increase the AM bandwidth [19,20] and possibly also the FM bandwidth. Moreover, devices showing a more-stable frequency-comb operation could tolerate a larger current modulation possibly resulting in a larger frequency deviation, ensuring a higher FM S/N. On the detection and demodulation side, by implementing the detection of the digital signal as modulation of the IBN around 7 GHz in place of the performed direct detection we expect a significant improvement in terms of S/N. In order to fully exploit the advantages of FM communication, a dynamic-range compressor amplifier could be added right before the DLD. This would make the demodulation schemes could be adopted in place of the DLD, e.g. based on a phase-locked loop. Finally, alternative transmission schemes able to exploit a balanced detection [61] could improve further the S/N of the communication by suppressing common-mode noise.

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Data availability. Data underlying the results presented in this paper may be obtained from the corresponding authors upon reasonable request.

Supplemental document. See Supplement 1 for supporting content.

References

- G. Modugno, M. Inguscio, and G. M. Tino, "Search for small violations of the symmetrization postulate for spin-0 particles," Phys. Rev. Lett. 81(22), 4790–4793 (1998).
- G. Modugno and M. Modugno, "Testing the symmetrization postulate on molecules with three identical nuclei," Phys. Rev. A 62(2), 022115 (2000).
- D. Mazzotti, P. Cancio, G. Giusfredi, M. Inguscio, and P. De Natale, "Search for exchange-antisymmetric states for spin-0 particles at the 10⁻¹¹ level," Phys. Rev. Lett. 86(10), 1919–1922 (2001).
- A. Shelkovnikov, R. J. Butcher, C. Chardonnet, and A. Amy-Klein, "Stability of the proton-to-electron mass ratio," Phys. Rev. Lett. 100(15), 150801 (2008).
- A. Cournol, M. Manceau, M. Pierens, L. Lecordier, D. B. A. Tran, R. Santagata, B. Argence, A. Goncharov, O. Lopez, M. Abgrall, Y. Le Coq, R. Le Targat, H. Alvarez Martinez, W. K. Lee, D. Xu, P. E. Pottie, R. J. Hendricks, T. E. Wall, J. M. Bieniewska, B. E. Sauer, M. R. Tarbutt, A. Amy-Klein, S. K. Tokunaga, and B. Darquié, "A new experiment to test parity symmetry in cold chiral molecules using vibrational spectroscopy," Quantum Electron. 49(3), 288–292 (2019).
- A. Fast and S. A. Meek, "Sub-ppb measurement of a fundamental band rovibrational transition in HD," Phys. Rev. Lett. 125(2), 023001 (2020).
- I. Kortunov, S. Alighanbari, M. Hansen, G. Giri, V. Korobov, and S. Schiller, "Proton-electron mass ratio by high-resolution optical spectroscopy of ion ensembles in the resolved-carrier regime," Nat. Phys. 17(5), 569–573 (2021).
- F. Adler, P. Masłowski, A. Foltynowicz, K. C. Cossel, T. C. Briles, I. Hartl, and J. Ye, "Mid-infrared Fourier transform spectroscopy with a broadband frequency comb," Opt. Express 18(21), 21861–21872 (2010).
- A. Foltynowicz, P. Masłowski, A. J. Fleisher, B. J. Bjork, and J. Ye, "Cavity-enhanced optical frequency comb spectroscopy in the mid-infrared application to trace detection of hydrogen peroxide," Appl. Phys. B 110(2), 163–175 (2013).
- D. A. Lashof and D. R. Ahuja, "Relative contributions of greenhouse gas emissions to global warming," Nature 344(6266), 529–531 (1990).
- 11. Harvard–Smithsonian Center for Astrophysics (CfA), V. E. Zuev Insitute of Atmosperic Optics (IAO), *HITRAN on the Web* (2019).
- 12. Harvard-Smithsonian Center for Astrophysics (CfA), The HITRAN database (2013).
- 13. L. Elterman, "Parameters for attenuation in the atmospheric windows for fifteen wavelengths," Appl. Opt. **3**(6), 745–749 (1964).

- J. Faist, F. Capasso, D. L. Sivco, C. Sirtori, A. L. Hutchinson, and A. Y. Cho, "Quantum cascade laser," Science 264(5158), 553–556 (1994).
- 15. J. Faist, Quantum cascade lasers (Oxford University Press, Oxford, UK, 2013). ISBN:9780198528241.
- M. Beck, D. Hofstetter, T. Aellen, J. Faist, U. Oesterle, M. Ilegems, E. Gini, and H. Melchior, "Continuous wave operation of a mid-infrared semiconductor laser at room temperature," Science 295(5553), 301–305 (2002).
- A. Bismuto, S. Riedi, B. Hinkov, M. Beck, and J. Faist, "Sb-free quantum cascade lasers in the 3–4 µm spectral range," Semicond. Sci. Technol. 27(4), 045013 (2012).
- R. Köhler, A. Tredicucci, F. Beltram, H. E. Beere, E. H. Linfield, A. G. Davies, D. A. Ritchie, R. C. Iotti, and F. Rossi, "Terahertz semiconductor-heterostructure laser," Nature 417(6885), 156–159 (2002).
- B. Hinkov, A. Hugi, M. Beck, and J. Faist, "Rf-modulation of mid-infrared distributed feedback quantum cascade lasers," Opt. Express 24(4), 3294–3312 (2016).
- B. Hinkov, J. Hayden, R. Szedlak, P. Martin-Mateos, B. Jerez, P. Acedo, G. Strasser, and B. Lendl, "High frequency modulation and (quasi) single-sideband emission of mid-infrared ring and ridge quantum cascade lasers," Opt. Express 27(10), 14716–14724 (2019).
- B. Hinkov, M. Beck, E. Gini, and J. Faist, "Quantum cascade laser in a master oscillator power amplifier configuration with Watt-level optical output power," Opt. Express 21(16), 19180–19186 (2013).
- L. Consolino, F. Cappelli, M. Siciliani de Cumis, and P. De Natale, "QCL-based frequency metrology from the mid-infrared to the THz range: a review," Nanophotonics 8(2), 181–204 (2018).
- S. Borri, G. Insero, G. Santambrogio, D. Mazzotti, F. Cappelli, I. Galli, G. Galzerano, M. Marangoni, P. Laporta, V. Di Sarno, L. Santamaria, P. Maddaloni, and P. De Natale, "High-precision molecular spectroscopy in the mid-infrared using quantum cascade lasers," Appl. Phys. B 125(1), 18 (2019).
- 24. I. Galli, S. Bartalini, P. Cancio, F. Cappelli, G. Giusfredi, D. Mazzotti, N. Akikusa, M. Yamanishi, and P. De Natale, "Absolute frequency measurements of CO₂ transitions at 4.3 μm with a comb-referenced quantum cascade laser," Mol. Phys. **111**(14-15), 2041–2045 (2013).
- 25. I. Galli, S. Bartalini, M. Barucci, P. Cancio, F. Cappelli, G. Giusfredi, D. Mazzotti, N. Akikusa, and P. De Natale, "Spectroscopic detection of radiocarbon dioxide at parts-per-quadrillion sensitivity," *Conf. on Lasers Electro-Optics* p. JF2K.5 (2016).
- 26. R. Santagata, D. B. A. Tran, B. Argence, O. Lopez, S. K. Tokunaga, F. Wiotte, H. Mouhamad, A. Goncharov, M. Abgrall, Y. Le Coq, H. Alvarez-Martinez, R. Le Targat, W. K. Lee, D. Xu, P.-E. Pottie, B. Darquié, and A. Amy-Klein, "High-precision methanol spectroscopy with a widely tunable SI-traceable frequency-comb-based mid-infrared QCL," Optica 6(4), 411–423 (2019).
- J. Faist, C. Gmachl, F. Capasso, C. Sirtori, D. L. Sivco, J. N. Baillargeon, and A. Y. Cho, "Distributed feedback quantum cascade lasers," Appl. Phys. Lett. 70(20), 2670–2672 (1997).
- A. Hugi, G. Villares, S. Blaser, H. C. Liu, and J. Faist, "Mid-infrared frequency comb based on a quantum cascade laser," Nature 492(7428), 229–233 (2012).
- 29. D. Burghoff, T.-Y. Kao, N. Han, C. W. I. Chan, X. Cai, Y. Yang, D. J. Hayton, J.-R. Gao, J. L. Reno, and Q. Hu, "Terahertz laser frequency combs," Nat. Photon. 8(6), 462–467 (2014).
- 30. M. Rösch, G. Scalari, M. Beck, and J. Faist, "Octave-spanning semiconductor laser," Nat. Photon. 9(1), 42-47 (2014).
- J. Faist, G. Villares, G. Scalari, M. Rösch, C. Bonzon, A. Hugi, and M. Beck, "Quantum cascade laser frequency combs," Nanophotonics 5(2), 272–291 (2016).
- S. Riedi, A. Hugi, A. Bismuto, M. Beck, and J. Faist, "Broadband external cavity tuning in the 3 4 μm window," Appl. Phys. Lett. 103(3), 031108 (2013).
- 33. S. Riedi, F. Cappelli, S. Blaser, P. Baroni, A. Müller, and J. Faist, "Broadband superluminescence, 5.9 μm to 7.2 μm, of a quantum cascade gain device," Opt. Express 23(6), 7184–7189 (2015).
- P. Friedli, H. Sigg, B. Hinkov, A. Hugi, S. Riedi, M. Beck, and J. Faist, "Four-wave mixing in a quantum cascade laser amplifier," Appl. Phys. Lett. 102(22), 222104 (2013).
- 35. J. B. Khurgin, Y. Dikmelik, A. Hugi, and J. Faist, "Coherent frequency combs produced by self frequency modulation in quantum cascade lasers," Appl. Phys. Lett. 104(8), 081118 (2014).
- N. Henry, D. Burghoff, Q. Hu, and J. B. Khurgin, "Temporal characteristics of quantum cascade laser frequency modulated combs in long wave infrared and THz regions," Opt. Express 26(11), 14201–14212 (2018).
- N. Opačak and B. Schwarz, "Theory of frequency-modulated combs in lasers with spatial hole burning, dispersion, and kerr nonlinearity," Phys. Rev. Lett. 123(24), 243902 (2019).
- D. Burghoff, "Unraveling the origin of frequency modulated combs using active cavity mean-field theory," Optica 7(12), 1781–1787 (2020).
- J. B. Khurgin, "Analytical expression for the width of quantum cascade laser frequency comb," Appl. Phys. Lett. 117(16), 161104 (2020).
- F. Cappelli, G. Campo, I. Galli, G. Giusfredi, S. Bartalini, D. Mazzotti, P. Cancio, S. Borri, B. Hinkov, J. Faist, and P. De Natale, "Frequency stability characterization of a quantum cascade laser frequency comb," Laser Photonics Rev. 10(4), 623–630 (2016).
- F. Cappelli, G. Villares, S. Riedi, and J. Faist, "Intrinsic linewidth of quantum cascade laser frequency combs," Optica 2(10), 836–840 (2015).
- 42. D. Burghoff, Y. Yang, D. J. Hayton, J.-R. Gao, J. L. Reno, and Q. Hu, "Evaluating the coherence and time-domain profile of quantum cascade laser frequency combs," Opt. Express 23(2), 1190–1202 (2015).

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- 43. F. Cappelli, L. Consolino, G. Campo, I. Galli, D. Mazzotti, A. Campa, M. Siciliani de Cumis, P. Cancio Pastor, R. Eramo, M. Rösch, M. Beck, G. Scalari, J. Faist, P. De Natale, and S. Bartalini, "Retrieval of phase relation and emission profile of quantum cascade laser frequency combs," Nat. Photon. 13(8), 562–568 (2019).
- 44. L. Consolino, M. Nafa, F. Cappelli, K. Garrasi, F. P. Mezzapesa, L. Li, A. G. Davies, E. H. Linfield, M. S. Vitiello, P. De Natale, and S. Bartalini, "Fully phase-stabilized quantum cascade laser frequency comb," Nat. Commun. 10(1), 2938 (2019).
- G. Villares, A. Hugi, S. Blaser, and J. Faist, "Dual-comb spectroscopy based on quantum-cascade-laser frequency combs," Nat. Commun. 5(1), 5192 (2014).
- 46. L. Consolino, M. Nafa, M. De Regis, F. Cappelli, K. Garrasi, F. P. Mezzapesa, L. Li, A. G. Davies, E. H. Linfield, M. S. Vitiello, S. Bartalini, and P. De Natale, "Quantum cascade laser based hybrid dual comb spectrometer," Commun. Phys. 3(1), 69 (2020).
- 47. M. S. Vitiello, L. Consolino, S. Bartalini, A. Taschin, A. Tredicucci, M. Inguscio, and P. De Natale, "Quantum-limited frequency fluctuations in a terahertz laser," Nat. Photon. 6(8), 525–528 (2012).
- 48. A. Shehzad, P. Brochard, R. Matthey, F. Kapsalidis, M. Shahmohammadi, M. Beck, A. Hugi, P. Jouy, J. Faist, T. Südmeyer, and S. Schilt, "Frequency noise correlation between the offset frequency and the mode spacing in a mid-infrared quantum cascade laser frequency comb," Opt. Express 28(6), 8200–8210 (2020).
- 49. L. Consolino, A. Campa, M. De Regis, F. Cappelli, G. Scalari, J. Faist, M. Beck, M. Rösch, S. Bartalini, and P. De Natale, "Controlling and phase-locking a THz quantum cascade laser frequency comb by small optical frequency tuning," Laser Photonics Rev. 15(6), 2000417 (2021).
- 50. K. Gürel, S. Schilt, A. Bismuto, Y. Bidaux, C. Tardy, S. Blaser, T. Gresch, and T. Südmeyer, "Frequency tuning and modulation of a quantum cascade laser with an integrated resistive heater," Photonics **3**(3), 47 (2016).
- 51. F. Capasso, R. Paiella, R. Martini, R. Colombelli, C. Gmachl, T. Myers, M. Taubman, R. Williams, C. Bethea, K. Unterrainer, H. Hwang, D. Sivco, A. Cho, A. Sergent, H. Liu, and E. Whittaker, "Quantum cascade lasers: ultrahigh-speed operation, optical wireless communication, narrow linewidth, and far-infrared emission," IEEE J. Quantum Electron. 38(6), 511–532 (2002).
- X. Pang, O. Ozolins, L. Zhang, R. Schatz, A. Udalcovs, X. Yu, G. Jacobsen, S. Popov, J. Chen, and S. Lourdudoss, "Free-space communications enabled by quantum cascade lasers," Phys. Status Solidi A 218(3), 2000407 (2020).
- 53. M. Piccardo, M. Tamagnone, B. Schwarz, P. Chevalier, N. A. Rubin, Y. Wang, C. A. Wang, M. K. Connors, D. McNulty, A. Belyanin, and F. Capasso, "Radio frequency transmitter based on a laser frequency comb," Proc Natl Acad Sci USA 116(19), 9181–9185 (2019).
- K. Faith, F. Labaar, A. Lance, and F. Mendoza, "Frequency switching speed measurements using a delay line discriminator," IEEE Trans. Instrum. Meas. 37(4), 620–625 (1988).
- 55. L. Tombez, F. Cappelli, S. Schilt, G. Di Domenico, S. Bartalini, and D. Hofstetter, "Wavelength tuning and thermal dynamics of continuous-wave mid-infrared distributed feedback quantum cascade lasers," Appl. Phys. Lett. 103(3), 031111 (2013).
- 56. L. Rothman, I. Gordon, Y. Babikov, A. Barbe, D. Chris Benner, P. Bernath, M. Birk, L. Bizzocchi, V. Boudon, L. Brown, A. Campargue, K. Chance, E. Cohen, L. Coudert, V. Devi, B. Drouin, A. Fayt, J.-M. Flaud, R. Gamache, J. Harrison, J.-M. Hartmann, C. Hill, J. Hodges, D. Jacquemart, A. Jolly, J. Lamouroux, R. Le Roy, G. Li, D. Long, O. Lyulin, C. Mackie, S. Massie, S. Mikhailenko, H. Müller, O. Naumenko, A. Nikitin, J. Orphal, V. Perevalov, A. Perrin, E. Polovtseva, C. Richard, M. Smith, E. Starikova, K. Sung, S. Tashkun, J. Tennyson, G. Toon, V. Tyuterev, and G. Wagner, "The HITRAN2012 molecular spectroscopic database," Journal of Quantitative Spectroscopy and Radiative Transfer 130, 4–50 (2013).
- 57. Atmospheric visibility is a parameter associated to the atmospheric conditions. The visibility range is the distance where the intensity of a beam of visible light decreases to 2% of its original value [58,59]. To correlate optical attenuation and visibility range several models have been developed. The one adopted here is an empirical model applying the Mie Scattering theory, according to the Kruse model [60].
- International Telecommunication Union, "Propagation data required for the design of terrestrial free-space optical links," Recommendation P.1817 (2012).
- International Telecommunication Union, "Prediction methods required for the design of terrestrial free-space optical links," Recommendation P.1814-0 (2007).
- 60. H. Kaushal, V. Jain, and S. Kar, Free space optical communication, (Springer, 2017).
- T. Gabbrielli, F. Cappelli, N. Bruno, N. Corrias, S. Borri, P. De Natale, and A. Zavatta, "Mid-infrared homodyne balanced detector for quantum light characterization," Opt. Express 29(10), 14536–14547 (2021).