

Article

Thermoelectric Generator Based on Kesterite ($\text{Cu}_2\text{ZnSnS}_4$) Synthesized via Sol–Gel Method

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Abstract

The need for sustainable and cost-effective thermoelectric materials has brought attention to earth-abundant and mineral compounds, like $\text{Cu}_2\text{ZnSnS}_4$ (CZTS). In this work, CZTS nanoparticles (NPs) were synthesized via the sol–gel method using environmentally friendly solvents based on water and ethanol mixtures. The resulting CZTS NPs were then processed into inks through ball milling to produce a thin-film thermoelectric generator (TEG). Structural and microstructural properties were investigated via X-ray diffraction and Raman spectroscopy, confirming the kesterite CZTS phase upon sintering. The chalcogenide exhibited *p*-type semiconductor behaviour, with a Seebeck coefficient reaching $\sim 69 \mu\text{V}/\text{K}$ at 385 K. Van-der-Pauw measurements of conductivity confirmed a non-degenerate semiconducting behaviour, achieving $\sim 1.77 \text{ S}/\text{cm}$ at 323 K. A two-leg CZTS thin-film TEG reaching a maximum power output of 32(3) nW at a $\Delta T \sim 160 \text{ K}$ was used, measured with a home-made setup. The volume-specific power of the generator reached $4 \times 10^{-4} \mu\text{W cm}^{-3} \text{ K}^{-2}$. These results point to an effective use of sol–gel-based techniques to produce a functional thermoelectric generator, providing a costless and environmentally friendly approach to CZTS NPs.

Keywords: CZTS; thermoelectric generator; sol–gel; ball milling; doctor blade; Seebeck coefficient; power factor



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1. Introduction

Kesterite $\text{Cu}_2\text{ZnSnS}_4$ (CZTS) is a quaternary semiconductor that was initially developed as a sustainable alternative to $\text{Cu}(\text{In,Ga})\text{Se}_2$ absorbers in thin-film photovoltaic devices [1]. In contrast to CIGS, CZTS is composed exclusively of earth-abundant and environmentally benign elements (Cu, Zn, Sn, and S) [2], exhibiting a high absorption coefficient ($\alpha > 10^4 \text{ cm}^{-1}$) and a near-optimal direct band gap of approximately 1.5 eV [3–6]. For this reason, recent advances in solution processing and interface engineering have led to CZTS-based solar cells with efficiencies reaching 15.1% [7]. Beyond photovoltaics, CZTS belongs to the broader family of copper-based chalcogenides investigated as efficient and potentially low-cost thermoelectric materials [8–12], offering a more sustainable alternative to established systems such as Bi_2Te_3 , Sb_2Te_3 , PbTe and GeTe , that depend on scarce and

expensive elements [13]. This interest is further supported by the fact that several key attributes of CZTS—its diamond-like structure, tuneable electronic band structure, and intrinsically low lattice thermal conductivity—are also advantageous for thermoelectric energy conversion [13,14].

The thermoelectric performance of a material is quantified by the dimensionless figure of merit

$$zT = \frac{S^2\sigma}{\kappa} T \quad (1)$$

where S is the Seebeck coefficient ($\mu\text{V K}^{-1}$), σ the electrical conductivity (S cm^{-1}), κ the total thermal conductivity ($\text{W m}^{-1} \text{K}^{-1}$), and T (K) the absolute temperature [15,16].

Initial efforts to maximize zT focused on bulk CZTS synthesized by mechanochemical routes. Baláž et al. [17] demonstrated the semi-industrial synthesis of Cu-based chalcogenides using an eccentric vibratory mill, reporting a maximum zT of 0.04 at 673 K for bulk CZTS. Subsequent studies showed that thermoelectric performance can be optimized through control of cation disorder, particularly by promoting the partially disordered tetragonal kesterite phase (space group $I-42m$) [18,19], which modifies carrier scattering and transport properties. By this route, further disorder can be achieved by stabilizing the cubic sphalerite phase ($F-43m$) with complete cation disorder [19–23].

More recently, attention has shifted toward thin-film thermoelectric devices, where optimization of the power factor ($S^2\sigma$) is particularly critical. For example, Kumar et al. [24] deposited CZTS thin films by chemical vapour deposition at different temperatures and reported a significant increase in power factor from 2 to 7 $\mu\text{W K}^{-2} \text{m}^{-1}$ with increasing deposition temperature. This enhancement was attributed to changes in carrier concentration associated with secondary phase formation.

Thin-film devices based on CZTS and related kesterites have demonstrated promising performance per unit cost, highlighting their potential for scalable thermoelectric applications [13]. Various deposition techniques have been explored for CZTS thin films, including sputtering, hot injection, ball milling-based ink processing, and spin coating [14,25]. While sputtering yields high phase purity and excellent stoichiometric control, it involves complex and costly equipment. Alternative solution-based approaches offer advantages in scalability and cost-effectiveness, though they often face challenges related to phase control, contamination, and particle size distribution [14,25].

Among these approaches, the sol–gel method represents a particularly attractive route due to its simplicity, compositional flexibility, and ability to produce homogeneous thin films from low-cost precursors [26]. Importantly, sol–gel processing enables precise control over stoichiometry and defect chemistry, which are key parameters governing charge transport in CZTS [27]. Previous studies have shown that adjusting sulphur content improves both the Seebeck coefficient and electrical conductivity by enhancing phase purity and reducing detrimental sulphur-related defects. Similarly, tuning the Sn precursor concentration has been demonstrated to significantly modify carrier concentration and transport behaviour, highlighting the strong influence of cation stoichiometry on intrinsic defect formation and scattering mechanisms [28]. In addition to compositional control, optimization of film thickness and microstructure through spin-coating strategies has been reported to enhance grain connectivity and reduce interfacial scattering, thereby improving the power factor [29]. Collectively, these findings demonstrate that sol–gel processing provides a versatile platform for tailoring composition and microstructure to optimize the thermoelectric performance of CZTS thin films.

The aim of this work is to fabricate a TE device based on sustainable CZTS material. We adopt a simple and environmentally benign sol–gel strategy using distilled water and ethanol as eco-friendly solvents for the synthesis of CZTS nanoparticles. These nanoparticles are subsequently formulated into inks for thin-film fabrication and thermoelectric device assembly. The structural, morphological, and electrical properties of the resulting CZTS-based thermoelectric generators are systematically investigated, with emphasis on the relationship between processing conditions, microstructure, and thermoelectric performance.

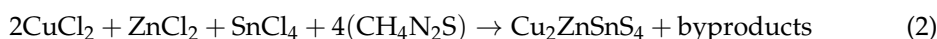
2. Experimental Section

2.1. Materials

Copper (II) chloride dihydrate ($\text{CuCl}_2 \cdot 2\text{H}_2\text{O}$; 99% pure), zinc chloride ($\text{ZnCl}_2 \cdot 2\text{H}_2\text{O}$; 98% pure), tin (II) chloride dihydrate ($\text{SnCl}_2 \cdot 2\text{H}_2\text{O}$; 98% pure), thiourea ($\text{CH}_4\text{N}_2\text{S}$; 99%), Oleylamine (70% pure), sulphur (99.5% pure), and ethanol (96%) were obtained from Sigma-Aldrich (St. Louis, MO, USA). Distilled water was used for the synthesis.

2.2. Synthesis of CZTS NPs and TEG Fabrication

The experimental procedure utilized chloride precursors of copper, zinc, and tin, alongside thiourea as the sulphur source. To prevent precipitation during the process and ensure the preparation of a stoichiometric $\text{Cu}_2\text{ZnSnS}_4$ solution, a 50% excess of sulphur was employed. Each precursor was dissolved separately in a water–ethanol (70–30%) mixture for 15 min. Subsequently, all four precursors were mixed under magnetic stirring at 50 °C for 2 h, resulting in a homogeneous yellow solution. The CZTS solution was filtered using a glass funnel to obtain CZTS NPs. The filtered material was then dried in an oven at 150 °C for 1 h. Following the drying process, the CZTS NPs were ground into a fine powder and subjected to sulfurization at 350 °C for 4 h under an inert atmosphere. The overall synthesis reaction followed the equation below:



The CZTS inks were prepared by combining 0.5 g of CZTS NPs with 2.5 mL of Oleylamine. The mixture was milled in a Fritsch P4 planetary ball mill (Fritsch, Laufen, Germany) for 15 min, with a main disk rotation speed of the main disk of 300 rpm (main-disk-to-planet ratio of -1.8). The solvent was separated from the NPs using a centrifugation process (MPW-380, MED-Instruments, Jacksonville, FL, USA) at 12,000 rpm for 10 min, after washing the mixture with an ethanol/toluene solution (12.5 mL:0.5 mL). Subsequently, an additional 0.5 mL of toluene was introduced, and the resulting solution was subjected to a sonication bath for 15 min to ensure thorough dispersion.

The CZTS thin-films were fabricated using the CZTS inks. First, soda-lime glass (SLG) substrates were cleaned. Then, a mask with specific dimensions was prepared using Kapton tape. Next, 0.1 mL of CZTS ink was deposited onto the substrates using a Finn pipette (Thermo Fisher Scientific, Waltham, MA, USA) to form the thin films by the doctor blade deposition method. Finally, the thin films underwent sulfurization at 400 °C for 1 h with a heating rate of 3 °C/min under a controlled atmosphere (N_2), with ~85 mg of sulphur placed next to the sample. The thin films were patterned as shown in Figure 1a, connected by thermally evaporated Ag thin-film contacts (see [8] for details).

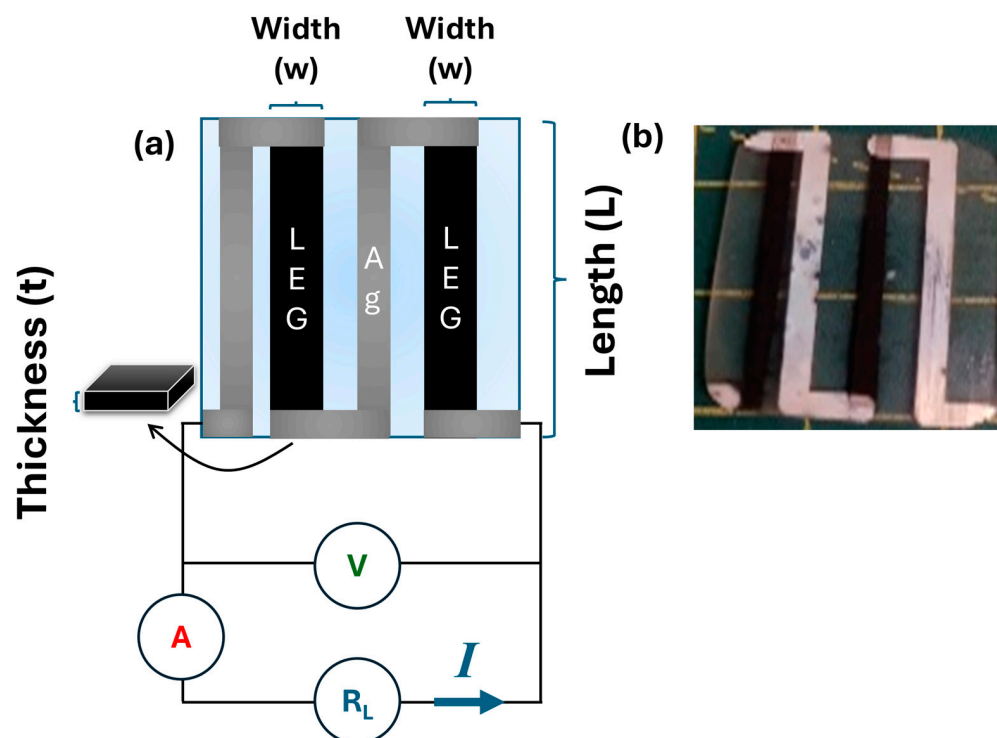


Figure 1. (a) Schematic representation of the thermoelectric generator (TEG) and the electrical circuit employed to acquire the current–voltage ($I - V$) characteristics. R_L denotes the external load resistance, A the ammeter used to measure the current, and V the voltmeter used to record the voltage. The dimensions of thickness, length, and width are indicated by the brackets and arrows. (b) Photograph of the CZTS-based TEG device used for the electrical measurements.

2.3. Characterization Techniques

The phase of the as-prepared CZTS NPs and thin films was investigated by X-ray diffraction (XRD Bruker D8 Discover, Bruker, Billerica, MA, USA) with $\text{CoK}\alpha$ radiation $\lambda = 1.78901 \text{ \AA}$ in Bragg–Brentano geometry. Phase identification used the PDF4+ database [30] and Rietveld analysis was performed using TOPAS 7 software [31,32]. Whole-powder pattern modelling employed the macros developed by Scardi et al. [33–35] to estimate crystallite size and distribution. The fundamental parameters approach was employed to deconvolve the instrumental resolution function [36]. Baseline fitting employed Chebyshev functions.

The Raman measurements were performed using a LabRAM Aramis (Horiba Jobin-Yvon, Palaiseau, France), which has an optical microscope and $100\times$ objective, with a 532 nm source wavelength. The conductivity was estimated using Van der Pauw measurements performed with a Linseis HCS-1 instrument (Linseis, Selb, Germany).

For electrical characterization, a specifically shaped CZTS thin film with square dimensions of $12.5 \text{ mm} \times 12.5 \text{ mm}$ with a thickness of 0.22 mm was fabricated, following the same conditions for the TEG deposition. Ohmic contact verification was performed before the measurement, varying the current from 0.1 to 1 mA for all 4 spring-load contacts, yielding a linear relationship between current and measured voltage. The magnetic field range was from -0.7 to 0.7 T , with measurements performed in a controlled N_2 atmosphere at an applied current of 1 mA. The thickness of the samples was measured by a Dektak3 Stylus Profilometer (Veeco Instruments, Plainview, NY, USA).

The TEG's power output was measured using a home setup to estimate the current–voltage–power curves ($I - V - P$). It consists of two thermocouples, one manually attached to two micro-heaters in the hot side, and another in the Peltier cooler to measure the cold side temperature. Upon generation of a temperature gradient by the micro-heaters,

a load resistance connected in series with the TEG was varied from 0 to 2 M Ω using a newly developed automated system [8]. The voltage (V) and current (I) were measured with a Keithley 2601A source meter (Keithley Instruments Inc., Cleveland, OH, USA), connected simultaneously in parallel and in series, as represented in the circuit of Figure 1a. Calibration of the resistance applied was estimated by Ohm's law ($R = V/I$). The power output was evaluated according to the standard relation: $P_{\text{OUT}} = V \times I = R \times I^2$. To ensure a stable acquisition of current and voltage values, the measurements were repeated three times. The uncertainty was estimated from the instrument precision specified by the manufacturer.

3. Results and Discussion

3.1. Structural and Microstructural Characterization

The XRD pattern of the as-prepared CZTS nanoparticles in Figure 2a confirms that the dominant crystalline phase corresponds to tetragonal kesterite CZTS (space group 121, $I-42m$, PDF #00-026-0575), representing approximately 73(3) wt% of the sample. A secondary phase identified as cassiterite SnO₂ (space group 136, $P4_2/mnm$, PDF #01-070-4176) is also detected. The occupation of Cu atoms at the $x = 0$, $y = 1/2$, and $z = 1/4$ positions and Sn atoms was refined, yielding values of 0.15 and 0.6, respectively, indicating that the CZTS structure is Cu- and Sn-deficient. The refined lattice parameters are $a = 5.390(1)$ Å and $c = 10.648(4)$ Å and the crystallite size is 9(2) nm on average with a standard deviation of 4(1) for a log-normal distribution as shown in the inset of Figure 2a. The value of the microstrain is $\epsilon_0 = 0.15(2)$ %.

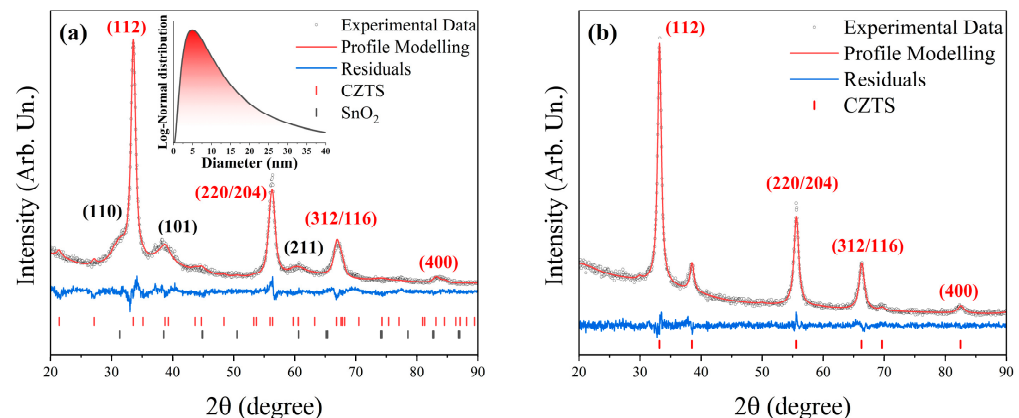


Figure 2. (a) XRD pattern and profile modelling of the CZTS NPs and (b) XRD pattern and profile modelling of the CZTS thin film. The black open dots represent the measured data, the red line the profile modelling, and the residuals are represented in blue. The black markers correspond to the Miller indices of the CZTS phase and the red ones to the SnO₂ phase.

After ink preparation, washing, and sulfurization, the XRD pattern of the CZTS legs, shown in Figure 2b, no longer exhibits detectable SnO₂ reflections. The reduction of the oxide phase may originate from multiple mechanisms. First, SnO₂ nanoparticles are known to interact with long-chain amine ligands such as Oleylamine (OLA), which can coordinate to metal-oxide surfaces and promote dispersion in organic media; ligand-assisted stabilization and phase transfer of SnO₂ nanocrystals in nonpolar solvents has been reported in the literature [37]. Such interactions may facilitate preferential removal of oxide-rich fractions during solvent washing and centrifugation. Second, SnO₂ can undergo sulfurization at moderate temperatures (≈ 350 – 450 °C), leading to the formation of SnS₂ under sulphur-rich conditions [38]. In multinary chalcogenide systems, transient SnS₂ has been reported to react further with Cu–S and Zn–S intermediates during annealing,

contributing to the formation of kesterite CZTS [39]. Therefore, the disappearance of SnO₂ reflections after sulfurization may result from a combination of partial physical removal during ink purification and chemical transformation during annealing. No additional crystalline secondary phases are observed within the detection limit of XRD, indicating that any intermediate sulphide species, if formed, are either consumed during the reaction or present below the instrumental sensitivity threshold.

Raman spectroscopy provides complementary information on short-range order and minor phases. The nanoparticle (NP) spectrum (Figure 3a) is dominated by the characteristic A₁ mode of kesterite CZTS, observed at ~332 cm⁻¹, slightly downshifted from the typical 337–338 cm⁻¹ position [21,40,41]. This shift is consistent with nanoscale effects and cation disorder, in agreement with the refined crystallite size and microstrain [11]. A weak feature near ~195 cm⁻¹ further supports the presence of the kesterite phase. A small band at ~471 cm⁻¹ suggests a minor Cu_{2-x}S contribution, which may remain undetected by XRD due to its low concentration or limited crystallinity [24]. A weak band at ~642 cm⁻¹ is consistent with Sn–O-related vibrations [42], corroborating the SnO₂ fraction observed in the NP XRD pattern.

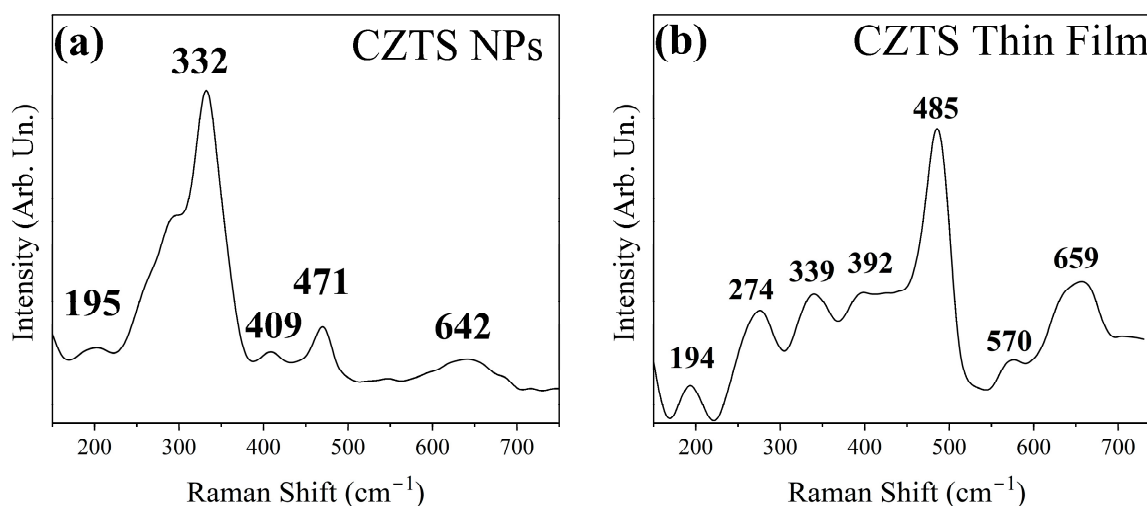


Figure 3. Raman spectrum of CZTS NPs (a) and CZTS thin film (b).

The Raman spectrum of the CZTS leg, in Figure 3b, exhibits kesterite-related modes at ~194, ~339, and ~392 cm⁻¹, indicating improved crystallinity and structural ordering after annealing [41]. Additional bands at ~570 and ~659 cm⁻¹ fall within the second-order (multiphonon) scattering region commonly reported for CZTS [43]. A pronounced band at ~485 cm⁻¹ appears after processing. Since this feature does not correspond to the most commonly reported Raman signatures of ZnS, CTS, or SnS_x phases, it is tentatively attributed to disorder-activated or defect-related lattice vibrations, or possibly to local structural modifications induced during sulfurization or laser excitation. Further power-dependent Raman measurements would be required to conclusively determine its origin.

Importantly, the Raman spectrum of the thin film shows an intense broad band centred near ~1400 cm⁻¹ (see Figure S1 of the Supporting Information File). This feature is characteristic of organic ligands, such as OLA, associated with CH₂ bending and C–C vibrational modes [44]. The high Raman activity of organic species explains their strong spectral contribution, even when present in relatively small amounts. In contrast, XRD is largely insensitive to such organic residues due to their low electron density and poor crystallinity. Therefore, while XRD confirms the predominance of crystalline kesterite CZTS after processing, Raman spectroscopy reveals that residual organic ligands remain in the film and contribute significantly to the vibrational spectrum.

Overall, the combined XRD and Raman analyses indicate that the sulfurized films are predominantly kesterite CZTS, with oxide impurities substantially reduced after ink preparation and annealing. Minor sulphide-related species and residual organic ligands may persist but remain below the detection limit of XRD.

3.2. Electrical Conductivity

The electrical conductivity (σ) of the CZTS thin film is shown in Figure 4. The overall variation in conductivity between room temperature and 200 °C is limited to approximately 3%, which is within the uncertainty range of the experimental measurement (5%). Above approximately 150 °C in the heating curves, a slight decrease in electrical conductivity is observed, which may be associated with disorder–order transitions in kesterite-type materials, as previously reported by Isotta et al. [45]. Upon cooling, the kinetics of this phase transformation are not observed, giving a constant value of conductivity throughout the entire temperature range. The results were compared with those reported for other chemical routes. The hot-injection ink-based TEGs developed by Syafiq et al. [14] reported an electrical conductivity of 0.28 S cm⁻¹; the present CZTS thin film shows an approximately sixfold-higher value. However, Sharma and Neeleshwar observed higher electrical conductivity for CZTS-based TE materials reaching up to 2.1 S cm⁻¹ [36]. Given the strong dependence of σ on the concentration of defects and on microstructural features, variations arising from different synthesis routes are to be expected, even in stoichiometric CZTS.

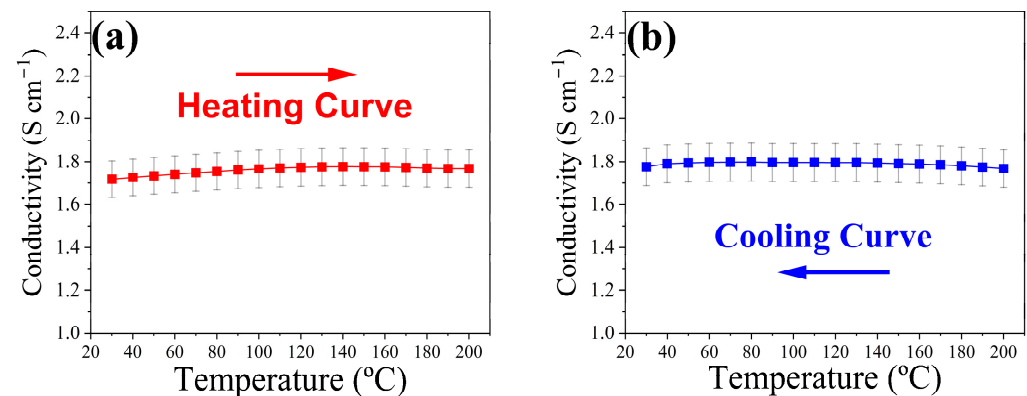


Figure 4. Electrical conductivity of the CZTS thin film measured in the Van der Pauw configuration during (a) heating and (b) cooling cycles.

3.3. Performance Characteristics of the CZTS-Based TEG

Since the film displayed a measurable conductivity, it was employed as a thermoelectric generator, as shown in Figure 1. The thermoelectric performance of the device was evaluated using a custom-built measurement setup described elsewhere [8], in which both current and voltage were recorded simultaneously under controlled temperature gradients. Heating was applied to the hot side (T_{hot}) using resistive heaters positioned inside an insulated chamber, while the cold side was stabilized by a Peltier module (T_{cold}) to ensure a constant reference temperature. Before each measurement, sufficient time was allowed to establish a stable thermal gradient across the device. T_{hot} was varied from 323 K up to 473 K while maintaining the T_{cold} 293 K.

Figure 5a shows the current–voltage (I–V) characteristics obtained by systematically varying the external load resistance between 0 and 2 M Ω in discrete steps. The resulting I–V curves exhibit a linear relationship, confirming the Ohmic behaviour of the device and the absence of significant contact resistance. The internal resistance (R_{int}) of the TEG was extracted from the slope of the I–V curves ($R_{int} = -\Delta V/\Delta I$), yielding values on the order of 4 k Ω (see inset of Figure 5a). This internal resistance includes both intrinsic material

resistance and contact contributions and therefore represents the effective resistance of the thermoelectric leg. Two limiting conditions define the electrical response: the open-circuit voltage (V_{OC}), measured when the external resistance approaches infinity and the current is zero, and the short-circuit current (I_{sc}), measured when the external resistance is zero and the voltage vanishes.

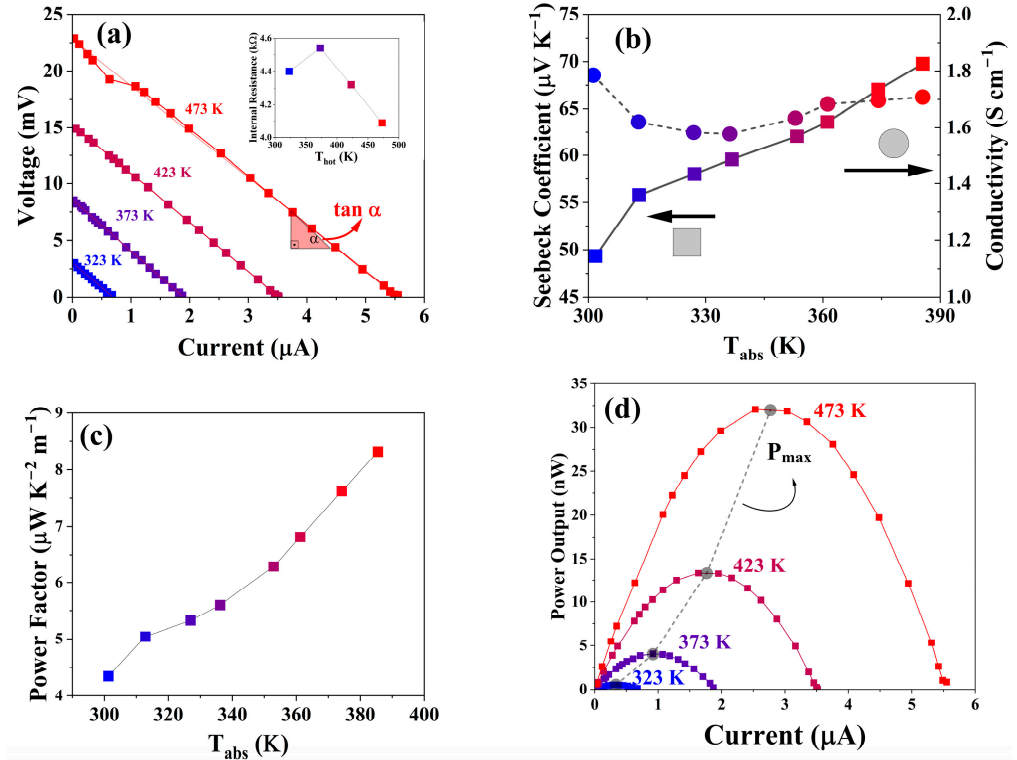


Figure 5. (a) I-V plots, (b) Seebeck (squares) and conductivity (spheres), (c) power factor, (d) power output obtained at different temperatures for the TEG. The maximum power output, P_{max} , is indicated by the dashed lines.

The electrical conductivity (σ) of the CZTS thin film was recalculated from the internal resistance using the geometrical dimensions of the device according to $\sigma = L/(R_{int} A)$, where L is the leg length and A the cross-sectional area determined from the film width and thickness (values are reported in Table 1). The temperature dependence of the conductivity is shown in Figure 5b, exhibiting values similar to those of the Van-der-Pauw method in Figure 4. The conductivity can be considered stable through the entire range of temperature, given the intrinsic error of the setup instrument. This reveals that the contact resistance of the TEGs is minimal and does not deteriorate the electronic properties during the measurement.

Table 1. TEG dimensions. t corresponds to the thickness, w to the width, L to the length, and A to the area of the device.

	t [μm]	w [cm]	L [cm]	A [cm^2]
Leg A	94	0.143	2.5	1.35×10^{-3}
Leg B	4.7	0.138	2.5	6.45×10^{-5}

The effective Seebeck coefficient was determined from the open-circuit voltage using $S = (1/N)(V_{OC}/\Delta T)$, where N is the number of thermoelectric legs—2 in this case, and $\Delta T = T_{hot} - T_{cold}$ is the applied temperature difference. The Seebeck coefficient increases

from approximately $50 \mu\text{V K}^{-1}$ at 323 K to nearly $70 \mu\text{V K}^{-1}$ at 385 K, confirming the *p*-type semiconducting nature of CZTS through its positive sign (see Figure 5b). Although these values are lower than those reported for mechanically alloyed or cold-pressed CZTS materials [14,18], they remain consistent with thin-film kesterite systems.

The power factor ($PF = S^2\sigma$) was calculated to assess the ability of the CZTS film to generate useful electrical power and is given in Figure 5c. A gradual increase in PF with temperature is observed, reaching a maximum of approximately $8 \text{ nW cm}^{-1} \text{ K}^{-2}$ at the highest measurement temperature. While this value is lower than those reported for optimized bulk kesterite and CZTSe materials, it reflects the effect of the organic ligand, OLA, which is attributed to decreasing the conducting channels throughout the material. The relatively modest *PF* indicates that further defect engineering or carrier concentration optimization would be required to enhance performance.

The electrical power output was calculated from $P = VI$ for each applied load resistance and is shown in Figure 5d. The power–current ($I - P$) curves exhibit the expected parabolic behaviour, with maximum power achieved when the external load resistance equals the internal resistance of the device. This condition is consistent with classical impedance matching in thermoelectric generators. The maximum power can also be estimated from $P_{max} = VOC I_{sc}/4$, reflecting the quadratic dependence of power on current. For the present CZTS TEG, a maximum output power of approximately 32(3) nW was obtained at the highest applied temperature gradient. Notably, the maximum power increases approximately with the square of the temperature difference ($P_{max} \propto \Delta T^2$), as expected from the proportionality of both *VOC* and *I_{sc}* to ΔT . This quadratic scaling confirms stable thermoelectric behaviour without anomalous transport effects or contact limitations.

4. Discussion

To enable a meaningful comparison with literature reports on thermoelectric generators (TEGs), the maximum power output must be normalized with respect to both the applied temperature gradient and the characteristic length scale of the device, as discussed in [9,13]. Since the maximum output power scales approximately with ΔT^2 , performance comparisons are most appropriately carried out using ΔT^2 -normalized metrics. Two complementary normalization strategies are commonly employed: (i) normalization by the total device volume, yielding the volumetric temperature power density (P_{volume} , in $\text{W cm}^{-3} \text{ K}^{-2}$), and (ii) normalization by the active surface area, yielding the areal temperature power density (P_{area} , in $\text{W cm}^{-2} \text{ K}^{-2}$).

When applying this normalization to CZTS- and CZTSe-based thin-film TEGs reported in the literature, clear trends emerge in both volumetric and areal performance. For the ink-processed devices reported by Syafiq et al. [14] at $\Delta T = 160 \text{ K}$, the ΔT^2 -normalized volumetric power densities (P_{volume}) are $8.4 \times 10^{-3} \mu\text{W cm}^{-3} \text{ K}^{-2}$ for CZTS via hot injection, called HI-CZTS, (2.0 μm thick), $1.53 \times 10^{-2} \mu\text{W cm}^{-3} \text{ K}^{-2}$ for CZTSe via ball-milling, called BM-CZTSe (4.8 μm thick), $2.19 \times 10^{-3} \mu\text{W cm}^{-3} \text{ K}^{-2}$ for BM-CZTSSe (4.1 μm thick), and $4.53 \times 10^{-3} \mu\text{W cm}^{-3} \text{ K}^{-2}$ for BM-CTS (5.1 μm thick), confirming CZTSe as the volumetric champion within that study. In terms of areal-normalized performance (P_{area}), the same devices yield 1.68×10^{-3} , 7.34×10^{-3} , 8.98×10^{-4} , and $2.30 \times 10^{-3} \text{ nW cm}^{-2} \text{ K}^{-2}$ for HI-CZTS, BM-CZTSe, BM-CZTSSe, and BM-CTS, respectively [14]. A higher level of integration and improved contact engineering was reported by Isotta et al. [13], whose CZTS and CZTSe planar thin-film TEGs ($\Delta T = 150 \text{ K}$) exhibit P_{volume} values of 2.72×10^{-2} and $4.13 \times 10^{-2} \mu\text{W cm}^{-3} \text{ K}^{-2}$, and P_{area} values of 6.80×10^{-3} and $1.24 \times 10^{-2} \text{ nW cm}^{-2} \text{ K}^{-2}$, respectively, representing a substantial improvement over early ink-based architectures. Overall, CZTSe devices consistently outperform CZTS counterparts in both volumetric and areal metrics, reflecting their higher electrical conductivity and optimized

microstructure. Nevertheless, the differences between P_{volume} and P_{area} rankings highlight the strong influence of film thickness and device geometry: while volumetric normalization favours thinner, better-integrated films, areal normalization emphasizes lateral device design and scalability. These comparisons demonstrate that, although absolute performance remains below bulk thermoelectric benchmarks, thin-film CZTS and CZTSe TEGs fabricated through solution-based or low-cost routes already achieve competitive ΔT^2 -normalized power densities within the emerging class of sustainable chalcogenide generators.

In contrast, the present CZTS TEG, fabricated via a low-temperature sol-gel route followed by simple doctor-blade deposition and sulfurization, achieves a maximum power output of approximately 35 nW at $\Delta T \sim 160$ K. Normalizing by volume and temperature gradient squared, $P_{volume} = 4 \times 10^{-4} \mu\text{W cm}^{-3} \text{K}^{-2}$, while by active surface area it achieves $P_{area} = 3.5 \times 10^{-4} \text{nW cm}^{-2} \text{K}^{-2}$. Overall, while CZTSe and optimized CTS systems may deliver superior volumetric performance [14], CZTS thin films processed via cost-effective chemical routes demonstrate that functional nanowatt-level thermoelectric generation can be achieved without vacuum-based deposition or complex device architectures. From a materials sustainability and manufacturing perspective, the simplicity of the sol-gel approach represents an advantage, even if further optimization of carrier concentration and defect engineering is required to approach the highest reported device efficiencies.

5. Conclusions

In this work, CZTS thin films were successfully fabricated through a fully solution-based and environmentally benign route combining sol-gel synthesis of nanoparticles, ink formulation via ball milling, and doctor-blade deposition followed by sulfurization. The following results were achieved:

- Structural analysis confirmed the formation of predominantly tetragonal kesterite CZTS, while XRD and Raman spectroscopy evidenced the substantial reduction of secondary SnO_2 phases after ink processing and annealing.
- Raman analysis further indicated improved structural ordering after sulfurization, although minor ligand residues and defect-related vibrational features remain detectable.
- Electrical transport measurements revealed stable *p*-type semiconducting behaviour, with a Seebeck coefficient reaching $\sim 69 \mu\text{V K}^{-1}$ at 385 K and electrical conductivity up to 1.77 S cm^{-1} .
- The resulting two-leg thin-film thermoelectric generator delivered a maximum output power of 32(3) nW under the applied temperature gradient, demonstrating reliable Ohmic behaviour and predictable ΔT^2 scaling.
- The generator achieves competitive ΔT^2 -normalized volumetric and areal power densities compared with previously reported CZTS-based thin-film TEGs fabricated through more complex or vacuum-assisted approaches, when normalized to device geometry and temperature gradient.

Although the absolute power factor remains modest relative to optimized bulk or CZTSe-based systems, the present results demonstrate that functional thermoelectric generation can be achieved using low-temperature, scalable, and low-cost processing routes based entirely on earth-abundant elements. The simplicity of the sol-gel/ink-based strategy represents a significant advantage for sustainable thermoelectric device manufacturing. Further improvements are expected through optimization of carrier concentration, defect engineering, ligand removal, and microstructural densification. Overall, this study confirms that solution-processed CZTS thin films constitute a viable platform for environmentally friendly thermoelectric generators and provides a practical pathway toward scalable and sustainable chalcogenide-based energy harvesting devices.

Supplementary Materials: The following supporting information can be downloaded at <https://www.mdpi.com/article/10.3390/ma19101971/s1>. Figure S1: Raman spectra for the kesterite thin film (red line) and glass substrate (SiO₂, given by the black line).

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